

AUSTRALIAN ATOMIC ENERGY COMMISSION

RESEARCH ESTABLISHMENT LUCAS HEIGHTS

NAIAD - A COMPUTER PROGRAM FOR CALCULATION OF THE STEADY STATE AND TRANSIENT BEHAVIOUR (INCLUDING LOCA) OF COMPRESSIBLE TWO-PHASE COOLANT IN NETWORKS

bу

G.D. TRIMBLE W.J. TURNER

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ABSTRACT

The three one-dimensional conservation equations of mass, momentum and energy are solved by a stable finite difference scheme which allows the time step to be varied in response to accuracy requirements.

Consideration of numerical stability is not necessary. Slip between the phases is allowed and descriptions of complex hydraulic components can be added into specially provided user routines. Intrinsic choking using any of the nine slip models is possible. A pipe or fuel model and detailed surface heat transfer are included.

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1. INTRODUCTION

Computer codes for simulation of loss-of-coolant accidents (LOCA) in nuclear power reactors have been written at many reactor research establishments throughout the world. All these codes suffer from one or more of the following problems.

1.1 Very Short Time Steps

In most codes, the time step is limited in size by numerical instability; usually the time step must be less than the time for a pressure or sound wave to cross one segment in the finite difference mesh. Because of the high speed of sound, particularly in sub-cooled water, this time step restriction can lead to very long computation times.

In 1971, Porsching et al. published an implicit finite difference method which allowed larger time steps for some calculations. This method was first incorporated in the FLASH-4 code [Porsching et al. 1969] and, more recently, in RELAP-4 [Moore & Rettig 1973] and FIREBIRD [Tong 1975]. In practice, the method has not removed the time step problem [Spinks 1975, Sullivan 1975].

The finite difference scheme used in NAIAD has been found stable for all calculations to date and the time step size can be chosen on accuracy considerations alone. This finite difference method is a further development of Turner [1969, 1972] and Turner & Trimble [1974(a)].

1.2 Slip Not Allowed

It is well known that in most two-phase flows the two phases move at different velocities. In such situations, flow (or mass velocity weighted) enthalpy

$$H = \langle uhd \rangle/G$$

is larger than static (or density weighted) enthalpy

$$H_s = < hd >/\rho$$
 .

Thus, enthalpy storage and flow cannot both be correct if slip is not allowed for. Nine slip models are available in NAIAD and others may easily be added.

1.3 Inadequate Momentum Flux Treatment

Several codes neglect this term of the momentum equation [Rettig et al. 1970; Porsching et al. 1969]. The treatment of Moore et al. [1973] has been incorporated in RELAP-4 [Moore & Rettig 1973]. However, this model (homogeneous flow and thermodynamic equilibrium) was shown by

Isbin et al. [1962] to give far too large an acceleration pressure drop for nearly choked two-phase flows, the flows at which momentum flux is important. The AECL code RODFLOW [Elliott 1968] satisfactorily includes momentum flux. The treatment in NAIAD gives reasonable agreement with the pressure drop measurement of Fauske [1962] for nearly choked flow [Turner & Trimble 1974(b)].

1.4 Inconsistent Choke Flow Treatment

The assumption of thermodynamic equilibrium and the specification of a slip model fully determine choke flow rate as a function of throat pressure and quality. In most codes both assumption and specification are made hence a choke flow function is fully determined. choke flow rate will only agree with experiment if the slip model is correct and the assumption of thermodynamic equilibrium holds. codes, the slip model assumed is homogeneous flow which is well known to give choke flow rates and pressure drops for near choke flow (those for which the acceleration pressure drop dominates) which do not agree with measured values. In these codes a choke flow rate function calculated using a more realistic flow model is applied as a boundary condition, and near choke flow is avoided as far as possible in the calculation. Thus one slip model is assumed everywhere except at the throat which requires another slip model. In NAIAD, the chosen slip model is used both to determine the pressure gradient near the throat and to calculate the choke flow rate. The fact that such a procedure can give agreement with steady state measurements of both choke flow rate and near choke pressure gradients was shown by Nahavandi & Von Hollen [1965], Turner & Trimble [1974(b)] and Tremblay & Andrews [1971].

GENERAL DESCRIPTION

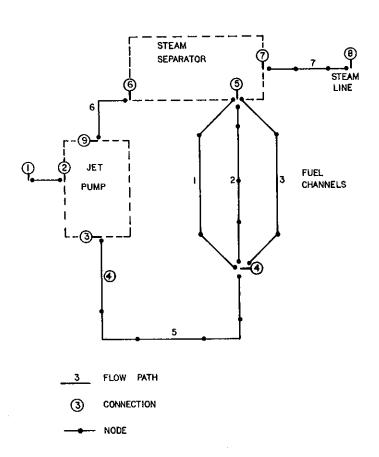
The NATAD code solves the one-dimensional equations of conservation of mass, momentum and energy by means of a stable finite difference scheme. The code is not specific to any reactor type or, indeed, to any particular pump, steam generator or other hydraulic component. The components required for any problem are, if possible, described in the input data in terms of variable area heated flow paths. The description of these flow paths requires specification of geometric parameters and selection of friction, slip and heat transfer options. Any components that cannot be so described must be described in special FORTRAN routines which are written by the code user for his particular problem or class of problems.

Thus a NATAD run consists of two parts. In the first part, the special routines written for the problem in hand are compiled and linked with the standard code to form a composite code. For many problems, the standard code suffices and this step can be omitted. In the second part, the required calculation is done using the composite code. The calculations that can be done with composite codes are described in Sections 2 and 3. The instructions for writing the special routines and using the composite codes are contained in Sections 4 to 11.

2.1 Geometry

The NAIAD code may be used to calculate steady state and transient behaviour in complex flow networks containing a compressible fluid in both single and two-phase states. In the calculation, the flow network is treated as a set of one-dimensional flow paths each of which begins and ends at a connection, e.g. the flow network shown in Figure 1 can be treated as seven flow paths and nine connections. Nodes are placed at the start, finish and at any required intermediate positions in each flow path so that an implicit finite difference method can be used for

FIGURE 1
Connections and Flow Paths for a BWR



each flow path. The only limits on network topology result from upper limits of 10, 10 and 49 on the number of connections, flow paths and nodes respectively. A possible arrangement for a Boiling Water Reactor (BWR) is shown in Figure 1.

There are four types of connections: internal, choked, explicit and implicit. Internal connections are points at which two or more flow paths join. Choked connections are placed at points where choked flow is expected to occur, e.g. at a pipe burst or pressure vessel crack. Explicit and implicit connections are placed at boundaries of the network, e.g. connections 1 and 8 in Figure 1, and at interfaces between complex hydraulic components and the network, e.g. connections 2, 3, 9, 5, 6 and 7 in Figure 1.

The behaviour at connections is governed by the connection equa-The standard code includes connection equations for choked and internal connections, and for explicit connections at which either the mass flow rate or pressure is constant. The connection equations describing the behaviour of more complex hydraulic components are not in the standard code; they must be put into finite difference form and coded into the special routines mentioned above. If these difference equations are explicit, then the connections on the interface are all specified as explicit in the input data. If the difference equations are implicit, then at least one interface connection must also be specified as implicit so that the difference equations can be solved iteratively. The sample problem (Figure 6) contains connections of all Special routines written for this problem represent a pump between two connections, and a pressure controller which varies the inlet flow at one connection in response to pressure at another connection.

The walls of flow paths in the network may be heated as in a reactor core, or cooled as in heat exchangers or steam generators. A single node fuel or pipe wall model is coupled to each hydraulic node. Heat energy flow is to or from this node.

2.2 <u>Initial Condition Calculations</u>

The initial condition section of the code is designed for setting initial values of mass flow rate, pressure, enthalpy and fuel temperature at every node. The steady state routine for this is useful as a code in its own right since a whole range of steady states can be calculated in a single NAIAD run. After all geometry, flow model and fuel

model data have been read, sets of data specifying the initial condition calculations are read. Each set of data initiates one initial condition calculation in one flow path. The flow paths may be done in any order and a particular flow path can be done more than once to give a series of steady state results. Three parameters must be specified in the input data set for each flow path calculation: pressure and enthalpy at one node, and either pressure at another node or mass flow rate in the flow path. These parameters can be specified by value or by selecting a node in another flow path at which the parameter has the same value.

For example, a steady state could be set up in the network of Figure 1 as follows:

- Flow, pressure and enthalpy at a node in flow path 3 is read and a steady state is calculated for this flow path. This gives the pressures at connections 4 and 5 and the enthalpies of fluid entering and leaving the flow path.
- Assuming the flow in both flow paths 2 and 3 is from connection 4 to 5, then the enthalpies of fluid streams entering flow paths 2 and 3 are equal. The steady state of flow path 2 is then calculated from this enthalpy and the pressures at connections 4 and 5.
- . The remaining flow paths are done in a similar way.

If required, a series of initial condition calculations on one or a group of flow paths can be done and the results printed after each calculation.

2.3 Transient Calculations

The transient calculation begins from the initial conditions and proceeds time step by time step. The user has complete control of the calculation in that major changes in the calculation may be made between time steps, e.g. input data, boundary conditions and heating power may be changed and output written. Each time step is essentially a new calculation in that entries are made into the special routines which exercise this control at critical points in the calculation. Any coupling between nodes not routinely allowed for in the standard code can be incorporated in these special routines. If no changes are made, then the time step is done using the same boundary conditions, heating power and input data as the previous time step. The node couplings in the standard code are the hydraulic conservation equations, constant pressure boundary conditions, constant flow boundary conditions, and choke flow

boundary conditions. Special routines are required for any other coupling. For example, a calculation of both primary and secondary coolant loops is done in the test problem. In the special routines for this problem, the heat flux from one loop to the other through the steam generator is calculated at the start of each time step using the temperatures from the previous time step. Powers for nodes within the steam generator are updated each time step and the correct coupling is made.

Special routines can easily be written to do a series of transient calculations, each one using results from earlier calculations. It is possible, for example, to do a search for the break size giving the largest peak fuel temperature.

2.4 Equation of State

Thermodynamic equilibrium is assumed at each node and an equation of state routine for light water is supplied with the code. This routine interpolates from tables calculated with the ASTEM code [Moore 1971].

2.5 Slip Model

A slip model must be specified for each node so that momentum flux, static enthalpy, and kinetic energy density and flux can be calculated. The slip models in the code are:

- . Homogeneous flow
- . CISE [Di Francesco et al. 1971]
- . Beattie [Beattie 1974]
- . Jones [Jones & Dight 1962]
- . Smith [1970]
- . Fauske [1962]
- . Moody [1965]
- Small Bubble [Beattie 1975]
- . Modified CISE

Other slip models may easily be added, including ones which take into account the local distribution of voidage and velocity. The only limitation is that at any axial cross section the average mass, energy and kinetic energy densities and fluxes can be calculated from the values of mass flow rate, pressure and mean enthalpy at that axial cross section plus all geometry information. This excludes counter current flow and any effect of upstream flow structure.

Some of these slip models are well known; others are new [Beattie 1974, 1975] or their use in this way is new [Fauske 1962; Moody 1965].

The selection of the slip model is largely a matter of experience and evaluation of steady state data. These slip models should only be used in a calculation if they are physically possible in the range over which they are applied.

A necessary condition for this is that all three characteristics of the three coupled first order partial differential conservation equations be real over this range. The regions over which the above slip models have complex characteristics are given by Trimble & Turner [1976]. Of course with homogeneous flow there are no complex regions. This is also true for Beattie slip for flows less than choke flow. Jones, Smith and CISE all have complex characteristics in the low quality, low pressure, near choke flow region. Our experience is that calculations in which many points of the finite difference time-space mesh are in the regions of complex characteristics are numerically unstable. However in actual calculations with Jones, Smith and CISE slip we find this to be a rare occurrence.

2.6 Heat Transfer

A heat transfer package must be specified at each node. This package gives the surface heat transfer coefficient as a function of the state and flow of the coolant, the fuel surface temperature and geometry data. At present, only one package is available; it is designed for flow in round tubes and contains correlations for most flow situations including condensation and post dryout.

2.7 Friction

A friction relation must be specified at each node for the frictional pressure gradient calculation. At present, eight options are available; these are described in Section 6. Selection of the friction model should be based on steady state data.

A friction coefficient to cover pressure drops due to valves, bends, etc. is required for each segment of flow path between two nodes. Either an equivalent length or an equivalent fL/D may be specified.

2.8 Choke Flow

Choke flow relations of two types are available from the literature. In the first type, the choke flow is given as a function of the stagnation conditions; in the second, it is a function of the throat conditions. Both types may be applied in NAIAD by means of the special routines. However, as discussed in Section 1.4, it is desirable that choking be a consequence of the conservation equations and the slip

model, not an external condition imposed on them. This intrinsic choking is used in NAIAD. Intrinsic choking occurs when either of the two sonic characteristic velocities is zero [Turner & Trimble 1975]. However, characteristic velocities are not continuous functions of thermodynamic quality across the single-phase two-phase boundaries. Nevertheless, it is clear that choking will occur where the characteristic velocity changes sign. This is the criterion for choke flow used in NAIAD. A node can be located at the break and the geometric information and slip model specified for that node used to determine the choke flow rate.

3. METHOD OF CALCULATION

Each flow path in the problem is treated as one-dimensional and all quantities are evaluated at every node. All nodes are treated identically. The finite difference equations are derived in Section 4 from the conservation equations for the coolant and fuel. These equations, together with the boundary conditions and the connection equations, fully determine the system. Once all the coefficients are calculated at the start of a time step, the difference equations, boundary conditions and connection equations are solved for the values of the mass flow rate, pressure, enthalpy and fuel temperature at every node in the network at the end of the time step. The details of the solution method are given in Section 5. With mass flow rate, pressure, enthalpy and fuel temperature at every node determined, the equation of state, slip, friction and heat transfer routines are used to find all the coefficients for the next time step and so calculation proceeds.

As with any numerical calculation, some errors are to be expected. Because of the implicit formulation of the finite difference equations, one effect of these errors is that mass and energy are not exactly conserved. Thus errors arising from the numerical methods can be evaluated by checking the departure from conservation. As usual, these errors may be reduced by decreasing the time and space steps.

The computation time per time step per node depends on the problem being solved. This is because the number of iterations required to converge the network solution at each time step varies, the various slip heat transfer and friction options take different times, and the number of choked connections is important. However, on the AAEC's IBM360 model 65, central processor times of 10 to 20 milliseconds per time step per node have been experienced. The sample problem, which has 22 nodes and

l choked connection, took 44 and 100 seconds respectively to complete the coarse and fine time step calculations of the first 25 seconds.

The NAIAD code is written in the IBM360 language FORTRAN IV except for the matrix inversion routine and the AAEC input routine SCAN [Bennett & Pollard 1967] which are in 360 assembler language. The plot program uses many local features of the AAEC computing network and transfer of this program is not envisaged.

4. CONSERVATION EQUATIONS

The set of equations solved in the standard code are derived in this section.

4.1 One-dimensional Conservation Equations

Following Meyer [1960], we assume that the pressure P is constant over any cross-section of the flow path perpendicular to the z or flow direction, and write the conservation equations in terms of averages <> of local density d, velocity u and enthalpy h over such cross sections.

$$a \frac{\partial \rho}{\partial t} + \frac{\partial W}{\partial z} = 0$$

$$\frac{1}{a} \frac{\partial W}{\partial t} + \frac{1}{a} \frac{\partial}{\partial z} \left(\frac{M}{a} \right) + \frac{\partial P}{\partial z} = -F - \rho g s i n \theta$$

$$\frac{\partial}{\partial t} (a \rho H_s - a P + \frac{M}{2a}) + \frac{\partial}{\partial z} (W H + K) = pq - W g s i n \theta$$
where
$$\rho = \langle d \rangle,$$

$$W = a \langle ud \rangle,$$

$$M = a^2 \langle u^2 d \rangle,$$

$$\rho H_s = \langle h d \rangle,$$

$$W H = a \langle h u d \rangle, \text{ and}$$

$$K = \frac{a}{2} \langle u^3 d \rangle.$$

Many finite difference representations of these equations are possible. Here, the finite difference equations have been chosen to be numerically stable and to give good agreement with steady state results. Local stability analysis was used as a guide to the finite difference form. This indicated that evaluation of all space derivatives at the end of each time step was necessary. Actual calculations indicated that numerical instability could arise also if the friction and heat flux terms were evaluated at the start of each time step. Thus, implicit finite difference forms of these terms have been used in the equations

described below. Further discussion of the form of the finite difference equation is given in earlier work [Turner 1972; Turner & Trimble 1975].

We first define a slip-equation-of-state combination as the specification of ρ , M, H, and K as single valued functions of both the geometric parameters (a, equivalent diameter, roughness, etc.) at this or other z locations, and of W, P and H. An example of a slip-equation-of-state combination is Jones' slip [Jones & Dight 1962] together with the thermodynamic equilibrium equation of state for light water. We can now write

$$\frac{\partial \rho}{\partial t} = \frac{\partial \rho}{\partial w} \frac{\partial w}{\partial t} + \frac{\partial \rho}{\partial P} \frac{\partial P}{\partial t} + \frac{\partial \rho}{\partial H} \frac{\partial H}{\partial t} \equiv \frac{\partial \rho}{\partial x} \cdot \frac{\partial x}{\partial t}$$
$$x \equiv \begin{pmatrix} w \\ P \\ H \end{pmatrix},$$

and write similar expressions for the time derivatives of the other variables.

Consider a one-dimensional flow path for which the spatial finite difference mesh is the set of nodes m to n inclusive where

$$z_{m} < z_{m+1} < \dots < z_{n}$$

and the finite time step is Δt . The finite difference equations for each conservation equation are described in the following sections.

4.1.1 Mass equation

This may be written

$$a \frac{\partial \rho}{\partial x} \cdot \frac{\partial x}{\partial t} + \frac{\partial w}{\partial z} = 0$$

The finite difference equation is:

$$\left(\frac{a_{i} + a_{i+1}}{2\Delta t} \right) \left[\begin{cases} R \left(\frac{\partial \rho}{\partial x} \right)_{i} + (\rho_{i} - \rho_{i+1}) \frac{\partial R}{\partial x_{i}} \end{cases} . \quad (x_{i}' - x_{i}') \\ + \left\{ (1 - R) \left(\frac{\partial \rho}{\partial x} \right)_{i+1} + (\rho_{i} - \rho_{i+1}) \frac{\partial R}{\partial x_{i+1}} \right\} . \quad (x_{i+1}' - x_{i+1}') \right] + \frac{w_{i+1}' - w_{i}'}{z_{i+1} - z_{i}'} = 0$$

for i = m, m + 1, ..., n - 1.

The interpolation factor R is one half, and all its derivatives are zero unless a change of sign of the thermodynamic quality x occurs within the

segment. Its value and derivatives in this circumstance are discussed in Section 4.1.4. The truncation errors [Richmyer & Morton 1967] are

$$O(\Delta z^2) + O(\Delta t)$$
.

4.1.2 Momentum equation

The finite difference equation is

$$\frac{1}{2\Delta t} \left(\frac{1}{a_{i}} + \frac{1}{a_{i+1}} \right) \begin{bmatrix} R(W_{i}' - W_{i}) + (1-R) & (W_{i+1}' - W_{i+1}) \\ + \left\{ \frac{\partial R}{\partial x_{i}} \cdot (X_{i}' - X_{i}) + \frac{\partial R}{\partial x_{i+1}} \cdot (X_{i+1}' - X_{i+1}) \right\} & (W_{i} - W_{i+1}) \end{bmatrix}$$

$$+ \left(\frac{1}{a_{i}} + \frac{1}{a_{i+1}}\right) \frac{1}{2(z_{i+1} - z_{i})} \begin{bmatrix} \frac{1}{a_{i+1}} \left\{ M_{i+1} + \left(\frac{\partial M}{\partial x}\right)_{i+1} \cdot (x_{i+1}' - x_{i+1}) \right\} \\ -\frac{1}{a_{i}} \left\{ M_{i} + \left(\frac{\partial M}{\partial x}\right)_{i} \cdot (x_{i}' - x_{i}) \right\} \end{bmatrix}$$

$$+ \frac{P_{i+1}' - P_{i}'}{z_{i+1} - z_{i}} + F + [R\rho_{i} + (1-R)\rho_{i+1}] gsin\theta_{i,i+1} = 0$$

for i = m, m + 1, ..., n - 1.

The treatment of the momentum flux term is such that, for incompressible flow (where the term depends only on the end points), exact agreement with the Bernoulli equation is obtained. The truncation error for the momentum equation is

$$O(\Delta z^2) + O(\Delta t)$$
.

The friction term F is the sum of contributions from wall friction and pressure losses due to fittings etc.

$$F = \left(\frac{W'|G|v_{f}}{a}\right)_{i} \left[\frac{f_{i}}{D_{ei}} + \left(\frac{f_{i}K_{ei} + K_{fi}}{z_{i+1} - z_{i}}\right)\right] + \left(\frac{W'|G|v_{f}}{a}\right)_{i+1} \left[\frac{f_{i+1}}{D_{ei+1}} + \left(\frac{f_{i+1}K_{ei} + K_{fi}}{z_{i+1} - z_{i}}\right)\right],$$

where $v_f = \langle u \rangle /G$.

Only one of the two types of friction coefficients can be used in any one segment, so that either K_{ρ} or K_{f} is zero.

4.1.3 Energy equation

It is essential that the finite difference energy equations are able to handle the change in number of boundary conditions that occurs when the direction of flow at either end of the flow path changes. Such a finite difference scheme was formulated for single phase acoustics by Rose et al. [1967]. The way in which the present scheme overcomes this problem is described by Turner & Trimble [1975]; the scheme is given in full below.

First, the mass equation is used to eliminate the density derivative from the energy equation. Then

$$a\rho \frac{\partial H_{s}}{\partial x} \cdot \frac{\partial X}{\partial t} + (H - H_{s}) \frac{\partial W}{\partial z} - a \frac{\partial P}{\partial t} + \frac{1}{2a} \frac{\partial M}{\partial x} \cdot \frac{\partial X}{\partial t} + \frac{W\partial H}{\partial z} + \frac{\partial K}{\partial z} = pq - Wgsin\theta .$$

We evaluate space derivatives over the segment of flow path on the upstream side of the node, and all other terms at the node. The difference equation is

$$\frac{\mathbf{a_{i}}\rho_{i}}{\Delta t} \left(\frac{\partial \mathbf{H_{s}}}{\partial \mathbf{x}}\right)_{i} \cdot (\mathbf{x_{i}'} - \mathbf{x_{i}}) + (\mathbf{H_{i}} - \mathbf{H_{si}}) \left(\frac{\mathbf{w_{i}'} - \mathbf{w_{j}'}}{\mathbf{z_{i}} - \mathbf{z_{j}}}\right)$$

$$- \mathbf{a_{i}} \frac{(\mathbf{P_{i}'} - \mathbf{P_{i}})}{\Delta t} + \frac{1}{2\mathbf{a_{i}}\Delta t} \left(\frac{\partial \mathbf{M}}{\partial \mathbf{x}}\right)_{i} \cdot (\mathbf{x_{i}'} - \mathbf{x_{i}}) + \mathbf{w_{i}} \left(\frac{\mathbf{H_{i}'} - \mathbf{H_{j}'}}{\mathbf{z_{i}} - \mathbf{z_{j}}}\right)$$

$$+ \frac{1}{(\mathbf{z_{i}} - \mathbf{z_{j}})} \left[\mathbf{K_{i}} + \left(\frac{\partial \mathbf{K}}{\partial \mathbf{x}}\right)_{i} \cdot (\mathbf{x_{i}'} - \mathbf{x_{i}}) - \mathbf{K_{j}} - \left(\frac{\partial \mathbf{K}}{\partial \mathbf{x}}\right)_{j} \cdot (\mathbf{x_{j}'} - \mathbf{x_{j}})\right]$$

$$= \mathbf{P_{i}} \mathbf{q_{i}'} - \mathbf{W_{i}} \mathbf{gsin} \mathbf{\theta_{ij}}$$

where j = i - 1 if
$$W_i \ge 0$$
,
= i + 1 if $W_i < 0$, or i = m and $W_m = 0$,

and i takes all integer values such that i and j are in the range m to n. Thus the number of energy equations depends on the flow directions at the flow path ends. This is further discussed in Section 4.3. The truncation error is

$$O(\Delta z) + O(\Delta t)$$
.

4.1.4 Boiling boundary

Some of the derivatives of ρ , M, H and K with respect to W, P and H, are not continuous across the single-phase two-phase boundaries, the zero quality discontinuity being of much greater significance than that at unit quality. Segments which include this zero quality discontinuity are given special treatment. The position of the boiling boundary is determined by assuming that

$$y = H - H_{\ell} = y(X)$$

varies linearly over the segment. The density and flow rate are assumed to be constant over the liquid portion of the segment and to vary linearly with z over the two-phase portion (Figure 2). Then, for the segment between nodes i and i+l for which

$$y_{i+1} > 0$$
 and $y_i \le 0$,

$$\int_{z_{i}}^{z_{i+1}} \rho dz = [R\rho_{i} + (1-R)\rho_{i+1}](z_{i+1}-z_{i}) ,$$

where
$$R = 1 - \frac{y_{i+1}}{2(y_{i+1} - y_i)}$$
.

Similarly, if $y_{i+1} \le 0$ and $y_i > 0$

$$R = \frac{Y_{i}}{2(Y_{i} - Y_{i+1})}.$$

If no boiling boundary is present, $y_i y_{i+1} \ge 0$ and R = 1/2. In general,

$$R = R(X_{i}, X_{i+1}) \qquad .$$

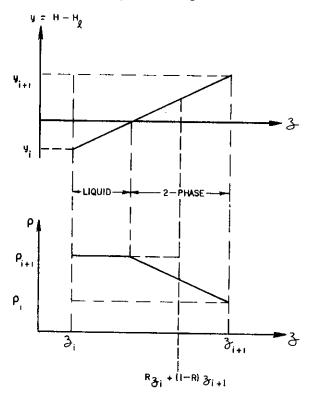
This defines the interpolation factor used in the mass and momentum equations.

The special treatment of the boiling boundary may be removed from the difference scheme either at the start of or during the calculation by including the statement

CALL GAMOFF

in the special routines MAIN or BCSETM. After execution of this statement, R is always one half and its derivatives are zero. It is also

FIGURE 2
Boiling Boundary



possible to include the variation of R at boiling boundary segments but set all the derivatives of R to zero. This is achieved by execution of the statement

CALL GAMDOF .

These options are included because we have, as yet, little experience with the special treatment of the boiling boundary. However, we recommend using the treatment unless problems arise.

4.2 Fuel

At each node, heat is assumed to be generated in a fuel node connected to the coolant by a surface heat transfer coefficient. If the heat capacity λ of this fuel node is zero,

$$q' = Q', T' = T'_{C}$$

and no fuel model calculation is required for this node as all generated power flows directly into the coolant. If λ is not zero, the finite difference equation (applicable to any geometry) is:

$$\frac{\lambda (T'-T)}{\Delta t} = Q'-q'$$

The surface heat flux is obtained from

$$q' = q + \frac{\partial q}{\partial x} \cdot (x' - x) + \frac{\partial q}{\partial T} (T' - T) ,$$
 where
$$q = h_T (T - T_C) = q(x, T) \text{ at each node,}$$
 and
$$h_T = \frac{h}{h + h_C} .$$

Two ways of obtaining the derivatives of q are available.

As $T_C = T_C(X)$, h = h(X,T) and h_F is constant,

$$\frac{\partial \mathbf{q}}{\partial \mathbf{x}} = (\mathbf{T} - \mathbf{T}_{\mathbf{C}}) \frac{\partial \mathbf{h}_{\mathbf{T}}}{\partial \mathbf{x}} - \mathbf{h}_{\mathbf{T}} \frac{\partial \mathbf{T}_{\mathbf{C}}}{\partial \mathbf{x}} ,$$

and
$$\frac{\partial \mathbf{q}}{\partial \mathbf{T}} = (\mathbf{T} - \mathbf{T}_{\mathbf{C}}) \frac{\partial \mathbf{h}_{\mathbf{T}}}{\partial \mathbf{T}} + \mathbf{h}_{\mathbf{T}}$$
.

The derivatives of $T_{\rm C}$ are carried in the equation-of-state tables. The derivatives of h are evaluated by numerical differentiation of the heat transfer coefficient correlation used to obtain h. This evaluation can be quite time consuming and is only necessary for stability when large time steps are used.

4.2.2 Explicit h

The parameter h is assumed constant in the determination of the derivatives of q. Then

$$\frac{\partial \mathbf{q}}{\partial \mathbf{x}} = -\mathbf{h}_{\mathbf{T}} \frac{\partial \mathbf{T}_{\mathbf{c}}}{\partial \mathbf{x}} \quad ,$$

and
$$\frac{\partial \mathbf{q}}{\partial \mathbf{T}} = \mathbf{h}_{\mathbf{T}}$$
.

This has no significant effect on accuracy as the difference equations already contain order Δt truncation errors. With this option the heat flux q* from the previous time step is used to obtain an estimate of the fuel surface temperature which is then used to obtain h.

$$T_{FS} = T - q*/h_{F}$$

$$h = h(X,T_{FS})$$
.

This procedure can be unstable for large time steps and hence it is not used with the previous option (Section 4.2.1). It can be suppressed by making $\mathbf{h}_{_{\mathbf{F}}}$ very large.

4.3 Connections

In general, several flow paths terminate at each connection. At some of these terminations the flow will be from connection to flow path; at others it will be from flow path to connection. The connection enthalpy η is defined as the enthalpy of fluid passing in the former direction. The second connection parameter ψ is called the connection control and may be pressure or mass flow rate.

Consider a flow path which begins at connection α node m and ends at connection β node n. The connections α and β are the flow path boundaries. The conditions within the flow path are calculated from the conservation equations described above and the flow path boundary conditions. The form of the boundary conditions depends on the sign of the flow at the flow path end and on the type of control at the connection. At the α connection, the boundary conditions are:

For
$$\alpha$$
 pressure controlled $P_m^{\, \text{!`}} = \psi_\alpha^{\, \text{!`}}$. For α flow controlled
$$W_m^{\, \text{!`}} = \psi_\alpha^{\, \text{!`}} \text{ .}$$
 Also
$$H_m^{\, \text{!`}} = \eta_\alpha^{\, \text{!`}} \text{ if } W_m > 0 \text{ .}$$

The number of boundary conditions depends on the direction of flow between the connection and the flow path - the enthalpy boundary condition is only used where flow is from the connection to the flow path.

Similarly, at the β connection

$$\begin{array}{lll} P_n' &=& \psi_\beta' & \text{ or } \\ W_n' &=& \psi_\beta' & \text{, and } \\ H_n' &=& \eta_\beta' \text{ if } W_n < 0 \end{array}.$$

Hence the number of boundary conditions may be two, three or four, depending on the flow directions. However, the total number of equations is always 4(n-m+1) as shown in Table 1.

TABLE 1
FLOW PATH EQUATIONS

Equation Type Number of Equations			ions
Mass		n-m	
Momentum	n-m		
		$w_n \ge 0$	w _n < 0
Energy	M > 0	n-m	n-m-1
	w _m ≤ 0	n-m+1	n-m
Fuel	n-m+1		
ψ boundary	2		
		$ w_n \ge 0$	w _n < 0
Enthalpy boundary	w _m > 0	1	2
	w _m ≤ 0	0	1
Total	4(n-m+1)		

Consider now a network containing n nodes and b connections. The 4n flow path equations have been described above. These contain up to 2b connection parameters, hence an equal number of connection equations are required. At internal connections, two equations are always required, while at other types either one or two equations are required depending on flow direction. The connection equations for each connection type are described below. Input parameters specify which connections are of each type.

4.3.1 Internal

An internal connection is pressure controlled and must be common to more than one flow path. Each of these flow paths has a node at the common connection and the ψ boundary conditions are such that the pressures at these nodes are all equal to ψ . The pressure losses at such junctions can be included in friction coefficients or area changes in the last segment of the flow paths joining the connection so it is unnecessary to include them here. Internal connections are assumed to have zero volume, hence there is at least one flow towards and one away from the connection.

Let node i be a node for which flow at the start of the time step is towards the connection and let node j be one for which flow is away from it. In this section, the convention is adopted that all flows out of the connection are positive; thus

$$W_{i} \leq 0$$
 and $W_{j} > 0$.

As connections have zero volume, the total flow is zero

$$\Sigma W' = 0$$

It is assumed that kinetic energy can be neglected and that the stagnation enthalpies of all fluid streams leaving the connection are equal. Therefore

$$\eta'_{\alpha} = \sum_{in} w'H'/\sum_{in} w'$$
.

These are the two equations for an internal connection.

4.3.2 Explicit

The connection equations in the standard code are:

$$\begin{pmatrix} \psi'_{\alpha} \\ \eta'_{\alpha} \end{pmatrix}$$
 = values at the start of the first time step,

where ψ is flow or pressure depending on whether connection α is flow or pressure controlled.

More complex equations giving $\psi_{\alpha}^{"}$ and $\eta_{\alpha}^{"}$ in terms of quantities at the start of each time step, (flows, pressures, qualities, void fractions, etc.) may be coded in the special routine BCSETM which is called at every time step. An example is provided in the test problem.

4.3.3. Implicit

The connection equations must be coded in the special routine BPWIMP. These connection equations are solved by an iterative method discussed in Section 5. The special routine BPWIMP(KC) is called once at each iteration for each connection KC declared implicit in the input data. Any boundary conditions may be applied to explicit and implicit connections. The difference is that in BPWIMP, the enthalpies, pressures and flows for all connections at the end of the time step are available, as well as all parameters at the start of the time step. Thus implicit connection equations can be used, e.g. in the test problem, a pump model is set up using one implicit and one explicit connection.

4.3.4 Choked

The conditions for choke flow or zero sonic characteristic velocity are [Turner & Trimble 1975]

$$\gamma \equiv \left[\frac{\partial}{\partial P} \left(P + \frac{M}{a^2} \right) \right]_{\star} = 0$$

To determine γ , consider

$$WH_O = WH + K$$

where ${\rm H}_{\rm O}$ is a stagnation enthalpy. This is a function of the components of X, i.e. W, P and H, hence

$$= \frac{\partial \mathbf{P}}{\partial \mathbf{A}} + (\mathbf{M} + \frac{\partial \mathbf{H}}{\partial \mathbf{K}}) \left(\frac{\partial \mathbf{B}}{\partial \mathbf{H}}\right)^*$$

which

= 0 as * means constant G, H_{O} and a.

Therefore,
$$\left(\frac{\partial H}{\partial P}\right)_{\star} = -\frac{\partial K/\partial P}{W + \partial K/\partial H}$$
.

 $\mathfrak{g}^{(0)}$ Similarly, M is a function of X hence

$$\gamma = 1 + \frac{1}{a^2} \frac{\partial M}{\partial X} \cdot \left(\frac{\partial X}{\partial P}\right)_{\star}$$

$$= 1 + \frac{1}{a^2} \left[\frac{\partial M}{\partial P} + \frac{\partial M}{\partial H} \left(\frac{\partial H}{\partial P} \right)_{\star} \right]$$

$$= 1 + \frac{1}{a^2} \left[\frac{\partial M}{\partial P} - \frac{\partial M}{\partial H} - \frac{\partial K/\partial P}{(W + \partial K/\partial H)} \right] .$$

The value of γ' at the end of the time step is evaluated as a function of the pressure ψ'_{α} at the choke connection. The choking pressure P'_{c} at the connection is the pressure at which γ' changes sign.

Now choking only occurs for throat pressures above the receiver pressure, hence the connection equation is

$$\psi_{\alpha}^{\prime} = \text{Max} (P_{r}^{\prime}, P_{c}^{\prime})$$
 ,

where the receiver pressure P_r^{\prime} is linearly reduced to the input pressure PAT over the break opening time TOPEN. These two quantities are the input data for a choked connection.

5. METHOD OF SOLUTION

The solution of the set of equations described in Section 4 is done in four stages.

5.1 Elimination of Surface Heat Flux and Fuel Temperature

The surface heat flux q' appears in the energy, the fuel and the surface heat flux equations, while T occurs in only the last two of these three equations. T is eliminated from these to give

$$\mathbf{q'} = \frac{\mathbf{q} + \frac{\partial \mathbf{q}}{\partial \mathbf{x}} \cdot (\mathbf{x'} - \mathbf{x}) + \frac{\partial \mathbf{q}}{\partial \mathbf{T}} Q' \frac{\Delta \mathbf{t}}{\lambda}}{1 + \frac{\partial \mathbf{q}}{\partial \mathbf{T}} \frac{\Delta \mathbf{t}}{\lambda}}$$

which is used to eliminate q' from the energy equation.

5.2 Solution of Flow Path Equations

Consider a flow path which includes nodes m to n and joins connections α and β . There are 4(n-m+1) equations; elimination of the fuel temperatures reduces this to 3(n-m+1). In this stage of the network solution, these equations are solved for all W', P', and H', which are then determined as linear functions of the four connection parameters. The first step is to write the 3(n-m+1) equations in groups of three in the following way:

Group m

 $\begin{array}{lll} \psi_\alpha^{\:\raisebox{3.5pt}{\text{\circ}}} & \text{boundary condition} \\ \text{Node m, m+1} & \text{momentum equation} \\ \eta_\alpha^{\:\raisebox{3.5pt}{\text{\circ}}} & \text{boundary condition if } W_{\atop{m}} > 0 \\ \text{Node m} & \text{energy equation if } W_{\atop{m}} \leqslant 0 \end{array}$

Only unknowns at nodes m and m+l appear.

Group j = m+1 to n-1

Node j-1, j mass equation

Node j, j+l momentum equation

Node j energy equation

Only unknowns at nodes j-1, j and j+1 appear.

Group n

Node n-1, n mass equation

 ψ_o^* boundary condition

 η_{R}^{r} boundary condition if $W_{n} < 0$

Node n energy equation if $W_n \ge 0$

Only unknowns at nodes n-1 and n appear.

This set is now written as the matrix equation:

where A, B, C are 3 by 3 matrices and the d vectors have 3(n-m+1) components. The A, B and C matrices are removed by forward elimination and back substitution to yield a solution of the form

$$x'_{j} = y_{j}^{(0)} + \psi_{\alpha}' y_{j}^{(1)} + \eta_{\alpha}' y_{j}^{(2)} + \psi_{\beta}' y_{j}^{(3)} + \eta_{\beta}' y_{j}^{(4)}$$

$$j = m, m+1, ..., n ; ...(1)$$

where the Y vectors have three components.

This procedure requires that all the B matrices be non-singular. The determinant of a B matrix is a cubic polynomial in $1/\Delta t$. We are not able to make any general statement about the roots of this polynomial, except that in some situations a positive root exists. This is only a

problem if W, P and H are nearly equal at several equally spaced, geometrically identical nodes, in which case the B matrices for these nodes will all be singular for almost the same values of Δt . In practice, the problem is easily overcome by using unequal node spacings, thus ensuring that, at most, only one near singular B matrix is present.

5.3 Connection Balance

Next, the connection equations are solved for the connection parameters

$$(\psi_{\alpha}^{\dagger},\eta_{\alpha}^{\dagger})$$
 for $\alpha=1,\ldots,b$

In the simplest case, the connection equations are those given in Section 4.3.2, and no solution procedure is needed as the values from the last time step are correct. For more complex explicit connections, solution is a simple matter as the connection equations are in terms of quantities at the start of the time step. These connection parameters are calculated in BCSETM at the start of the time step.

The remaining connections are solved by an iterative procedure beginning with the values from the last time step. Each internal, choked and implicit connection is considered in turn. For internal and choked connections, the equations in which ψ_{α}' or η_{α}' occur are the matrix equations for flow paths connecting to connection α and the connection equations for connection α . At each iteration these equations are solved for ψ_{α}' and η_{α}' using the most recent values of ψ_{β}' , η_{β}' for $\beta \neq \alpha$. For each implicit connection the statement

CALL BPWIMP (KC, N, ω)

is executed, where KC is the connection number, N the iteration number and ω the relaxation parameter. Thus, the solution procedure for implicit connection equations must be coded into the special routine BPWIMP.

In general, ψ'_{α} and η'_{α} for an implicit connection will not only appear in the matrix equations for flow paths connecting to α and the connection equations for α , but also in the connection equations for other connections. For example, in Figure 1 the connection parameters for connection 5 would appear in the connection equations for connections 5, 6 and 7. In this case, one of the group of connections is designated implicit in the input data and the rest explicit, and all the connection equations for the group are solved together when the special routine BPWIMP is called with KC equal to the number of the implicit connection.

This is essentially a Gauss-Siedel iteration procedure and can be expected to converge provided the connection groups are sufficiently independent. The degree of independence will determine the convergence, or sometimes divergence of the iteration scheme. If the latter occurs, it is assumed that the connections are made more independent by reducing the time step. Either a time step is found which gives convergence, or the final estimate of the connection parameters for the minimum time step is accepted and an error message printed. We have found the convergence to be almost linear, i.e. if

$$\psi(n) = \begin{bmatrix} \psi_1' \\ \eta_1' \\ \vdots \\ \vdots \\ \psi_b' \\ \eta_b' \end{bmatrix}$$
stimate of the con

is the estimate of the connection parameters after n iterations and

$$\Delta \Psi^{(n)} = \Psi^{(n)} - \Psi^{(n-1)} ,$$

then for n sufficiently large,

$$|\Delta \Psi^{(n)}| = e |\Delta \Psi^{(n-1)}|$$

where e < 1, constant and independent of n.

Hence, an acceleration scheme is used which corresponds to successive line over-relaxation in the iterative solution of linear systems. The relaxation is applied to the calculated change in the parameters. If at any iteration the change in ψ_{α} is $\Delta\psi_{\alpha}$, the relaxed ψ_{α} is $\psi_{\alpha}+\omega\Delta\psi_{\alpha}.$ As highly non-linear conditions may be present at implicit connections, this acceleration scheme is only applied to internal and choke connections. The relaxation constant ω is available in the special routine BPWIMP for use at implicit connections and should be used unless convergence problems arise.

The network convergence criterion is that

$$\sum_{\alpha=1}^{b} \left[\left(\frac{\Delta \psi_{\alpha}^{i}}{\psi_{\alpha}^{i}} \right)^{2} + \left(\frac{\Delta \eta_{\alpha}^{i}}{\eta_{\alpha}^{i}} \right)^{2} \right] < \text{TOL}^{2}$$

Both TOL and the maximum number of network interations are input parameters. Values of 10^{-5} to 10^{-10} have been used for TOL and up to 50 iterations allowed. The error in the internal connection balances can easily be checked from the output. If significant errors are found, the value of TOL should be reduced.

In general, any convergence problems will be due to too close a coupling between connection groups. This can be overcome by keeping connections well separated. If this is not feasible, the closely coupled connections can be placed in a group and solved together in the special routine BPWIMP.

5.4 Back Substitution

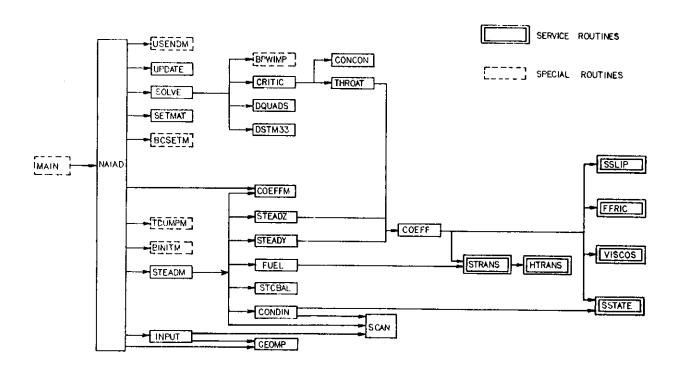
Once all the connection parameters are knowm, the values of W', P' and H' at all nodes are determined from Equation (1). Then the fuel temperatures are determined from the fuel equation (Section 4.1.2).

6. DESCRIPTION OF ROUTINES

6.1 Introduction

The general structure of the code is shown in Figure 3. The routines which make up the code are either service routines, special routines or ordinary routines.

FIGURE 3
Structure of NAIAD



The service routines are those which supply data on the coolant in the flow network and on the heat transfer. The standard service routines are for light water. They may need to be changed or extended if problems outside the range of these standard routines are to be done. For example, all service routines would need to be replaced for NAIAD to be used for sodium-cooled networks. The standard service routines and the requirements for new service routines are described in Section 6.2.

The special routines are intended to provide a simple way to expand the standard code in order to:

- compute and write additional output either for printing or plotting;
- impose time dependent boundary conditions;
- include any coupling between connections or nodes not routinely allowed for in the standard code, e.g. as required to simulate pumps, steam generators, pressure controllers, etc; and
- . include NAIAD runs as part of a higher calculational procedure, e.g. a search for maximum fuel temperature as a function of break size.

The sample problem includes examples of all special routines except MAIN. The special routines supplied with NAIAD are given in Section 6.3 together with requirements for new special routines.

The remaining routines are briefly outlined in Section 6.4.

6.2 Service Routines

6.2.1 SSTATE

This is the equation of state routine. The assumption of thermodynamic equilibrium enables the coolant state to be determined from the pressure and static enthalpy which are the input data. The output is the remaining items in the call statement argument list which is given below:

$$\begin{array}{l} P,S,\rho,x_{_{S}},\left(\frac{\partial\rho}{\partial H_{_{S}}}\right)_{p},\left(\frac{\partial\rho}{\partial P}\right)_{H_{_{S}}},T_{_{C}},T_{_{S}},\alpha,H_{_{\ell}},H_{_{V}},\\ \\ \rho_{\ell},\rho_{_{V}},S_{\ell},S_{_{V}},\left(\frac{\partial T_{_{C}}}{\partial P}\right)_{H_{_{S}}},\left(\frac{\partial T_{_{C}}}{\partial H_{_{S}}}\right)_{p},H_{_{S}},\left(\frac{d\rho_{_{V}}}{dP}\right)_{q},\\ \\ \frac{d\rho_{\ell}}{dP},\left(\frac{dH_{_{V}}}{dP}\right)_{q},\left(\frac{dH_{_{\ell}}}{dP}\right)_{q},\left(\frac{dT_{_{C}}}{dP}\right)_{q},\\ \end{array}$$

where \mathbf{x}_s is static thermodynamic quality $\frac{H_s^{-H} \ell}{H_v^{-H} \ell}$. In the two-phase

region, all output quantities except S, \mathbf{x}_{S} , α , ρ and its derivatives, are functions of P only. In the single phase region, the saturation line values are those at pressure P.

The standard SSTATE is for light water over the range 0.045 to 16 MPa and 273 to 1073 K. The required quantities are obtained by interpolation in tables read from disk at the first entry to SSTATE. The tables have been generated from the state equations for water [Meyer 1967] by means of the ASTEM code [Moore 1971]. The pressure/temperature mesh used is given in Table 2.

TABLE 2
EQUATION OF STATE MESH

Pressures (MPa)							
0.045 to 0.09	at intervals of 0	.003					
0.090 to 0.200	at intervals of 0	.005					
0.20 to 0.7	at intervals of 0	.01					
0.7 to 1.50	at intervals of 0	.02					
1.50 to 4.00	at intervals of 0	.05					
4.00 to 10.0	at intervals of 0	.1					
10.0 to 16.0	at intervals of 0	. 2					
Temperatures (K)							
At each pressure point, 5 liquid and							
6 vapour points.							
Liquid: T	Vapour: T						
Tsat-10) ^T sat	+10					
Tsat-20) ^T sat	+20					
Tsat-50) ^T sat	+50					
273.15	T sat	+100					
	1073	.15					

6.2.2 Viscos

The liquid and vapour viscosities μ_f and μ_g are calculated in this routine. Two formulae of Bruges & Gibson [1969] are used. The first (Equation 2, United Kingdom Committee on the Properties of Steam 1970) specifies the viscosity of saturated liquid light water as a function of

temperature in the range 375 to 623 K. The viscosity of both saturated and compressed water are calculated from this formula. For compressed water, the values are accurate to within a few per cent over the above temperature range and up to 16 MPa.

The second formula (Equation 4, United Kingdom Committee on the Properties of Steam 1970) specifies the viscosity μ_g of superheated steam, including saturated steam, as a function of the vapour density ρ_g and the temperature $T_{_{\rm C}}$ in the ranges 375 to 600 K and 0 to 16.5 MPa.

The call statement argument list is

$$T_{c}, \rho_{q}$$
 (input data)

and μ_{f}, μ_{g} (output).

6.2.3 FFRIC

This is the friction routine; in it, the Fanning friction factor is evaluated. The input data defining the state of the coolant and the node i at which friction is to be evaluated is in the call statement argument list which is

i, w, x, 1-x,
$$\rho_{\ell}$$
, ρ_{v} , μ_{f} , μ_{g} , μ_{h} , T_{s} , f,

where here x equals physical quality which is always in the range 0 to 1, and μ_h is the viscosity used in the homogeneous friction model (see below).

The input data defining the geometry at node i (roughness, equivalent diameter, flow area, etc.) and the required friction model are taken from BLANK common. This division of input must be adhered to when new friction models are added because sometimes W \neq W and $\rho_{\rm v} \neq \rho_{\rm vi}$ etc.

The standard routine contains four basic models and an optional boiling wall correlation to allow for regime changes at high quality, making eight friction options in all. In each of the basic models, the viscosity used to obtain the Reynolds number is calculated in a different way:

(a) Homogeneous

The viscosity used here is calculated in COEFF.

$$\frac{1}{\mu} = \frac{1}{\mu_h} = \frac{\mathbf{x}}{\mu_g} + \frac{1-\mathbf{x}}{\mu_f} \quad .$$

(b) Bubbly Flow [Beattie 1973]

$$\mu = \mu_f \left[1 + \frac{x}{v_f \rho_v} + \frac{1.5 \mu_v}{\mu_q + \mu_f}\right]$$
,

where $v_f = \frac{x}{\rho_v} + \frac{1-x}{\rho_{\ell}}$ is the flow specific volume.

(c) <u>Wavy gas liquid interface</u> [Beattie 1973] e.g. annular flow

$$\mu = \frac{1}{v_f} \left[\frac{1-x}{\rho_\ell} \quad \mu_f + \frac{x}{\rho_v} \mu_g \right] \quad .$$

(d) Small bubble [Beattie 1973]

$$\mu = \mu_{\ell} \left(1 + \frac{2.5 \text{ x}}{v_f \rho_v} \right) .$$

In all cases, the Reynolds number R_e is $R_e=\frac{D_e\;G}{\mu}$, and the Fanning friction factor f is the solution of one of the following equations depending on the value of R_a .

<u>Turbulent Flow</u> $R_{e} \ge 4000$

$$\frac{1}{\sqrt{f}} = -4 \log_{10} \left[\frac{\varepsilon}{3.7D_{e}} + \frac{1.255}{R_{e}} \sqrt{f} \right]$$

Interpolation Region $4000 > R_e > 2000$

$$f = \frac{16}{R_e} + (f_* - 0.004) \left[\frac{R_e}{2000} - 1 \right]$$
.

where f_* is the solution of the turbulent flow equation.

<u>Laminar Flow</u> $R_e \le 2000$

$$f = \frac{16}{R_{e}} .$$

Near Zero Flow $R_e \le 0.0016$

$$f = 10000$$

Often, a Reynolds number of 2000 is taken as the boundary between laminar and turbulent flow. The interpolation region removes the discontinuity that such a procedure entails and gives reasonable agreement with the pressure drop measurements of Nikuradse [1933].

The optional boiling wall friction factor [Beattie 1973] is the solution of

$$\frac{1}{\sqrt{f}}$$
 = 14 log₁₀ (f W_e) -10.4 ,

where the Weber number W_e is $\frac{G^2D_ev_f}{\sigma}$ and σ is the surface tension. Rather than solve this equation at every node and time step, the solution has been evaluated for Weber numbers in the range 10^2 to 10^{22} and fitted as a quadratic polynomial in log W_e . The liquid surface tension is calculated from a straight line fit to the saturation line surface tension as a function of saturation temperature. If the boiling wall extension is specified, then f is the lower of the friction factors from the boiling wall extension and the specified basic model.

6.2.4 SSLIP

This is the slip routine. The call statement argument list is

i,
$$\rho_{\ell}$$
, ρ_{v} , H_{ℓ} , H_{v} , W , P , H , T_{s} , μ_{f}

which are input data, and

$$\rho$$
, κ , M , H_s , α , k , $< u_{\ell} >$

which are the output. The parameter k is the slip ratio $< u_V^{>/<} u_Q^{>}$. These input data define the state of the coolant. Input data defining the slip model and the geometry at node i are taken from BLANK common. This division of input must be adhered to when new slip models are added because sometimes $P \neq P_i$, $H_V \neq H_{Vi}$, etc. The slip routine is only called for two-phase non-homogeneous flow. Single phase and homogeneous flow are done in COEFF. In all cases, the quality is

$$x = \frac{H-H_{\ell}}{H_{\nu}-H_{\ell}}$$

and the flow void fraction is

$$\beta = \frac{x \rho_{\ell}}{x \rho_{\ell} + (1-x) \rho_{V}}.$$

The standard slip routine contains eight slip models. In seven of these, separated flow is assumed *i.e.* all vapour moves at velocity $\mathbf{u}_{\mathbf{V}}$, all liquid at velocity \mathbf{u}_{ℓ} . For these either the slip ratio k or the void fraction α is calculated.

(a) <u>Jones</u> [Jones & Dight 1962]

The void fraction α is the solution of

$$\frac{\alpha}{\beta} = E + (1-E)\alpha^{T} ,$$
where $r = 3.33 + 0.577 \frac{P}{P_{C}} + 4.735 \left(\frac{P}{P_{C}}\right)^{2} ,$
 $E = 0.71 + 0.29 \frac{P}{P_{C}} ,$ and
 $P_{C} = \text{critical pressure (22.12 MPa)}.$

(b) CISE [Di Francesco et al. 1971]

where
$$V = \frac{\beta}{1-\beta} ,$$

$$E = 1.578 \left(\frac{\mu_{\ell}}{GD_{e}}\right)^{0.19} \left(\frac{\rho_{\ell}}{\rho_{v}}\right)^{0.22} ,$$

$$C = 0.0273 \frac{G^{2}D_{e}}{\sigma_{\varrho}\rho_{\ell}} \left(\frac{\mu_{\ell}}{GD_{e}}\right)^{0.51} \left(\frac{\rho_{v}}{\rho_{\varrho}}\right)^{0.08} , \text{ and }$$

 σ_{ℓ} = liquid surface tension.

(c) Beattie

The Beattie slip model is defined as follows. Let

$$c = (3.125f)^{-\frac{1}{2}}$$
,
 $w = \ln[2(e^{C}-1)] - \frac{1}{2}$, and
 $s = 1+2 \exp[-w^{-\frac{1}{2}}]$.

Then the void fraction is the solution of

$$(1-\beta) = (1-\alpha) \left[w + \frac{1}{2} + \ln \left(\frac{s-\alpha}{2} \right) \right]$$

In the original form of the model [Beattie 1974], s = 1 and w = c, which gave small errors at x = 0 and 1. The forms of s and w given above eliminate these errors.

(d) Fauske [1962] $k = (\rho_0/\rho_y)^{\frac{1}{2}}$

(e) <u>Moody</u> [1965]

$$k = (\rho_{\ell}/\rho_{v})^{1/3} .$$

(f) Modified CISE

The modified CISE correlation is the same as the CISE correlation except that the dimensionless parameter c is replaced by

$$\frac{2E}{\pi}$$
 arctan $\frac{\pi c}{2E}$

where $E=\frac{\sqrt{1+4V}-1}{2V}$ and V and c are defined in Section 6.2.4(b) for the CISE model. This modification was made because of the non-hyperbolic character of the conservation equations with CISE slip in conditions of high quality and mass velocity [Trimble & Turner 1976]. The modification has greatly reduced this non-hyperbolic region and, from a preliminary assessment, not altered the agreement of the correlation with measured density data.

(g) Smith [1970]

$$k = E + (1-E) \sqrt{\frac{x \rho_{\ell} + E(1-x) \rho_{v}}{x \rho_{v} + E(1-x) \rho_{v}}},$$

where E = 0.4

Thus either algorithms for α or k have been given for each of the seven separated flow slip models. If α has been calculated then

$$k = \frac{x(1-\alpha)\rho_{\ell}}{(1-x)\alpha\rho_{\chi}}$$

If k has been calculated then

$$\alpha = \frac{x \rho_{\mathbf{v}}^{k}}{x \rho_{\mathbf{v}}^{k+(1-x)} \rho_{\ell}} .$$

Then the velocities are given by

$$u_v = \frac{Gx}{\alpha \rho_v}$$
, $u_{\ell} = \frac{G(1-x)}{(1-\alpha)\rho_{\ell}}$,

and
$$\rho = \alpha \rho_{V} + (1-\alpha) \rho_{\ell} ,$$

$$H_{S} = [\alpha H_{V} \rho_{V} + (1-\alpha) H_{\ell} \rho_{\ell}]/\rho ,$$

$$K = \frac{a}{2} \left[\alpha \rho_{V} u_{V}^{3} + (1-\alpha) \rho_{\ell} u_{\ell}^{3} \right] ,$$

$$M = a^{2} \left[\alpha \rho_{V} u_{V}^{2} + (1-\alpha) \rho_{\ell} u_{\ell}^{2} \right] .$$

Separated flow is not assumed in the Small Bubble slip model hence none of these relations apply.

(h) Small Bubble [Beattie 1975]

 $C_0 = 1 + 2.65\sqrt{f}$

The Beattie small bubble correlation should only be used for void fractions below about 0.6. It is defined as follows:

$$\alpha = \frac{\mathbf{x}\rho_{\ell}}{\left[\mathbf{x}\rho_{\ell} + (\mathbf{1} - \mathbf{x})\rho_{\mathbf{v}}\right]C_{0}},$$

$$K = \frac{G^{3}}{2\rho_{\ell}^{2}} \left[1 + \mathbf{x}\left(\frac{\rho_{\ell}}{\rho_{\mathbf{v}}} - 1\right)\right]^{2} \left[1 - (C_{0}^{2} - 1)\mathbf{x}\left(\frac{\rho_{\ell}}{\rho_{\mathbf{v}}} - 1\right)\right],$$

$$M = \frac{G^{2}}{\rho_{\ell}} \left[1 + \mathbf{x}\left(\frac{\rho_{\ell}}{\rho_{\mathbf{v}}} - 1\right)\right] \left[1 - (C_{0} - 1)\mathbf{x}\left(\frac{\rho_{\ell}}{\rho_{\mathbf{v}}} - 1\right)\right].$$

The density ρ and static enthalpy H are given by the separated flow equations. This slip model should be used with the small bubble friction option so that the evaluation of the Fanning friction factor f is consistent with the slip model.

6.2.5 HTRANS and STRANS

The surface heat transfer coefficient is calculated in HTRANS. The call statement argument list is h and THT which are the output data, and

$$D_e$$
, p, G, P, T_c , T_s , T_{FS} , ρ_l , ρ_v , x, H_l , and H_v

which are the input data. In general, a number of correlations are needed to cover all situations. Such a group of correlations together with the rules determining which is to be used, is a heat transfer package. The indicator THT is available to record which correlations are used in determining the heat transfer coefficient. The package is used to evaluate the heat transfer coefficient h and its derivatives by five entries for slightly different values of W, P, H and T_{FS} . As the package rules only need to be evaluated once, the input value of THT determines whether these rules are to be used or not. If it is positive, then the rules are used. If it is negative, the correlation used on the last entry to HTRANS is used again.

The standard routine contains only one package which is designed for round tubes. Steady state heat transfer correlations expressed as a function of local conditions (coolant temperature, pressure, quality and pipe surface temperature) were assumed applicable to transient situations. This is the local condition hypothesis of Hassid & Rychlicki [1971] which predicts dryout somewhat sooner than it occurred in experiments.

Heat transfer correlations included to cover all situations are summarised in Table 3. Condensing heat transfer correlations are needed when steam or two-phase coolant heated in a heated pipe section moves to sections in which the pipework is considerably cooler. They may be required also if the pipe model is used to simulate a heat exchanger. Akers et al. [1959] gives a Reynolds number of 50000 for the boundary between their two condensing heat transfer correlations. However, the correlations actually intersect at 76228 and this value is used here. Also the correlations of Akers et al. only cover the two-phase region. McAdams [1954 p.351] recommended that heat flux is almost independent of superheating of the vapour; this is followed in the superheat region. In the boiling non-dryout region, the heat transfer coefficient corresponding to three correlations is calculated and the largest taken. This ensures continuity of surface heat transfer rate in this region as a function of coolant conditions and fuel surface temperature.

At each heat transfer coefficient calculation, pipe surface temperature and coolant state are known from the previous time step. The heat transfer coefficient is evaluated assuming no dryout then, if appropriate, tested to determine if the dryout limit is exceeded. If this is so, the appropriate dryout heat transfer coefficient is calculated. In the steady state calculations, heat flux is known and surface temperature must be determined. Here an iterative procedure is used. In those steady state calculations where dryout and non-dryout solutions exist, the latter solution is used.

An indicator THT is available to show which heat transfer correlations are used in evaluating the heat transfer coefficient:

$$THT = \ell + m + n ,$$

where ℓ and m indicate the correlation used to evaluate the coefficient assuming no dryout, m indicates the dryout correlation used to determine if the dryout limit is exceeded, and n the post-dryout correlation used.

In some situations, a dryout correlation is not appropriate, e.g. during condensation; zero values of m and n indicate these. Possible values of ℓ , m and n and their meanings are given in Table 3. For example, THT = 313 would indicate that:

- the heat transfer coefficient, assuming no dryout, was evaluated using the nucleate boiling correlation;
- . the dryout limit on the heat transfer coefficient was determined from Bernath's correlation;
- . dryout was found to be present; and
- Beattie's post-dryout correlation was used to evaluate the actual heat transfer coefficient used.

HTRANS is coded in Btu, foot, second units and conversion to SI units is carried out in STRANS.

 $\frac{\text{TABLE 3}}{\text{CORRELATIONS FOR SURFACE HEAT TRANSFER COEFFICIENT}}$ Heat Transfer Indicator THT = ℓ + m + n

Surface temperature (T _c)	Quality x			
s	x ≤ 0	1 > x > 0		x ≥ 1
Ts < Tsat	Liquid Dittus & Boelter [1930] \$\ell = 1, m = 0, n = 0\$	Akers et al. [1959] m = 0, n = 0	ndensir R < 7 Re > 7	6228 £ = 3
	Bo Non-dryout	oiling		Superheat
	Largest of - subcooled lique Dittus & Boelter [1930] - nucleate boil: Tong [1965 p.118] - forced convect Shrock & Grossman [1959] from Wright [1961]	ing zion boiling	£ = 5 £ = 3 £ = 2	
T ≥ T sat	$\frac{\text{Dryout}}{x \le 0, \text{ Bernath } \sim \text{Tong [1965]}}$ $x > \rho_{V}/(\rho_{V} + \rho_{T}), \text{ Biasi et al}$		m = 10	### ##################################
	low quality high quality		m = 30 m = 40	n = 0
	Other x, linear interpolati Bernath & Biasi low quality	on on x between	m = 20	
	Post Dryout - Dryout limit	not exceeded	n = 0	
	$x \le 0.05$ - Tong [1965 p.129] low speed flow high speed flow high speed flow $\frac{h s D_e}{k} = 0.0192 \left(\frac{GD_e}{\mu_V}\right)$	v V	n =100 n =200	
	evaluated at t Beattie [1972]	emperature T	n =300	

6.3 Special Routines

6.3.1 General

All quantities likely to be used within the special routines are in the various common blocks described in Tables 4 to 11. The required common blocks must be included in the special routines. The labelled commons BCONDS and POWER\$ are of particular importance in that they contain the connection parameters ψ and η in the arrays BPW and BHO, and the flow path powers in PGENP, these quantities being those most likely to be required. Any service routine may be called by any special routine. All arithmetic and REAL arrays are double precision hence an

IMPLICIT REAL*8 (A-H, O-Z)

card must be included in each subroutine. All the standard special routines consist of the three cards

SUBROUTINE name

RETURN

END

except for MAIN which is

CALL NAIAD

STOP

END

Thus, an actual NAIAD run is initiated in the special routine MAIN by this CALL statement. Each standard special routine except MAIN is automatically inserted in composite codes if no other routine of that name is provided. The standard MAIN is included by using the CINSERT card described in Section 9.1 and demonstrated in the sample problem.

6.3.2 BINITM

This routine is called once by NAIAD just before the transient calculation commences (i.e. after INPUT and STEADM). Any initial conditions (which are W, P, H and T at all nodes) or input data for the transient calculation may be modified here. The effects of these modifications appear in the time zero and subsequent output. Also any time independent constants in the connection equations should be calculated here.

6.3.3 TDUMPM

This routine is called each time step immediately after the state of the flow network has been completely evaluated and printed. Instructions for any additional output required should be placed in TDUMPM. The call statement argument list contains one variable, IPRINT which indicates

(continued p.42)

TABLE 4
BLANK COMMON

Variable name and dimension	Description	Algebraic symbol
Z(50)	Displacement along pipe	Z
A(50)	Pipe area	a
DE(50)	Pipe equivalent diameter	D _e
GSINT (50)	Value of g (acceleration due to gravity) times Δ height/ $\!\Delta z$	gsinθ
ORIK(50)	Value of friction coefficient if ≥ 0 K	K e
	< 0 K _f	K f
PD(50)	Power distribution (per unit length per unit heated wetted perimeter) (m^{-2})	Q/PGENP
CP(50)	Heat capacity of fuel (J m^{-2} K^{-1})	λ
HP(50)	Heated wetted perimeter	p
HF (50)	Fuel centre to suface heat transfer coefficient	h _F
W(50)	Mass flow rate	W
P(50)	Pressure	P
Н (50)	Flow enthalpy	н
HS (50)	Static enthalpy	Hs
X(50)	Thermodynamic quality	x
RHO (50)	Static density	ρ
VOLF (50)	Flow specific volume	${\tt v_f}$
DEFR(50)	Fanning friction factor divided by equivalent	f/D _e
	diameter	C
FK(50)	Kinetic Energy Flow (W)	К
BM(50)	Area times Momentum Flow (kg m ³ s ⁻²)	M
s (50)	Entropy	S
CT(50)	Coolant temperature	T _C
CTS (50)	Coolant saturation temperature	Ts
TFC (50)	Fuel temperature	T
TFS (50)	Fuel surface temperature	$^{\mathtt{T}}_{\mathtt{FS}}$
HT (50)	Fuel centre to coolant heat transfer coefficient	հ _ሞ
THT (50)	Heat transfer indicator	THT
FLUX (50)	Fuel to coolant heat fluxes (W m^{-2})	q
CRSL(50)	Saturation liquid density	$\rho_{m{\ell}}$
CRSV (50)	Saturation vapour density	$\rho_{\mathbf{v}}$
CHSL(50)	Saturation liquid enthalpy	<u>н</u> о

TABLE 4 (continued)

Variable name and dimension	Description	Algebrai symbol
CHSV (50)	Saturation vapour enthalpy	H
CRLP (50)		dρ _ℓ /dÞ
CRVP (50)		dp <mark>v</mark> /d₽
CHLP (50)		dH _l /dP
CHVP (50)		dh_/dp
CSSL(50)	Saturation liquid entropy	s
CSSV (50)	Saturation vapour entropy	${\tt s}_{\tt v}$
DRDW (50)		9b/9 w
DRDP (50)		9b/9 b
DRDH (50)		96/9н
DMDW (50)		9w∕9m
DMDP (50)		9W/9b
DMDH (50)		9м∕9н
DHDW (50)		∂H _s /∂W
DHDP (50)		дн°/9ь
DHDH (50)		ан _s /ан
DKDW(50)		∂ K/ ∂W
DKDP (50)		9K\9b
DKDH (50)		9к∕9н
DTDW(50)		∂Τ _С /∂W
DTDP (50)		∂τ _с /∂₽
DTDH(50)		∂т _С /Эн
DOTDR(50)		x6\q6.x
DOTDM (50)		x.9w/9x
DOTDH (50)		х.дн _s /дх
DOTDK (50)		x.9K/9X
DOTDT (50)		х.Эт _с /Эх
DQDW(50)		W6\p6
DQDP (50)		9 4 \9b
DQDH (50)		9 д √9Н
DQDT (50)		76\p6
DOTDQ(50)		x6/p6.x
AMAT(3,3,5	50) A matrices	А
BMAT(3,3,5	50) B matrices	В
CMAT(3,3,5	50) C matrices	С
DVEC (3,50	,5) Y vectors	Y

TABLE 5

LABELLED COMMON "BCONDS"

Variable Name and Dimension	Description	Algebraic symbol
BPW(10)	Connection control (pressure or flow)	ψ
BHO(10)	Connection enthalphy	η
TIME	Time	t
DELT	Time step	Δt
LC(10)	Connection type indicator	
LCB (10)	Boundary condition type indicator	
NSN (10)	Node number of the first node in flow path	
NFN (10)	Node number of the last node in flow path	
NSC (10)	Connection at start of flow path	
NFC (10)	Connection at finish of flow path	
MMAT(10,10)	Connection matrix	
NP	Number of pipes	
NC	Number of connections	
N	Total number of nodes	

NOTES:

- (i) The LC and LCB indicators correspond to LC and LCB in the connection input data.
- (ii) NSN and NFN are the node numbers at the beginning and end of each pipe, e.g. pipe 3 begins at node NSN(3) and ends at node NFN(3).
- (iii) NSC and NFC are the connection numbers at the ends of each pipe, e.g. pipe 3 joins connections NSC(3) and NFC(3).
- (iv) The connection matrix defines the network. The first index is the pipe number and the second the connection number. If MMAT(KP,KC) is +1 then the first node of pipe KP joins connection KC. If -1 then the last node of pipe KP joins connection KC.
 - (v) The subscripts of BPW, BHO, LC and LCB are connection numbers. The subscripts of NSN, NFN, NSC, NFC are pipe numbers.

TABLE 6

LABELLED COMMON "POWER\$"

Variable name and dimension	Description
PTCTP (10)	Power to coolant in each pipe
PGENP(10)	Power generated in each pipe
PTCT	Total power to coolant
PGEN	Total power generated

TABLE 7

LABELLED COMMON "COEFF\$"

Variable name and dimension	Description	Algebraic symbol
ROUGH (50)	Surface roughness	ε
IFR(50)	Friction correlation indicator	
ISLIP(50)	Slip correlation indicator	
IHTRAN(50)	Implicit heat flux indicator	

TABLE 8

LABELLED COMMON "BAL\$"

Variable name and dimension	Description
TOL	Tolerance in network balance
ITMAX	Max number of iterations in network balance

TABLE 9

LABELLED COMMON "CRIT\$"

Variable name and dimension	Description	Algebraic symbol
PAT(10)	Pressure of choke flow	Pr
TOPEN(10) PTHROT(10)	Break opening time Throat pressure	P
LT(10)	Logical variable which is TRUE if choke flow is present at the connection	С

TABLE 10

LABELLED COMMON "HEAD\$"

Variable name and dimension	Description
HEAD (10)	Title card
DAY (3)	Date (single precision REAL*4 array)

TABLE 11
LABELLED COMMON "STEP\$"

Variable name and dimension	Description
SMIN	Minimum time step allowed
SMAX	Maximum time step allowed
CTAR	Target fractional change
CMAX	Maximum fractional change
RCMAX	Maximum fractional density error at
	a phase change
CINC	Maximum time step increase
DELTR	Estimated time step to achieve "RCMAX"
TMAX	Maximum time required in transient
	calculation
TSFP	Time at which full printout commences
KW (50)	Print mask for partial printout
KFP	Full printout frequency
KKW	Number of nodes in partial printout

the type of listing that has been produced by the standard code for this time step:

IPRINT > 0 full listing

IPRINT = 0 no listing at all

IPRINT < 0 partial listing.

Its use is illustrated in the sample problem.

6.3.4 BCSETM

This routine is called at the start of every time step, including those time steps which are later rejected for reasons given in Section 7. Any explicit connection equations should be solved here. Also boundary conditions and heating power at the end of the time step (at time TIME+DELT from common BCONDS) can be changed.

6.3.5 BPWIMP

All implicit connection equations must be solved in this routine. BPWIMP is called once in each network balance iteration for each connection declared implicit. However, as described in Section 5.3, conditions at connections declared explicit may also be updated in BPWIMP. This allows connections with strong coupling to be handled together, one being declared implicit and the others declared explicit. The pump model in the sample problem is an example of this.

The call statement argument list is

- α the connection number,
- N the network balance iteration number, and
- ω the relaxation constant.

It is recommended that the relaxation constant should be used as described in Section 5.3.

Consider node n which is at the junction of flow path k and connection α . W', P' and H' are given as a function of the connection parameters in a component of Equation (1) (Section 5). During the network iteration at connection α , the connection parameters ψ_{β} , η_{β} at the other end of flow path k are assumed constant. Therefore from Equation (1), Section 5;

$$\mathbf{X}_{\mathbf{n}}^{\prime} = \begin{pmatrix} \mathbf{W}_{\mathbf{n}}^{\prime} \\ \mathbf{P}_{\mathbf{n}}^{\prime} \\ \mathbf{H}_{\mathbf{n}}^{\prime} \end{pmatrix} = \begin{pmatrix} \mathbf{W}_{\star} \\ \mathbf{P}_{\star} \\ \mathbf{H}_{\star} \end{pmatrix} + \psi_{\alpha}^{\prime} \begin{pmatrix} \mathbf{W}_{\mathbf{p}} \\ \mathbf{P}_{\mathbf{p}} \\ \mathbf{H}_{\mathbf{p}} \end{pmatrix} + \eta_{\alpha}^{\prime} \begin{pmatrix} \mathbf{W}_{\mathbf{H}} \\ \mathbf{P}_{\mathbf{H}} \\ \mathbf{H}_{\mathbf{H}} \end{pmatrix}$$

where
$$\begin{pmatrix} W_{\star} \\ P_{\star} \\ H_{\star} \end{pmatrix} = Y_{n}^{(0)} + \psi_{\beta}^{*} Y_{n}^{(3)} + \eta_{\beta}^{*} Y_{n}^{(4)}$$
,

$$\begin{pmatrix} W_{P} \\ P_{P} \\ H_{P} \end{pmatrix} = Y_{n}^{(1)} \quad \text{and} \begin{pmatrix} W_{H} \\ P_{H} \\ H_{H} \end{pmatrix} = Y_{n}^{(2)}$$

are all constant.

A subroutine CONCON is available for calculation of these coefficients. The argument list for CONCON is k and α (input) and

n,
$$W_{\star}$$
, W_{p} , W_{H} , P_{\star} , P_{p} , P_{H} , H_{\star} , H_{p} , H_{H} , ϕ (output) and & m, where & m is an error return if pipe k and connection α do not connect, ϕ is an orientation parameter, and the other arguments are the coefficients. The value of ϕ is 1 if node n is the beginning of pipe k and -1 if node n is the end of pipe k. Thus ϕ W'_{α} is positive for flow out of connection α into pipe k regardless of whether α is at the start or finish of pipe k.

6.3.6 USENDM

This routine is called only at the end of a NAIAD run. Instructions for any final output should be coded into this routine.

6.4 Other Routines

NAIAD

This subroutine, is essentially the mainline of the code. It controls time stepping and output.

INPUT

Most of the input data are read here. Also the power distributions are normalised and various other variables and arrays set to their initial values.

STEADM

The initial condition input data are read and the required initial states calculated.

GEOMP

A table of the geometry of the system is printed. Also the connection variables ψ and η are evaluated from the initial condition of the network.

COEFFM

This is an interface between NAIAD and COEFF. The output at each time step is printed.

SETMAT

The matrix equations are set up from the difference equations and boundary conditions.

SOLVE

The matrix equations are solved and the network is converged.

UPDATE

All the new values of W, P and H are calculated from the converged set of connection parameters. Also heat fluxes, power to coolant and new fuel temperatures are calculated.

CONDIN

This is called by STEADM to read the initial condition input data. STCBAL

This is called by STEADM to balance a connection.

FUEL

Steady state fuel temperatures are generated.

STEADY

The steady state hydraulic difference equations are solved along a flowpath.

STEADZ

The steady state, zero-flow mass and momentum conservation equations are solved along a flowpath. The enthalpy at each node is required input data for this.

DSTM33

The block tridiagonal system of 3 \times 3 matrices with multiple right hand sides is solved.

DQUADS

The relaxation constant ω is calculated.

CRITIC

The connection equations for choked connections are solved.

THROAT

The choke flow parameter γ is calculated.

CONCON

This has been described in Section 6.3.5.

COEFF

All the parameters required to evaluate the coefficients in the matrix equations are calculated. This includes the partial derivatives of h, $H_{\rm S}$, K, ρ and M which are evaluated by repeated calls to the heat transfer and slip routines. Thermodynamic equilibrium of the coolant is assumed.

7. TIME STEP CONTROL

It is usually desirable to take long time steps when conditions are changing slowly, and small time steps when rapid changes are occurring. This is done in NAIAD in the following way. Let

$$\Phi = \max_{\substack{\text{all} \\ \text{nodes}}} \left(\left| 1 - \frac{P'}{P} \right|, \left| 1 - \frac{H'}{H} \right| \right).$$

The next time step is chosen to meet the following criteria:

- . Always in the range SMIN to SMAX.
- . Never larger than CINC * (t'-t) ,
- . Otherwise equal to $\frac{\text{CTAR}}{\Phi}$ * (t'-t) .

Thus a new t' is defined and the old t' becomes t.

This time step is then completed to the stage at which the new Φ has been calculated. If Φ < CMAX or t' - t = SMIN, then the time step is accepted, otherwise the time step is redone with

$$\Delta t = Max [SMIN, \frac{CTAR}{\Phi} * (t'-t)].$$

There is another factor in the time step control which has been omitted in the above discussion. To derive the finite difference equations, the approximation

$$\frac{\partial \rho}{\partial t} \approx \frac{\partial \rho}{\partial x} \cdot \frac{(x'-x)}{\Delta t}$$
 ... (2)

is used. This is a very poor approximation if there is a change from liquid to two-phase or vice versa during the time step. If such a change of phase is estimated to have occurred during the time step, an estimate of ρ_e , the fractional density error resulting from use of Equation (2), is obtained as follows. At this stage, all W', P' and H' are known but no other parameters have been evaluated at the end of the time step. First, we obtain estimates (indicated by a superscript *) of the saturation line densities and enthalpies at the end of the time step.

$$H_{\ell}^{\star} = \frac{dH_{\ell}}{dP} (P'-P) + H_{\ell}, H_{V}^{\star} = etc.$$

We also define

$$\rho^* = \frac{\partial \rho}{\partial x} \cdot (x' - x) + \rho ,$$

$$x^* = \frac{H' - H^*_{\lambda}}{H^*_{\lambda} - H^*_{\delta}} .$$

and

If xx* \geq 0 then there is no liquid two-phase transition and ρ * is a good estimate of ρ '. Then $\rho_e = \left|1-\rho*/\rho'\right| \approx 0$.

Two-phase to liquid
$$x > 0$$
, $x* < 0$

Because water is nearly incompressible

$$\rho^{\bullet} \approx \rho_{\ell}^{\star}$$

hence

$$\rho_{\rm e} = |1 - \rho * / \rho_{\ell}|$$

Liquid to two-phase
$$x < 0, x^* > 0$$

We estimate ρ' from

$$\frac{1}{\rho^*} \approx \frac{x^*}{\rho_v^*} \; + \frac{(1-x^*)}{\rho_v^*} \quad . \label{eq:rho_v}$$

This rather poor estimate (slip is neglected) is sufficiently accurate for the present purpose. Then

$$\rho_e = |1-\rho*/\rho'|$$
.

Now if we assume that x varies linearly during the time step the error in density only occurs (in time) after the phase change has taken place. The time at which the phase change takes place is estimated as

$$t_0 = t + \frac{(t'-t)x}{x-x*}$$

and it is assumed that $\rho_{\mbox{e}}$ varies linearly with time in the interval t_0 to t'. If

$$\rho_e$$
 < RCMAX

the time step density error is deemed acceptable. Otherwise a time step of

$$(t_0-t) + (t'-t_0) \frac{RCMAX}{2\rho_e}$$

is tested against the criteria described in the first paragraph. All the time step control parameters (given in capitals above), are specified

in the input data and stored in the labelled common STEP\$.

8. INPUT DATA

8.1 General

The finite difference method relies on the continuity of the derivatives being approximated by finite differences. Unfortunately, some problems involve discontinuities (or, at least, large spatial derivatives) in some parameters, e.g. flow path area, surface heat flux. Such discontinuities are best treated by placing nodes close to either side of and at equal distances from the discontinuity.

Except for the title card, all data are read with the AAEC free input subroutine SCAN [Bennett & Pollard 1967]. Data may be punched anywhere on a card with numbers separated by blanks or commas. Alphabetic comments may be placed amongst the data and comment cards (* in column one) may be used. Other features are:

. repeat notation

e.g. 5*1.47

is equivalent to 1.47 1.47 1.47 1.47

fixed increment notation

e.g. 1. (0.5)3.

is equivalent to 1. 1.5 2. 2.5 3.

- . a decimal point is not needed for real numbers with integral values
- the data cards are listed in the output.

All input and output data are in SI units.

8.2 Title Card

HEAD Title card

8.3 Controls

tNP NP is the number of flow paths in the network. If negative a printout of the pipe data interpreted by the code is produced after each set of pipe data is read.

NC Number of connections in the network.

GRAV Acceleration due to gravity, $g (m s^{-2})$.

TOL Tolerance in network balance (suggest 10⁻⁸).

ITMAX Maximum number of iterations in network balance (suggest 15).

SMIN Minimum time step to be used (s)

SMAX Maximum time step to be used (s).

CTAR Target maximum fractional change in P or H in a time step (suggest 0.01).

CMAX Maximum allowable fraction change in P or H in a time step (suggest 0.02)..

RCMAX Maximum allowable fractional density error in a time step involving a liquid to two-phase transition (suggest 0.01).

CINC Maximum ratio between successive time steps, e.g. 2 means at most double the time step.

TMAX Maximum time required (s).

KFP Frequency of full listing, e.g. 5 means full listing every 5th time step and partial listing at other time steps. A partial listing covers only those nodes specified below in KW, and, if KFP positive, details of time step change and elapsed CPU time. If no nodes are specified below and KFP is negative, then no partial listing appears.

TSFP A full printout will be produced at every time step after this time.

A dashed line is printed in the output immediately after this data is read and printed.

8.4 Flow Path Data

The following set of data is required for each pipe or flow path beginning with pipe 1. Each set must start on a new card:

NODE Number of nodes in flow path.

KA Connection at start of flow path.

KB Connection at end of flow path.

PGENP Total power generated in this flow path (W).

ISLIP NODE integers specifying which slip relation is to be used at each node:

- 0 homogeneous (no slip)
- 1 Jones
- 2 CISE
- 3 Beattie
- 4 Beattie small bubble
- 5 Fauske
- 6 Moody
- 7 modified CISE
- 8 Smith.

IFR NODE integers specifying which friction relation is to be used:

±1 - homogeneous

 ± 2 - bubbly

±3 - annular

±4 - small bubbles.

If negative the boiling wall extension will be used.

IHTRAN NODE integers specifying differencing of the surface heat transfer coefficient:

0 - explicit heat transfer coefficient

1 - implicit heat transfer coefficient.

KW NODE integers specifying which nodes are to be printed in partial printouts:

0 - do not print the node

1 - print it.

NODE value of the axial displacement (m) along the flow path (arbitrary origin but strictly increasing).

SINT NODE-1 values of $\frac{\Delta \text{ height}}{\Delta z}$ for segment of flow path between adjacent nodes.

DE NODE values of equivalent diameter D_e (m).

If DE(I) is given as zero it is calculated from A(I) assuming a round pipe.

A NODE values of cross sectional area a (m^2) .

If A(I) is given as zero, it is calculated from DE(I) assuming a round pipe.

HP NODE values of heated wetted perimeters p (m).

If HP(I) is given as zero, it is calculated from DE(I) assuming a round pipe.

ROUGH NODE values of absolute surface roughness ε (m).

CP NODE values of fuel heat capacity λ (J m⁻² K⁻¹).

If zero, generated power passes directly to coolant.

HF NODE values of fuel centre to fuel surface heat transfer coefficient h_F (W m⁻² K⁻¹).

If given as zero set to 10^{10} (zero resistance).

ORIK NODE-1 values of friction coefficient.

If > 0 then taken as K_e If < 0 then taken as $-K_f$. NODE values specifying an equivalent heat flux distribution P_D (m⁻²) (the given distribution is normalised within the code so that $P_{Di} = Q_i/PGENP$ hence NODE-1

$$\sum_{i=1}^{\Sigma} (P_{Di}p_i + P_{Di+1}p_{i+1}) (z_{i+1} - z_i) = 1 .$$

A dashed line is printed in the output after each set of flow path data is read and printed.

8.5 Connection Data

The following set of data is required for all connections which are not internal. The first set must start on a new card.

KC CONNECTION number

LC CONNECTION type

0 - explicit

1 - internal

2 - choked (burst)

3 - implicit.

Only one flow path can connect to a choked connection, and the connection number allotted to a choked connection must be larger than that of the connection at the other end of the flow path involved.

If LC = 0 or 3

LCB type of connection control

1 Flow

2 Pressure.

If LC = 2

PAT Pressure downstream of burst (final receiver pressure)
(Pa)

TOPEN Opening time of break(s).

A zero marks the end of the connection data. A dashed line follows the listing of the connection data in the output.

8.6 Initial Condition Data

The initial condition calculation is a series of steady state calculations. Any number of calculations may be done on one flow path. If a steady state calculation is not done on a flow path, then its initial state must be specified in the special routine BINITM.

In the following specification an asterisk as a superscript on a flow, pressure or enthalpy means that it is specified by two numbers IP and Q where

IP = Pipe number or zero.

If IP = 0

Q = COND = the actual value of the quantity.

If IP \neq 0

Q = Node number. Then the condition is set equal to its counterpart at the Qth node of pipe IP.

Also, this node number Q is the position in the flow path containing the node, not the node number appearing in the code output. This allows extra nodes to be added with minimum impact on the input data.

The following set of data is required for each steady state calculation:

KP Pipe number.

IOPT Specifies which type of calculation is required.

- ±1 Conditions known at a single node.
- ±2 Pressure known at two nodes and enthalpy at one of them.
- ±3 Zero flow in flow path, enthalpy known at all nodes, pressures determined from gravitational head.
- the Conditions known in all but one of the flow paths terminating at a connection. The enthalpy and flow at the terminal node of the unknown flow path are chosen to give zero net enthalpy and mass flow through the connection. The pressure is chosen equal to that of the other terminal nodes. This process is called a connection balance.
- to Constant gravity head pump with head and flow path power determined to give connection balances at both ends of flow path.

If IOPT is negative, a listing of the calculated steady state for this flow path will be produced before any further initial state calculations are done.

The data required for each option is given below:

 $IOPT = \pm 1$

J Node number

 W_J^* Flow at node J in flow path KP (kg s⁻¹)

- P_{T}^{*} Pressure at node J in flow path KP (Pa)
- H_J^* Enthalpy or thermodynamic quality at node J (J kg⁻¹) in flow path KP.

For example, to produce a steady state in pipe 1 with flow 4 Mg s^{-1} and at node 2 a pressure and quality of 8.6 MPa and -0.14 respectively

The steady state finite difference equation will then be solved along flow path KP in both directions from node J.

 $IOPT = \pm 2$

J Node number

 W_J^* Estimate of flow at node J in flow path KP (kg s⁻¹)

P* Pressure at node J in flow path KP (Pa)

 H_{J}^{*} Enthalpy or thermodynamic quality at node J (J kg^{-1}) in flow path KP

K Node number

 P_{K}^{*} Pressure at node K in flow path KP (Pa).

The flow will then be found by iteration and finally the steady state finite difference equations solved along flow path KP:

 $IOPT = \pm 3$

J Node number

 $P_{,T}^{\star}$ Pressure at node J of flow path KP (Pa)

H Enthalpies at each node of flow path KP (J kg-1)

TFC Fuel temperatures at each node of flow path KP (K).

The steady state zero flow momentum equation will then be solved along flow path KP. Any fuel temperatures specified to be zero are set equal to the coolant temperature for that node:

 $IOPT = \pm 4$

KC Connection number.

Conditions at the node of flow path KP which connects to connection KC will be determined from a connection balance. The steady state finite difference equations will then be solved along flow path KP:

 $IOPT = \pm 5$

J Segment at which pump is to be inserted (between node J and node J + 1).

Conditions at each end of the pipe will be determined by connection balances. The head (Δ height) at segment J and power to coolant are then varied to achieve the required conditions.

Zero marks the end of steady state data. This is the end of the standard NAIAD input data. Any input for special routines will follow.

9. CATALOGUED PROCEDURES

All three procedures are used in the test problem.

9.1 NAIADCLG

This procedure is used to compile, link and execute a composite program. The procedure has four steps:

(a) PREP - In this step common blocks and other groups of cards can be copied into the FORTRAN source decks by including "CINSERT" cards in the FORTRAN program. By this means, the tedious process of duplicating common blocks is avoided. The format of this card is

where the letters CINSERT followed by a blank must be in columns 1 to 8 of the source card and COM1, COM2, COM3, etc. are the names of common blocks required in the FORTRAN routine. Any of the common blocks of Tables 4 to 11 may be inserted in this way, BLANK common having the label "BLANK". The default MAIN routine (Section 6.3.1) may be generated with a

CINSERT MAIN

card and the IMPLICIT REAL*8 statement with a CINSERT R8

card.

- (b) FORT FORTRAN H compilation step
- (c) LKED linkage editor
- (d) GO execution of the composite program load module. This load module may be placed in a private program library by including two symbolic parameters LIB and MEM. LIB is the private library name and MEM the member name under which the load module is to be stored.

9.2 NAIADG

In this procedure composite program load modules specified by the

two symbolic parameters LIB and MEM can be executed. The GO step is the only step in this procedure. If LIB and MEM parameters are not specified, then the standard NAIAD code is executed.

9.3 PROGRUN

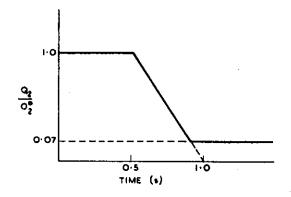
The plot program is executed in this procedure. It is used as shown in the sample problem.

10. SAMPLE RUN

The sample problem is the blowdown of a CANDU nuclear reactor through a small hole in an outlet header. The problem has been chosen to illustrate the use of most NAIAD options, to show simple examples of special routines and to demonstrate the advantages of NAIAD. In the sample run, a composite program for this problem is formed, the problem is run twice with target fractional changes of 10 and 1 per cent, and graphs are plotted comparing the coarse and fine solutions.

Both primary and secondary coolant loops are simulated together with their pumps and the steam generator coupling the two loops. The flow network is shown in Figure 6. A very coarse spatial mesh is used to keep the computation time and output quantity small. The lower part of the primary system simulates the flow path from inlet header to outlet header through the feeders and fuel channels. The conditions in the 195 channels are assumed identical so that they can be combined into one flow path. Similarly, in the upper part of the primary system, the 15600 steam generator tubes are combined into one flow path and the sets of six pipes connecting the headers to the steam generators are combined into one flow path. The secondary system includes a constant volume pump, the steam generator and a heat sink. The transient is initiated

FIGURE 4
Nuclear Power for Sample Problem



by a burst in the outlet header and it is assumed a reactor shutdown occurs.

10.1 Nuclear Power

The power generated in the fuel, which is in flow path 2, is assumed to vary as shown in Figure 4. It is calculated at the beginning of each time step in the special routine BCSETM.

10.2 Heat Sink

The heat removed from the secondary coolant within the heat sink, which is in pipe 6, is

$$Q_6 = \left(\frac{\overline{T} - T_{sink}}{\overline{T}^0 - T_{sink}}\right) Q_6^0 ,$$

where \bar{T} = mean temperature of coolant in the heat sink. Q_6 is calculated in the special routine BCSETM at the start of each time step using the coolant temperatures from the previous time step.

10.3 Steam Generator

The heat flow from the primary side (pipe 1) to the secondary side (pipe 5) of the steam generator is assumed proportional to the mean temperature difference:

$$Q_5 = -Q_1 = \left(\frac{\overline{T}_p - \overline{T}_t}{\overline{T}_p^0 - \overline{T}_t^0}\right) Q_5^0 ,$$

where \bar{T}_p = mean temperature of primary coolant in the steam generator, and \bar{T}_t = mean temperature of secondary coolant in the steam generator. The heat flows are calculated each time step in the special routine BCSETM.

10.4 Makeup System

The makeup or pressure control system is simulated by a flow boundary condition at connection 1. The flow ψ_1 is a function of the outlet header pressure ψ_3 as shown below, except that a limit of 20 kg s⁻² is imposed on the rate of increase or decrease in the flow rate. It is calculated at the start of each time step in BCSETM.

10.5 Secondary Pump

The secondary pump is represented by connection 7 which is an implicit flow controlled connection, connection 5 which is an explicit flow controlled connection, and a special routine BPWIMP. As the flow is always from connection 5 to 7 via 6, the quantities to be determined from the pump characteristic are

$$\psi'$$
, ψ' , and η' ,

which are the connection parameters for connections 5 and 7. Let n be the node at connection 5 and m be that at connection 7. Then the pump equations are

$$\psi_5' = \psi_7'$$

(the pump is assumed to have zero volume),

$$\frac{\psi^{i}}{\rho_{m}^{i}} = \frac{w_{m}^{0}}{\rho_{m}^{0}}$$

(the volumetric flow into the pump is constant), and

$$\eta_{5}^{1} - H_{m}^{1} = H_{n}^{0} - H_{m}^{0}$$

(the enthalpy change across the pump is constant).

To avoid calling SSTATE and SSLIP to obtain the density of the coolant at the end of the time step, the following estimate is used

$$\rho_{m}^{\prime} = \rho_{m} + \left(\frac{\partial \rho}{\partial W}\right)_{m} \left(W_{m}^{\prime} - W_{m}\right) +$$

$$\left(\frac{\partial \rho}{\partial P} \right)_m (P_m^{\,\prime} - P_m^{\,}) + \left(\frac{\partial \rho}{\partial H} \right)_m (H_m^{\,\prime} - H_m^{\,}) \quad .$$

These four equations have introduced new unknowns W_m^{\dagger} , P_m^{\dagger} and H_m^{\dagger} . The required three equations are components of Equation (1) Section 6. These equations are obtained by the CONCON subroutine. The seven equations are solved for ψ_1^{\dagger} , ψ_2^{\dagger} , and η_2^{\dagger} in BPWIMP.

Note that although connection 5 is an explicit connection its boundary conditions are set in BPWIMP and hence updated each network balance iteration.

10.6 Primary Pump

The primary pump is represented as a fixed gravitational head. The head required is determined in the steady state calculation. No special routines are needed. The gravitational potential energy from this head appears in the primary coolant, hence in the initial steady state the heat flow through the steam generator is the sum of this gravitational power and the nuclear power.

FIGURE 5
Makeup System for Sample Problem

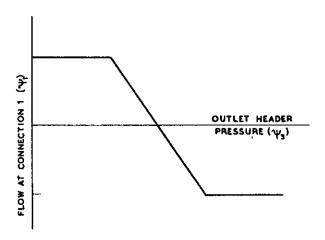
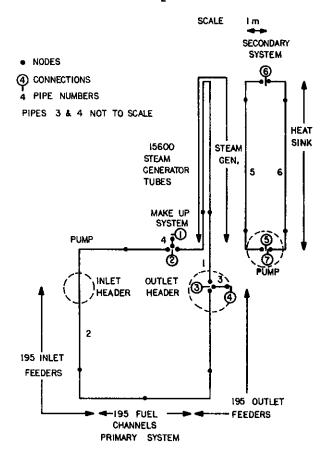


FIGURE 6
Sample Problem



10.7 Initial State

In the coarse time step calculation, the first two initial state calculations demonstrate the use of initial state options 3 and 2 respectively; the rest set the initial state for the transient calculation using the other three options.

10.8 Transient Calculation

The transient is initiated by a burst of area 1 cm2 in the outlet header, pipe 3 being entirely within the header. The receiver pressure at the break is reduced to atmospheric over the first ten milliseconds. The difference between the receiver pressure and the header pressure is not sufficient to cause choking until the quality at the break becomes positive. Once choking occurs, the throat and receiver pressures are no longer equal and the throat pressure only changes slowly as shown in Figure 7. The time step in the fine time step calculation (Figure 8) which is at the lower limit of 0.1 milliseconds before choking, increases rapidly soon after choking until reaching 0.1 seconds. Roughly, this time step is maintained for the remainder of the calculation. at the break rises rapidly to 4 kg s⁻¹. This is not sufficient to depressurise the system and, in fact, after 0.5 seconds the makeup system flow has exceeded the break flow. However, the drop in nuclear power which begins at 0.5 seconds, coupled with the nearly constant heat removal to the heat sink, causes a slow decrease in system pressure.

The flow at the break from the coarse and fine calculations are compared in Figure 9. The difference is quite tolerable for many applications and illustrates the value of coarse NAIAD calculations where precise results are not required.

10.9 Sample Run Output

A listing of the sample run output is given in Appendix A. A listing of the input is not given as the SCAN input subroutine lists all input data as it is read and the FORTRAN compiler lists all the user routines. Note however that the IMPLICIT and COMMON statements in these routines were inserted by the CINSERT procedure described in Section 9. The graphs from this sample run are shown in Figures 7, 8 and 9.

11. PLOTTING

Graphs are easily produced by including special instructions for writing a plot data file in the special routine BINITM and TDUMP, and executing the NAIAD plot catalogued procedure PROGRUN which reads this plot file and plots the graphs specified in the input data. In this

FIGURE 7
Pressure at Break (Fine Calculation)

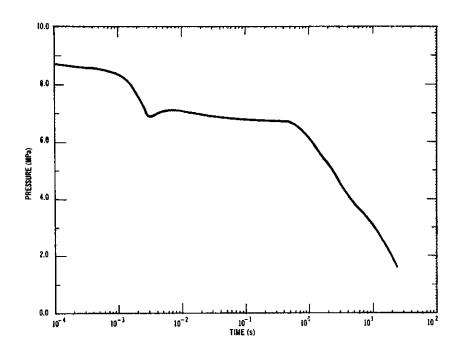
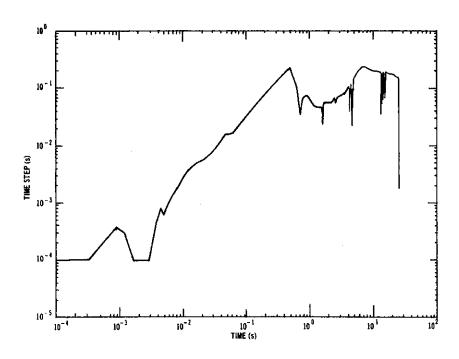
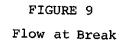
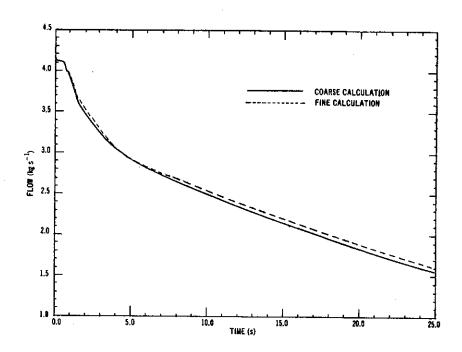


FIGURE 8
Time Step (Fine Calculation)







way, graphs can be produced of data in the plot data file from one or more NAIAD runs, e.g. pressure at a node as a function of time can be plotted for a number of NAIAD runs, pressure can be plotted against flow rate, etc. Points or lines of many kinds can be used and labelled, axes can be labelled and headings written. Also, experimental data can be plotted for comparison with the NAIAD results.

11.1 Plot Data File

This file is a sequential data set, each record corresponding to a single time step and containing up to 200 double precision words.

Output from several runs may be assembled into one long file with the start of each pseudo file marked by TIME equal zero. Thus, the first item of each record must be the variable TIME.

Two routines are available for writing the plot data file. They are TAPSET which locates the required pseudo file and TDUMP which writes a record.

The routines are used as follows:

(i) CALL TAPSET(IU,NF)

where IU - FORTRAN logical unit number;

NF - output control;

NF #0 - replace NF th pseudo file with results of this run. All old files past this point will then be lost; and

NF=0 - add output of this run to the end of existing file.

TAPSET is normally called from the special routine BINITM.

(ii) CALL TDUMP(IU, TIME, 1, A1, N1, A2, N2, ..., AK, NK)

where IU - FORTRAN logical unit number;

Al,..., AK - variable or array names; and

N1,...,NK - number of data items to be written from the corresponding variable or array.

TDUMP is normally called from the special routine TDUMPM. Both these routines are used in the sample problem.

11.2 Plot Program

This program is separate from NAIAD and runs in two phases, a data input phase and a plot phase. The data input phase requires the following data (all read using SCAN):

IFN - pseudo file number;

NENT - number of items in each record (max. 200);

NG - number of these items actually required for plotting (max. 20); and

NP - an array of NG numbers giving the positions in the NAIAD output record of the NG items required.

For example, considering the sample problem (Section 10) which has 22 nodes and 6 flow paths, the record structure is

TIME	position 1
w	positions 2 to 23
P	positions 24 to 45
х	positions 46 to 67
PTCTP	positions 68 to 73
PTCT	position 74
DELT	position 75

The input data needed to retrieve TIME and flow, pressure and quality at the break (node 12) are:

1 75 4 1 13 35 57.

The total number of words that can be retrieved in this way is limited by storage to 5000.

When we enter the plot phase, the items retrieved are identified by the numbers 1 to NG rather than their initial positions in the NAIAD record. Thus in the above example, TIME is now item 1, flow at the break is item 2, etc. The plot phase is controlled by 'keyword and data' cards. Each individual plot is initiated by a card beginning with the keyword PAPE (a PAPE card). This card defines the graph paper required. The PAPE card defines the plot axes explicitly or implicity. The explicit form is:

PAPE XP YP XMIN XMAX YMIN YMAX

where |XP| = 1ength of x-axis in inches,

|YP| = length of y-axis in inches,

XP, YP < 0 linear scale, or

> 0 log scale,

XMIN = minimum x (usually time) value required,

XMAX = maximum x value required,

YMIN = minimum y value required, and

YMAX = maximum y value required.

The implicit form is:

PAPE O XP YP $\begin{cases} O & XMIN & XMAX \\ NX & IX & (1), \dots, IX & (NX) \end{cases}$ $\begin{cases} O & YMIN & YMAX \\ NY & IY & (1), \dots, IY & (NY) \end{cases}$

where the upper and lower lines inside the braces are alternatives. In this form, the x limits are either on the PAPE card (NX = 0) or they are the range of a specified subgroup of the items loaded in the input phase. NX is the number of items in this subgroup and the item numbers are IX. The y limits are specified in the same way. Returning to our example

PAPE 0 -10 -8 1 1 1 4

specifies 10 in. \times 8 in. linear graph paper with the \times limits given by the range of TIME values and the y limits given by the range of quality values. If NY is negative, the cards

LINE IX(1) IY(1) 1

LINE IX(1) IY(2) 2

LINE IX(1) IY(NY) NY

are executed after the PAPE card. If the x limits were given explicity, then IX(1)=1 is assumed. The LINE card is described below.

Once the PAPE card has been read, data may be plotted on the graph by means of LINE or PNTS cards.

LINE IX IY NL

connects the given data by straight lines

where IX = item number for the x variable,

IY = item number for the y variable,

NL = line type, where NL=1 is a solid line, and $NL=2,\ldots,10$ are distinct dashed lines.

LINE O NP NL X(1),Y(1),...,X(NP),Y(NP)

connects the NP given data pairs by straight lines of type NL. This card can be used for plotting experimental results.

PNTS IX IY NS ISIZE

draws a plot symbol at each data point,

where IX = item number for the x variable,

IY = item number for the y variable,

NS = plot symbol number, where NS=1,...,8 are distinct
 plot symbols,

ISIZE = plot symbol size, where the actual size is 0.04*
ISIZE inches.

PNTS 0 NP NS ISIZE X(1), Y(1),...,X(NP), Y(NP) draws a plot symbol at each given data point.

The plot can be labelled by the following keyword cards:

HEAD "TITLE FOR PLOT"

NAMX "X AXIS LABEL"

NAMY "Y AXIS LABEL"

where the required labels are placed between the double quote characters ".

Legends describing the plot symbols or lines (see Figure 8) may be added to the plot by use of further keyword cards. Each legend entry consists of two fields. The first field contains a symbol or line while the second contains text.

LEGO HO VO

defines the top left hand corner or 'origin' of the legend. HO and VO are given in inches from the bottom left hand corner of the graph. If this statement is omitted, HO = |XP| + 0.5, VO = |YP| - 0.5 is assumed.

LEGH "HEADING" provides a heading for the legend.

LEG K "LABEL" generates a legend entry, where

K = 1 gives the latest plot symbol followed by LABEL

- = 2 gives the latest line type followed by LABEL
- = 0 just gives LABEL starting in the symbol or line field.

LEGS SIZE WIDTH SPACE

enables changing of the controlling parameters.

SIZE - character height in inches (default 0.14)

WIDTH - width of symbol or line field in inches (default 1.2)

SPACE - entry spacing (vertical) in inches (default 0.25).

The data may also be scaled for plotting by use of the statement

CONS A B

where all 'y-ordinate' data will be scaled so that the plotted data is A + By. The scale remains in force until a new PAPE or CONS card is read.

Finally, it is possible to return to the data input phase by use of the command

DATA

This allows a new set of data to be read which can then be plotted either on a new graph or on the previous graph if no new PAPE card has been read.

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APPENDIX A

PARTIAL LISTING OF SAMPLE RUN OUTPUT

```
JOB 322
//GDTNAIAD JOB ( PHI 13994/P22PHLCA , N1), G.D. TRIMBLE,
                                       CLASS=B.
                                           TIME=7
11
// EXEC NAIADCLG, PARM. FORT= 'OPT=2', LIB= 'GDT. NAIADP', MEM=TESTCASE
//SYSIN CO * GENERATED STA
ILF142I - STEP WAS EXECUTED - COND CODE 0000
                                                                                        GENERATED STATEMENT
1EF3731 STEP /PREP / START 76145.1752
1EF3741 STEP /PREP / STUP 76145.1752 CPU
                                                                                                                                       OMIN 02.30SEC MAIN 40K LCS
                                                                                                                                                                                                                              OΚ
                                                               OOO(HEX)
*** CCNDITION CODE =
IEF1421 - STEP WAS EXECUTED - COND CODE 0000
1EF3731 STEP /FURT
1EF3741 STEP /FURT
                                                             / START 76145.1752
/ STOP 76145.1752 CPU
                                                                                                                                                                                                                              0K
                                                                                                                                      JMIN 04.64SEL MAIN 24UK LCS
*** CONCITION CODE =
                                                              000(HEX)
1EF1421 - STEP WAS EXECUTED - COND CODE 0000
| 1721 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 1722 | 
                                                                                                                                     OMIN 03.52SEC MAIN 96K LCS
                                                                                                                                                                                                                              0к
//SYSIN DD * GENERATED STATEMENT

IF142I - STEP WAS EXECUTED - COND CODE 0000

IEF373I STEP /GU / START 76145.1753

IEF374I STEP /GO / STOP 76145.1759 CPU OMIN 54.80SEC MAIN 220K LCS

*** CCNDITION CODE = 000(HEX)

// EXEC NAIADG, MEM=TESTCASE, LIB="GDT.NAIADP"

//GC.FTC8F001 DD DSN=GDT.IESTPLT, DISP=OLD

GENERATED STATEMENT
 //SYSIN DO * GENERATED STATI
IEF142I - STEP WAS EXECUTED - COND CODE 0000
IEF373I STEP /GO / START 76145.1759
IEF374I STEP /GO / STOP 76145.1804 CPU
                                                                                    GENERATED STATEMENT
                                                                                                                                        2MIN 07.52SEC MAIN 220K LCS
 *** CGNDITION CODE = 000(HEX)
 // EXEC BUFFPROG.PRG=VPLUT
//SYSUT2 CD SYSUUT=F
 / START 76145.1804
/ STUP 76145.1804 CPU
                                                                                                                                         OMIN 00.245EC MAIN 14K LCS
                                                                                                                                                                                                                               ÚΚ
  *** CCNCITION CODE = OOO(HEX)
 // EXEC PROGRUN, LIB= GOT. NATAD , MEM=PLOTPGM, REGGO=140K
  //GC.AEPLOT DD SYSOUT=F
  //GO.FTO8FOO1 DD DSN=GDT.TESTPLT.DISP=SHR
                                                                                           GENERATED STATEMENT
                                  DD +
  //SYSIN
  IEF1421 - STEP HAS EXECUTED - COND CODE 0000
  IEF3731 STEP /GU
IEF3741 STEP /GD
                                                     / START 76145.1804
/ STOP 76145.1804 CPU
                                                                                                                                     OMIN 05.22SEC MAIN 110K LCS
  *** CONDITION CODE = 000(HEX)
  IEF375I JOB /GOTNAIAD/ START 76145.1752
IEF376I JOB /GDTNAIAD/ STUP 76145.1804 CPU 3MIN 18.24SEC
  HIGHEST CONDN CODE = 000(HEX)
```

DATE 75.349/14.03.42		THE ORIGINAL DECK. THEY HAVE BEEN INSERTED DURING THE PREP STEP IN	RESPONSE TO THE CINSER! CARD.		
05/360 FOPTRAN H	NAME:= MAIN.OPT=02.LINECNT=55.SI7F=D000K. Source.frcoic.nolist.nodeck.load.nomap.nofdit.noid.noxref	•	NAME = MAIN.OPT=02.LINFCNI=55.SIZE=0000K.	SGURCF.EGGDIC.NOLIST.NODECK.LOAD.NOMAP.NOFOIT.NOID.NOAREF	S = 3 *PROGRAW SIZF = 1A0
	COMPILER OFIIONS - NAMES SOURCE CINSERT MAIN	CALL NATAD STOP END	A H H H H H	SGURCF .E	SOURCE STATEMENTS =
LEVAL 20.1 (AIV. 7])	COMPILER OF	ISN 0003 ISN 0003 ISN 0004	elder in Shorter	WINDITCHS IN EFFECTS	DAUGS SOLISTICS

107K BYTES OF CORE NOT USED

OSTATISTICS NO STACHOSTICS GENERATED

sesses END OF COMPILATION sesses

05/360	

DATE 75,349/14,03,45		~					_	THESE CARDS WERE NOT INCLUDED IN	THE ORIGINAL DECK. THEY HAVE BEEN	INSERTED DURING THE PREP STEP IN	RESPONSE TO THE CINSERT CARD.										>															
20.1 (AUG 71) OS/360 FORTRAN H	OPTIONS - NAM SOU		COMMON TAKK 1V COMMON Z(SO) +A(SO) +OE(SO) +GSINT(SO) +ORIK(SO)	M (05) 4 (05) M	HS (50) • X (50) • F	HT (50) +THT (50)	COREGO: CRSL (50) * CRSV (50) * CRSV (50) * CRV (50) * CRV (50) * CRV (50)	CSSL (50) CSSV		DADA (DC) NOMO	COMMON DECIMATION OF THE CONTRACTOR OF THE CONTR	DTD#(50) + DTDP	DOTOR (50) . DOT		STATE OF THE PROPERTY OF THE P	ON /BCONDS/ BPW(A NTN(10) •NSC(10) •NFC(10) •MMAT(10•10) •NP•NC•N	•	. LT 	VETERS CATE.	COTECT TO COTE TO CO	COMMON /BALS/ TOL·ITMAX	*NO-TSINK-FLIM-IC-CMASS -FLIM-NEV-NOT-BRIEF	e	(5) / (TFC (NP1)	PS # PGENP (6) / (CT (NK	35 35		 The Cart of the Ca	TAPSET (8.NFP)	HETUP	C FATON BUNETA	IF (TI)	i	C MAKE FERT VALVE CONTROLLED DI DOLLE.	
LEVEL 20.1 (4		ISN 0002 ISN 0003		9000 NSI		15N 0008					ISN 0015		8100 XVI	000	ISN 0020	į	U	15N 0022		ISN 0024	15N 0025	2	15N 0027		EXOD XXI			15N 0032	15N 0034		1500 NS1		15N 0058			2400 NSI

IF (IPPINT.NE.0) WRITE(3.120)

120 FORMAT(*0:*T 6.*TIME**TI9.*THROAT PRESSURE AND GUALITY.*

\$ T50.*INLET FDR FLOW AND GUALITY.*TT9.*CHOKED*/)

40 WRITE(3.100) TIME*P(NR)*X(NR)*W(7)*X(7)*LT(4)

100 FORMAT(IPSG15.6.4X.63.0)

IF (IIME*EQ.0.D0) GO TO SO

CMASS=DFLT**500*(W(NH)*WSTRT) +CMASS

50 CALL TDUMP(8.TIME.1.W.N.P.N.X.N.PTCTP.NP 6 .PTCT.1.DELT.1)

PGENP (2) =PNUC*DMAX1 (DMIN1 (1.00.2.00-2.00*TEND).7.0-2)

TEND=TIME+DELT

0052 0.053 0054

N N N

ISN 0051

RETUPN

U

ISN 0055 ISN 0056 ISN 0058

0059 0900

0061 0063 SN 0064 ISN 0065

C NUCLEAR POWER

ENTRY TOUMPH(IPRINT)

C POWER THROUGH STEAM GENERATOR
PGENP(5) = PK*(TFC(NP1)+TFC(NP2)-TFC(NS1)-TFC(NS2))
PGENP(1) = -PGENP(5)

IF (DARS (DW) .LE .WFVRC) GO TO 20 WFV=W (NFV) +DSIGN (WFVRC+DW)

20 HP# (1)=WFV

WF VHC=WRMAX+DELT

DESENTATE (NEV)

ISN 0044 ISN 0044 ISN 0045 ISN 0045 ISN 0047

SN 0049 15N 0050 PGENP(6) = PS*(CT(NK1)-CT(NK2)-TSINK+2.00)

C POWER TO HEAT SINK

STATISTICS NO DIAGNOSTICS GENERATED

oobsee END OF COMPILATION sesses

1832 69 PROGRAM SIZE SOURCE STATEMENTS = *STATISTICS*

SOURCE. EBCDIC. NOLIST. NODECK. LOAD. NOMAP. NOEDIT. NOID. NOXREF

NAME = MAIN.OPT#02.LINECNT#55.SIZE#0000K.

OPTIONS IN EFFECT *OPTIONS IN EFFECT*

ENTRY USENDM WRITE (3.110) CHASS FORMAT(*0101AL STEAM IN CONTAINMENT **1PG15.5.* KG*)

RETURN END

U

ISN 0070

110

ISN 9066 ISN 9067 ISN 9068 ISN 9069

o

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DATE 75.349/14.03.51
     OS/340 FORTRAN H
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DATE 75.349/14.03.51			THESE CARDS WERE NOT INCLUDED IN THE ORIGINAL DECK. THEY HAVE BEEN INSERTED DURING THE PREP STEP IN RESPONSE TO THE CINSERT CARD.	>			
OS/340 FORTRAN H	LER OPTIONS - NAME= MAIN,OPT=02.LINECNT=55.SIZE=0000K. SOURCE.ERCPIC.NOLIST.NOPECK.LOAD,NOMAP.NOEDIT.NOID.NOXREF Subpoutine Apwimp(KI.NIT.OM) CINSERT P8.RLANK.BCONDS IMPLICIT REAL** (A-H.O-Z)	COMMON MARK IV COMMON PD(50)+A(50)+DE(50)+GSINT(50)+ORIK(50) COMMON PD(50)+CP(50)+HF(50) COMMON WE(50)+P(50)+H(50) COMMON WE(50)+X(50)+H(50) COMMON ME(50)+X(50)+RF(50)+FF(50)+FF(50)		COMMON ACCIONS COMMON ACCONDS/ NFN (10) *NSC (10) OGICAL L NT PUMPED FROM IOW CONNECTION SATA KO-L*KPI*KP 5**FALSE** 6*5	A A HIR.	L = .TKUE. 6PWRH (RHO(NI)+ORDW(NI)) L +DRDH(NI)+(HOI-H(NI)) L +DRDH(NI)+(HOI-H(NI)) L / (KWO. DRDW(NI)+WPWI-DRDH(NI)) BPW(KI)=BPW(KI)+OW*(BPWB-RPW(KI)) BHO(KO)=HOI-HH+RPW(KI) BHO(KO)=BPW(KI)*WORD RETURN STOP 111	NAMER MAIN.OPT.02.LINECNT.55.SIZF#0000K.
(AUG 71)	COMPILER OPTIONS SUBRO CINSERT REF	NOMMOO		T IOO D D D D D D D D D D D D D D D D D D	പര്ഷ് ബ്ബ് പ	100	I EFFECT*
20.1	C 0000	40000 40000 70000	00000000000000000000000000000000000000	0021 0022 0022		0030 0032 0033 0035 0035	NI SN
LEVEL 3	N N N				1 S 1 S 1 S 1 S 1 S 1 S 1 S 1 S 1 S 1 S	N	*OPTIONS

MAIAD SAMPLE PROBLEM - COAPSE CALCULATION

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READ
A S
CATA
TUPNI

PES 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	V = -	VIDE 1 OUTLET WEADER THROUGH STEAM GEN TO COLN 3 TO 2 POWER -8.3D+8 SICTION 4*4 IMTRAN 4*0 P P P P P P P P P P P P P P P P P P P	TPE 2 MAKE-UP SYSTEM JCT THKOUG ************************************	3 9POKFN PIPF FROM OUTLET HEADER 0 1 1 4 4 0 0 0 0 0 1 0 .308 0 4 MAKE-UP SYSTEM 0 1 1 4 4 4*0 0 .01 -1 .1 .308	PE 5 S OES CO CTION 4 DZ *5 1 3*0 P	SECONDARY CONN & TO CONN & TO THEAN SECONDARY	MAECTION 1 1 1.05	* INITIAL CONDITIONS * FOLLOWING STEADY STATES DEMONSTRATE IMMEDIATE PRINT OPTION * THE STEADY STATES FOR THESE PIPES ARE REDEFINED LATER * THE STEADY STATES FOR THESE PIPES ARE REDEFINED LATER * THE STEADY STATES FOR THESE PIPES ARE REDEFINED LATER * THE STEADY STATES FOR THESE PIPES ARE REDEFINED LATER * THE STEADY STATES FOR THESE PIPES ARE REDEFINED LATER * THE STEADY STATES FOR THESE PIPES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE STEADY STATES FOR THE STATES ARE REDEFINED LATER * THE ST
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	TFC THT FLUX	544.03 0.0 0.0 535.73 0.0 -9.41043D+04 506.53 0.0 -9.41043D+04 506.49 0.0 0.0			TFC THT FLUX	366.25 0.0 0.0 400.00 0.0 0.0	
)	544.03 535.73 506.53 506.49			-	366.25 366.25	
80+000	RHO	768.80 782.62 834.22 833.90	3.905		Вно	962.98	
OLANT -8.300	ALPHA	-0.085423 -0.105925 -0.177021 -0.174394	THALPIES 243	POWER TO COOLANT 0.0	ALPHA	-0.002597	8.606 x IS
WER TO CO	×	-0.110836 -0.140000 -0.234647 -0.231742	.90+5 .EN	DWER TO CO	×	-0.006658 -0.006712	PRESSURE 0 4 SSURE 1 1 05
-8.300000+08 POWER TO COOLANT -8.300000+08	I	1.19009D+06 -0.110836 -0.085423 1.14862D+06 -0.146000 -0.105925 1.00336N+06 -0.234647 -0.177021 1.00332N+06 -0.231742 -0.174394	PRESSURE IS 0 .90+5 .ENTHALPIES 2+3.905		x	3,900000+05 -0,006658 -0,002597 3,900000+05 -0,006712 -0,002619	4 FLOW IS 0 4179 PRESSURE 0 8.606 X IS 014 2 0 1 1 4 1 4 1 4 1 4 1 5 1 1 1 1 1 1 1 1 1
POWEP GENERATED -8.	Q.	8.608990+06 8.500000+06 8.40000+06 8.332230+06		POWER GENERATED 0.0	Q.	8.990560+04 9.000000+04	1 AT NODE 4 FLO 1 AT NODE 2 0 1 AT NODE 1 FLO 5 SEGMENT 2 1 NODE 1 0 3000 4 CONNECTION 6 STATE
-	3	4.44427D+03 4.44427D+03 4.44427D+03 4.44427D+03	PIDE 4 OPTION -3 AT NODE 2 FUEL TEMPS 0 400.	4	3	0.0	1 OPTION 4 OPTION 2 OPTION 5 OPTION 6 OPTION 6 STEADY
IN PIPE	7	0 4 4 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6)o I d	IN PIPE	7	0.0	
	NODE	ലവനാ ഭ			NODE	14	·

*** OUTPUT ON PSEUDO FILE 1

CALCULATION
- COARSE
PROBLEM .
SAMPLE
NAIAD

PIPE NUMBER 1 JOINS CONNECTIONS 3 AND 2

CONN	CONNECTION	3 18	INTEPNAL									÷		
CORN	CUMMECTION	2 15	INTEHNAL											
MODE	2	0.	ů.	₫	8	Ī	c a	ROUGH	SLIP	FRIC	PH -	SEG	GSINT	ORIF
⊶ √ 1 m 4	0.0 4.00 18.00 22.00	0.3680 0.0100 0.0100	00 4 00 4 00 4 00 4 00 4 00 8	4.4700-01 1.2200+00 1.2200+00	0000	1.0000+10 1.0000+10 1.0000+10 1.0000+10	0.0 1.1340-04 1.1340-04 0.0	4.000+05 4.000+05 4.000+06 4.000+06	2005 2005 2005 2005 8	7885 7885 7885 7885	0000	NI M	9.8040+00 0.0 -4.9020+00	000
PIPE	NUMPER	~	JOINS CONNECTIONS	ณ	E UNE									
NOU	CONNECTION	2 15	INTERNAL											
CONN	CONNECTION	m	IS INTERNAL											
NODE	2	DE	ā	∢	ون	Ŧ	Od.	ROUGH	SLIP	FRIC	DHT	SEG	GSINT	ORIF
o c	0.0	0.3080	5.80 5.80	4.4700-0	00	1.0000+10		4.000-4	SNOC	SBBL	00		0.0	0.0
~ 30	11.50	0.0500	30.20		2.2100+04	1.0000+10	0.0	4.000.4	2000	SBBL	> ⊶ -	0 - a	-2.941D+02 -2.941D+00	50.00
20	26.00	0.0500	30.20	3.7400-0	2.2100+04 2.2100+04	1.0000+10		4.00D-05 4.00D-05	SNOS	SBBL SBBL SBBL	n and and		9.804D+00 9.804D+00	00
3d1 d	NUMBER	m	JOINS CONNECTIONS	L,	AND 4									
CONNE	CONNECTION	3 IS IN	IS INTERNAL											
CONNE	CONNECTION	4 WILL	PURST	OPEN OVER TI	TIME 1.00-02	Z TO PRESSURE		1.0000000000						
NODE	~	DE	đ	⋖	පී	¥	9	ROUGH	SLIP	FRIC	DH1	SEG	GSINT	914
=======================================	0.0	0.3080	5.80 0.04	4.470D-01 1.000D-04	0.0	1.0000+10	00	4.000-04 4.000-04	SNOC	SBBL SBBL	60	=		0.0
3010	PIPE NUMBER	4 JOINS	S CONNE	CONNECTIONS 1 A	AND 2									
CONNE	CONNECTION	I HAS G	GIVEN MASSFLOW	ď	DUNDARY CONDITION	NOI.	ŧ							
CONNE	CONNECTION	2 IS IN	IS INTEPNAL											
MODE	~	ÐĒ	ā. I	⋖	9	Ŧ	0	ROUGH	SLIP	FRIC	THG	SEG	GSINT	ORIF
13	0.0	0.1000	5.80	7.854D-03 4.470D-01	00.0	1.0000+10	00.0	4.000-05 4.000-05	SNOC	SBBL SHBL	0.3	13 -6	-9.804D+00	C . 0

PIPE NUMBER 5 JOINS CONNECTIONS 5 AND 6

CONNECTION 5 HAS GIVEN MASSFLOW SOUNDARY CONDITION

	ORIF	000				ORIF	000
	6 GSINT	4.902D+00 9.804D+00 4.902D+00				SEG GSINT	19 -4,9020+00 20 -9,8040+00 21 -4,9020+00
	546	15 16 17				S	200
	DH1	0000				PHT	0000
	FRIC DHT	1885 1885 1885 1885				FRIC DHT	588L 588L 588L 588L
	SLIP	SNOU SNOU SNOU SNOUS				SLIP	SNOUS
	POUGH	4.0000-04 4.0000-04 4.0000-06				ROUGH	4.0000-4 4.0000-4 4.0000-4
	ç	0.0 2.268D-04 2.268D-04 0.0				ç	0.0 2.268D-04 2.268D-04 0.0
	ij.	1.0000+10 1.0000+10 1.0000+10 1.0000+10			NOITION	L I	1.0000+10 1.0000+10 1.0000+10 1.0000+10
	a O	0000	6 AND 7		CONNECTION 7 HAS IMPLICIT MASSFLOW BOUNDARY CONDITION	a .	0000
	⋖	5.80 4.4700-01 496.00 3.0000+00 490.00 3.0000+00 5.80 4.4700-01			MASSFLOW	₫	5.80 4.470D=01 490.00 3.000D+00 490.00 3.000D+00 5.80 4.470D=01
LTEPNAL	ā	44 (1.004 (1.004 (1.000) (1.000)	UTPE NUMBER 6 JOINS CONNECTIONS	NTFFNAL	TMPL ICIT	đ	490.06 490.06 490.00
CONMECTION 6 IS INTERNAL	D.E	0.0 0.3080 7.00 0.0500 9.00 0.0500 11.00 0.3080	6 301	CONNECTION 6 IS INTERNAL	7 HAS	D.	0.0 0.3080 2.00 0.0500 9.00 0.0500 11.00 0.3080
NOTE	MMDE 7 DE	9.00	34417	NOLLO	NOLL	^	2.00 9.00 11.00
CONVEC	MODE	15 16 17 18	ر ت ت	CONNE	CONNE	None 7	70 H N

		FLUX	0.0 -9.41043D+04 -9.41043D+04 0.0	0.0 0.0 0.0 7.468560+05 0.0	00 00	0.0 1.882090+05 1.882090+05 0.0	0.0 -1.682090+05 -1.882090+05 0.0	
		THT	0000		00 00	0000	0000	
		TFC	573.54 565.75 535.74 535.74		573.54 573.52 535.73 535.73	527.00 539.36 542.77 541.64	541 542 527 527 50 527 50 50	
		-	573.54 565.75 535.74 535.73	2000 2000 2000 2000 2000 2000 2000 200	573.52 573.52 535.73	527.00 539.36 542.77 541.64	541.64 542.58 527.00 527.00	
38310+06	00000000000000000000000000000000000000	RHO	711.63 729.27 782.70 782.62	o o o u o o	711.63 711.60 782.62 782.62	794,77 774,24 282,05 276,95	276.95 335.88 794.76 794.74 ITY CHOKE	ज्युज्युज्युज्युज्युज्युज्युज्युज्युज्यु
COOL ANT -7.336	M	ALPHA	-0.005998 -0.032535 -0.107915	004640	-0.005998 -0.004082 -0.105922	-0.036372 -0.009084 0.656796 0.664060	0.654060 0.584200 -0.036377 -0.035744	-0.191751 -0.191750 -0.191750 -0.191856 -0.191876 -0.191816 -0.191876 -0.191876 -0.191929 -0.191754 -0.191754 -0.191754
OWFR TO COO	OWER TO COOL OWER TO COOL OWER TO COOL OWER TO COOL OWER TO COOL	×	-0.007265 -0.038750 -0.142468 -0.140000	-0.140000 -0.139775 -0.191751 -0.028483 -0.0013195	-0.007265 -0.004943 -0.139997	_	0.126033 0.085467 -0.050006 -0.049160 ET FDR FLO	80.000 90.0000 90.000 90.000 90.000 90.000 90.000 90.000 90.000 90.0000 90.000 90.000 90.000 90.000 90.000 90.000 90.000 90.0000 90.
33831D+04 PC	0000+008 6620+008 0000+009 0000+008	r	.347240+06 .303140+06 .148660+06 .148620+06	.14862h+06 .14862h+06 .15049b+06 .34728b+06 .34728b+06	.347240+06 .347140+06 .148620+06 .148620+06	.104640+06 .166140+06 .381250+06 .380890+06	380890+06 319770+06 104660+06 104640+06 ITY INL	
ATED -7.	1750 -8.30 1750 8.22 1750 0.0 1750 8.30	a	85217D+06 1 85217D+06 1 66381D+06 1 60000D+06 1	.59426D+06 .99426D+06 .91841D+06 .93838D+06 .85947D+06	5947D+06 1 8729D+06 1 9991D+06 1	00000+06 1.7780+06 1.73630+06 1.75610+06 1.	7661D+06 1, 5727D+06 1, 0057D+06 1, 7854D+06 1, URE AND QUAL	-4.9428340-03 3.291380-03 3.2991380-03 4.2384250-03 8.7874120-03 1.9959330-02 2.8479440-02 2.847940-02 3.7203460-02 4.326460-02 4.326460-02 5.6735570-02 5.673550-02 5.7729920-02
POWES GENER	POWER GENERAL POWER	3	9000+03 R. 9000+03 R. 9000+03 R.	## ## ## ## ## ## ## ## ## ## ## ## ##	0000+00 8.8 0000+00 8.7 0000+00 8.5	00000+03 F	0000+03 5.2 0000+03 5.4 0000+03 5.4 1HPOAT PRESS	R.787295D+06 8.526476D+06 8.526676D+06 8.426086D+06 8.196903D+06 6.932976D+06 6.932936D+06 6.932936D+06 6.932936D+06 6.932936D+06 6.932936D+06 6.87291000+06 6.745612D+06 6.745426D+06
0.0	MAN	2	0.0 4.17 4.00 4.17 18.00 4.17 22.00 4.17	44444	0.0 1.00 0.01 1.00 0.01 1.00	00.60	11.000 W.000	00000000000000000000000000000000000000
TIME		NOOF	HNM4	ט ער פט אייטי	13 12 14		15 20 21 22 1	0.0 1.00001 7.00000 1.50000 6.50000 6.0016 6.0016 6.0016 6.11779 6.11399 6.50000 6.11399

		FLUX	0.0 -9.29389D+04 -9.29389D+04 0.0	0.0	3146D+0	5.359270+05 4.207900+04 1.625340+04			1.858780+05		0.0 -1.862090+05 -1.882090+05 0.0						÷						
		Ħ	0000	• •		22.00			000	• •	0000												
		1FC	572.49 565.54 535.65 535.63	40.4	Š	573,35 573,29 573,43	* r-	35.7	527.00	45.1 41.6	542,57 542,57 527,00 527,00												
		-	572.49 565.54 535.65 535.65	9	2	570.73 572.67 572.49		35.7	527.00	42.7 41.6	541.63 542.57 527.00 527.00	6											
⊢ ⊢ ⊢	690+08 210+08 350+08 210+08	ž ž	688.07 728.10 782.38 782.35		່ຜ	718.55	0 0	82.2	794.77	નું ભ	277.00 335.84 794.76	TY CHOKE											
-0,190121 -0,182685 -0,179385	ANT -2.311 ANT -8.197 ANT 0.0 ANT 0.0 ANT 0.0	•	0.038179 -0.022420 -0.097183 -0.095243	42560	13484	-0.023911 -0.004392	03817	09500	-0.036332	.65868 .66399	0.663995 0.584266 -0.036359 -0.035725	OW AND QUALI	-0.177245	-0.151086	-0.131677	-0.100625	-8.5410330-	-6.2459100-	-5.403357D-	-5.2184330-	14.7012940	-4.597220D-	-4.6530H7D-
85.27 02.10 08.74	POWER TO COO POWER TO COO POWER TO COO POWER TO COO	0 0 ×	6 -0.023110 6 -0.027214 6 -0.129049 6 -0.126593	6 -0.12659	6 -0.12635 6 -0.17724	6 -0.028443	6 0.00311	-0.12625	00	0.12309	6 0.125982 6 0.085484 6 -0.049982 6 -0.049135	INLET FOR FL(211.07	224.70	223.41	247.71	239.71	218,35	208.36	209.14	2 6	90	58.
180-02 41 240-02 42 340-02 42	.573100+08 P	#0+0E00	1.345760+00 1.302550+00 1.148280+00	.14818D+0	.14817D+0 .15013D+0	1.330990+0	345760+0	.338100+0 .14862D+0 .14827D+0	.1046	1,38104D+0 1,380790+0	1.380790+06 1.319760+06 1.10466n+06 1.104630+06	QUALITY I	6.6743030-02 4 6.0664320-02 4	•	•	•	•	• •	•	4		•	•
5.8261 6.2371 6.3896	ATED -8. ATED 3. ATED 0. ATED 0.	PATED +8	.513530+06 .504640+06 .312580+06 .247370+06	247370+0	24097D+0 56021D+0	13917	513530+0	.38476540 .246620+0 .247375+0	.498630+0	.47238D+06	.376200+06 .456550+06 .499910+04 .477880+06	SSURE AND				_			_				
6,720588D+06 6,532791b+06 6,452606D+06	DOWER GE 3 POWER GE 4 POWER GE 5 POWER GE	S POWER GE	4.144610+03 B. 4.186270+03 B. 4.197030+03 B.	782D+03	19540+03	4.17155D+03 9.	34290+03	.987330+00 .288530+00 .290520+00	.000000+03	2.992250+03 5 2.99021D+03 5	2.990210+03 5 2.990110+03 5 2.999980+03 5 3.000000+03 5	THROAT PRE	6.3847580+	5.8070170+	5.4247320+	4.9172470+	4.666012D+	4.426938D+	3.9647210+	3.9054720+06	3.816787U4 2 474531D4	3,4580610+	3,230186D4
.553746 .660441 .698255	E84E NI	12 N	0.0 4.00 18.00	0.0	2,50	16.50	0.00		0 1	11:00	11.00	TIME	.77348	1.2248	5415	3009	7572	4050		4.81486	1597	1448	9019
000	T I E	NOOF	⊸ 01.00 4	ŝ	40	· ac or	2 :	2 51	21.	17.1	2002		9 (•									

	<u>.</u>	•	-5.2426BD+0	-5.24268D	• •	0.0	6.022	40.40	6.25352D+0	0 0	• -	000	0.0		0.0	0.0	.89979D+0	•					
	<u>;</u>			000	0 0	0	0.		42.00	0 0		0	0 0	9	0.0		000						
	TFC	ָ ע ע	14.00 14.00	490.87	49.64	499.27	505.30	515.83	506.28	505.20 489.85	E .00		488.83 497.93	97.0	Ň	464.24	486.21 488.83						
	-	, A	05.3	499.86 499.52	99.5	~	100 t	· ~	_	505.11	2	66	488.83	96.9	_	ውው	488.20 488.82	c					
210-02 470-02 110-02 110-02 1382810+08 624040+08 624040+08	AC+08	516.76	623,22	831.71 728.47	716.94	663.76		816.96		516.76	4	•	846.12 709.20	σ	4	247.49		CHOKE					
1-4-77 1-5-97536 1-5	OLAN: -6.	0.38007	0.24841	f	0.140431	20002.0	-0.0151	-0.00010	0.38007	0.380078	.721	71975	-0.013753 0.152287	•6956	.7151	71515	55	W AND QUALITY	-5,3596730+02	-5.7360220-02	-5.9323790-02	-6.1031800-02 -6.3074360-02	-6.5002680-02
220 2210 2214 2214 2214 2214 2014 2014 2	- ×	0.017	9000	90	0.0		220.0- 3	-0.000	5 0.017	6 0.017441 6 0.053537	6 0.096592	61960-0	90	0.079	0.085		-0.0199	INLET FOR FLOW	4264.07				
090-02 950-02 950-02 990-02 970-02 970-02 970-02 902230+08 624040+08 624040+08	I			9.791670+05	9.799680+08	364700+0	20050+0	26330+0	30920+0	1-030920+00	1.148620+06	•14915D+0	9.238910+05 9.726190+05	17350+0	•10/250+0	1.107260+06 1.057310+06		QUALITY IN	37380-02 42 02010-02 42		ر 1000	ָטָ פָּי	20-
A TEO C TA C C C C C C C C C C C C C C C C C	5	90+06+669	.91088D+06 .68737D+06	•61647D+06	2.61647D+06	011580+04	62881D+06	212410+06	899490+06	.89949D+06	90+050+09	16470+0	2.51748D+06 2.53557D+06	493220+06	90+761+96	2.36415D+06 2.46669D+06 5.510145+04	+R665D+06	SURE AND GL	5,35	5.19	5.0		4.70
0.000	3	•••	4.24129D+03 2.	1 (1)	.25151D+03	.26407D+03	*25597D+03	*246320+03	• < 3281D+03	.001900+00 2 .012430+00 2	.000000+01 2.	00385D+01 Z.	3.193870+03 2.3.180580+03 2.3.	4536D+03		3.211680+03 2.3 3.227790+03 2.4 3.205550+03 2.4		THPOAT PRESS	2.1731900+06	1.9269910+00	1.8651420+0(1.626248D+06	1.532632D+0f
10.118 10.118 10.064 10.064 10.706 11.706		0	0 0	2.00	6 V	9 09	4	0 0	2	0.0	0	7 10.	0000			M M M	0	TIME	17.0485 18.2822	.5077	. A.C.A.	9416	.0390
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TIME 2.500000+01 POWER GENERATED -7.810700+08 POWER TO COOLANT -7.33R195+08
IN PIPE 1 POWER GENERATED -4.532335+08 POWER TO COOLANT -4.532335+08
IN PIPE 2 POWER GENERATED 5.75863D+07 POWER TO COOLANT 1.04838D+08

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60.695

TOTAL STEAM IN CONTAINMENT

	THT FLUX	0.0 0.0	0.0 -5.138700+04		0.0 0.0	0.0		1.00 6.223790+04				42.00 5.91051D+04	0.0 0.0	0.0		0.0 0.0	0.0 0.0			0.0 1.027740+05	0.0 1.027740+05	0.0	0.0 0.0		0.0 -1.901720+05			
	TFC	483.80	484.24	477.96	475.12	475-12	474.26	4 K 3. 35	488.81	40.004	1000	484.79	483.80	470.20	; ; ;	474.32	475.12	1	467.14	476.34	475.04	469.67	19.69*	473.74	466.50	467.14		
	►	483.75	484.10	477.95	475.05	475.05	474.18	402.30	70.7	100 00 7	70.40	483.75	483.75	470.16	•	474.25	475.05		467.14	476.28	474.98	469.59	469.59	473.68	466.50	467.14	۵	
0.0 0.0 4.53233D+08 -8.38657D+08	ОНа	483.21	65.064	849.09	564.45	561.92	502 46	00 900	44	70 440	40.0	483,21	483,21	280.78	•	153.05	155,00		A71.76	707.90	231,15	204.53	204.53	275.48	872.47	871.68	ту снокер	-02 T
	ALPHA	0.437888	0.273894	080000	0.348838	F 13636 A	400000	244	0911000	004120*0-	*0.00484	0.437888	0.437888	CAECRA.0		0.830473	0,824133		-0.011847	0.179575	0.739011	0.771069	0.771069	0.687505	-0.011583	-0.009656	FLOW AND GUALITY	-6.6927500-02
POWER TO COOLANT POWER TO COOLANT POWER TO COOLANT	×	0.015678	004150			040000	100110					0.015678	0.015678	0.045882		0.150598	0.149270		-0.02000-			0.084625	0.084625	0.050617	-0.019617	-0.016407	INLET FOR FLO	4271.75
32330+08 96570+08	τ	9,307110405	0.152488405	24080040 a	8.77186D+05	. 0	20.000000000000000000000000000000000000	000000000000000000000000000000000000000	C24011046*0	50+02456746	49+025192462	9+302110+05	9.302110+05	204070405	CA+G:0:03*6	1.148620+06	1.149210+06		8,255060+05	8,726020+05	1.002190+06	1,001580+06	1,001580+06	9.527620+05	8,226300+05	8.255060+05	QUALITY INL	
GENERATED 0.0 GENERATED 0.0 GENERATED 4.5: GENERATED 4.5:	a .	404050404	1 04 40 00 40 1	1.741900+06	1.617810+06		_		_	2.696040406	2,290220+06	1,930950+06	1.930950+06	1 460000	90406000	1.590980+06	1.61781D+06		ø	•		1.442480+06	442480+06	.57220D+06	4614590+04	1.590400+06	AND	106 4.5882240-02
3 POWER GE 5 POWER GE 6 POWER GE	>	F0+083465.4											1.556210+00			2.000000+01						3.27149D+03	E0+0671227E			3,290490+03	THROAT PRESSURE	1.460095D+06
TN PIPE TN PIPE TN PIPE	2	•									21,50		6.0	_		0.0			0.0	0			0.0	_			TIME	25.0000
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NAIAD SAMPLE PROBLEM - ACCURATE CALCULATION

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7 CONNECTIONS CRAVITY 9.8043 EPPING CONTROLS TS .0001 40. CHANGE TARGET .01 CREASE 2 TMAX 25 FULL PRINT EACH	* PIPE 1 OUTLET WEADER THROUGH STFAM GEN TO MAKE-UP SYSTEM JUNCTION 4 ***ODES CONN 3 TO ? POWER -8.3D+8 SLIP 4*1 FRICTION 4*4 IHTRAN 4*0 PRINT 0 0 0 Z 0 4 18 22 DH/DZ 1 05 DE .308 Z***OI .308 A .447 Z*1.22 .447 HP 5.8 Z*490 5.8 ROUGH 4.D-5 Z*4.D-6 4.D-5 CP 4*0	PIPE 2 NIOTIEN 2 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	PE 3 AFOKEN PIPE FF 4 0 1 1 4 4 0 0 0 0 PE 4 MAKE-UP SYSTEP 2 0 1 1 4 4 4*0 0	PIPE S SECONDARY LOOP PUMP THROUGH STEAM GEN NODES COND S TO 6 POWER 8,3D+8 SLIP 4*1 FRICTION 4*4 IHTRAN 4*0 PRINT 4*0 Z 0 Z 9 11 DH/DZ 85 1 *5 DE 308 Z**,05 308 A 447 Z*3 *447 P 5.8 Z*490 5.8 ROUGH 4*D-4 Z*4,0D-6 4*D-4 CP 4*0 HF 4*0 RIK 3*0 PD 0 1 1 0	* PIPE 6 SECONDAFY LOOP HEAT SINK TO PUMP 4 NODES CONN 6 TO 7 POWER -8.3D+8 SLIP 4*1 FPICT 4*4 IHTPAN 4*0 PRINT 4*0 7 0 2 9 11 DH/DZ5 -15 DE .300 2*.05 .308 A .447 2*3447 HP 5.8 2*490 5.8 POUGH 4.D-4 2*4.D-6 4.D-4 CP 4*0 HF 4*0		* INITIAL CONDITIONS PIPE 1 OPTION 1 AT NODE 4 FLOW IS 0 4179 PRESSURE 0 8.6D6 X IS 014 PIPE 4 OPTION 1 AT NODE 2 0 1 1 4 1 4 PIPE 3 OPTION 1 AT NODE 1 FLOW IS 4 1 PRESSURE 1 1 X IS 1 1 PIPE 2 OPTION 5 SEGMENT 2 PIPE 5 OPTION 1 NODE 1 0 3000 0 5.5D6 005
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PIPE 6 OPTION 4 CONNECTION 6 0 END STEADY STATE *

(ADD TO DATA FILE) OUTPUT CONTROL 0

*** OUTPUT ON PSEUDO FILE

15 DEC 75

PIPE NUMBER 1 JOINS CONNECTIONS 3 AND 2

CONNECTION 3 IS INTERNAL

CONNECTION	\$1 2 N	INTERNAL											
NODE 7	90	Ţ	∢	ð	Ť	Q	ROUGH	SLIP	FRIC	D#1	SEG	651NT	ORIF
1 0.0 2 4.00 3 18.00 4 22.00	0.3080 10 0.3100 10 0.0100	6.00 00 00 00 00 00 00 00 00	4.470D-01 1.220D+00 1.220D+00	0000	1.0000+10 1.0000+10 1.0000+10 1.0000+10	0.0 1.1340-04 1.1340-04	4.000-05 4.000-06 4.000-06 4.000-05	JONS JONS JONS JONS	3881 3881 3881 3881	0000	N W	9.804D+00 0.0 -4.902D+00	000
PIPE NUMPER	~	JOINS CONNE	CONNECTIONS 2	E ONA									
CONVECTION	2 15	INTERNAL											
CONNECTION	3 IS	INTERNAL											
2 300N	DE	Ţ	4	9	1	o d	ROUGH	SLIP	FRIC	DH1	SEG	GSINT	ORIF
5 0.0			4.4705-01	0.0	0.0000-1		4						•
	0 0,3080		4-4700-01	0.0	1.0000+10	00	4.000-05 4.000-05	CONS	888L 888L	0	o o	0.0 2.115D+02	000
8 16.50	0.0010	~	4.9100-01		1.0000+10	0.0	4.000-05	SNOC	SBBL	-		-2.941D+00	50.00
		(4)	3.7400		1.0000+10	0.0	4.000106	SNOC	588L			Z.941D+00	0.0
26.0	0 0.0500		3.740D-01		1.0000+10	0.0	4.00D-05	SNOS	SBBL	-		00+0+00	0.0
PIPE NUMBER	m	JOINS CONNECTIONS	CT TONS 3	AND 4									
CONNECTION	m	IS INTERNAL											
CONNECTION	*	WILL BURST OPEN OVER	-	IME 1.00-02	TO PRESSURE		1.000000+05						
NODE 7	DE	<u>q</u>	∢	g.	IA.	6	ROUGH	SLIP	FRICD	DHT	SEG	6SINT	ORIF
11 0.0 12 0.01	0.3080	00.0	4.4700-01 1.0000-04	000	1.0000+10	00	4.000-04 4.000-04	CONS	SBBL	••	11 0	0.0	0.0
PIPE NUMBER		4 JOINS CONNECTIONS	-	AND 2									
CONNECTION	-	HAS GIVEN MASSFLOW	900	NDARY CONDITION	ION								
CONNECTION	~	IS INTERNAL											
NODE: 7	90	ā	∢.	ą,	¥	04	ROUGH	SLIP	FRIC O	DHT	SEG	GSINT	900
13 0.0	0.1000	5.80	7.8540-03	0.0	1.0000+10	00.0	4.00D-05 4.00D-05			90	139	-9.8U4D+00	0.0

PIPE NUMBER 5 JOINS CONNECTIONS 5 AND 6

CONNECTION 5 HAS GIVEN MASSFLOW BOUNDARY CONDITION

CONNECTION 6 IS INTERNAL

		FLUX	0.0 -9.4]043D+04 -9.4]043D+04 0.0	0.0 0.0 0.0 7.466560+05 0.0	00 00		0.0 -1.88209D+05 -1.88209D+05 0.0	
		THT	9990	0.0 0.0 1.00 15.00 1.00	00 00		0000	
		TFC	573,54 565,75 535,74 535,74	536.43 536.33 536.33 573.57 573.57	573.54 573.52 535.73 535.73	27.0 39.3 41.6	541.64 542.58 527.00 527.00	
		-	573.54 565.75 535.74 535.73	53 33 33 33 34 35 34 35 35 35 35 35 35 35 35 35 35 35 35 35	573.54 573.52 535.73 535.73	27.04 4.04 4.04 4.1.64	541.64 542.58 527.00 527.00	
33831D+06	000+08 620+08 000+08 000+08	8H0	711.63 729.27 782.70 782.62	782.62 784.62 712.57 711.82	7111.63 7111.60 782.62	44NA 100	276.95 335.88 794.76 794.74	मा ज्याच व्याच व्याच व्याच व्याच व्याच व्याच व्याच व्याच मा
-	-8.300 0.0 0.0 0.0 8.300	ALPHA	0.005998 0.032535 0.107915 0.105925	0.105925 0.105744 0.149330 0.024196 0.010958	0.005099 0.004082 0.105922 0.105925	0363 0000 6567	.664060 .584200 .035377 .035744	-0.191749 -0.191747 -0.191747 -0.191547 -0.191538 -0.191538 -0.191656 -0.191920 -0.192056 -0.192056 -0.192249 -0.192249
ER TO COOLANT	7000 7000 7000 7000 7000 7000 7000 700	×	0.007265 - 0.038750 - 0.142468 -	0.139775 -0 0.139775 -0 0.191751 -0 0.028483 -0 0.013195 -0	0.007265 -0 0.004943 -0 0.139997 -0 0.140000 -0	.050000 .012340 .123181	1.126033 0 1.085467 0 1.050006 -0 1.049160 -0	
D+06 POWE	POWER	1	17240+06 -0 13140+06 -0 18660+06 -0	96520+06 -0 96520+06 -0 90490+06 -0 7350+06 -0 7280+06 -0	7240+06 -0 7140+06 -0 8620+06 -0 8620+06 -0	1 1	380895-06 0 319775-06 0 104665-06 0 104645-06 0	
-7,338310+06	8.2266204 0.0 0.0 0.0 8.3000004				5 1.34 5 1.34 7 1.14 5 1.14	AAAA	HHHH NO	1428340-0 2752600-0 2752600-0 2752600-0 275220-0 143110-0 275280-0 270870-0 270870-0 270870-0 270870-0 270870-0
GENERATED	GENERATED GENERATED GENERATED GENERATED GENERATED	۵	8.85947D+06 8.85217D+06 8.66381D+06 8.600000+06	8.600000+006 8.546260+00 9.918410+00 9.492750+00 9.038380+00	8.85947D+00 8.787290+00 8.599910+00 8.600000+00		5.376610+06 5.457270+06 5.500570+06 5.478540+06 FF SSURE AND	
a a a a a a a a a a a a a a a a a a a	200 d d d d d d d d d d d d d d d d d d	3	4.179000+03 4.179000+03 4.179000+03	180000+03 180000+03 180000+03 180000+03 180000+03	00+000000 00+000000 00+000000	00000+03 00000+03 00000+03	00000+03 00000+03 00000+03 1440AT P	8.7872950+0 8.6125490+0 8.516490+0 8.5370390+0 8.6283740+0 8.066790+0 8.066790+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0 7.9788770+0
0.0		7	000000000000000000000000000000000000000	7.50 4.11 7.50 4.11 7.50 4.11 7.11.50 4.11 7.11.50 4.11	00 00 00 00 00 00 00 00 00 00 00 00 00		0.00 3.00 0.00 3.00 0.00 3.00 11.00 3.00	0.0 1.000000000000000000000000000000000
1 CME		MODE	~~~*	46 t 2 46 U	12 24	1 0 - 3	19 20 21 22 32 71	

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			FLUX	0.0 -5.165930+04 -5.165930+04 0.0	0.0 0.0 6.07423D+04 0 7.11511D+04 0 5.94106D+04 0 5.87025D+04	00 00	0.0 1.033190+U5 1.033190+05 0.0	0.0 -1.899300+05 -1.899300+05 0.0
			71	0000	001114	00 00	0000	9000
			TFC	485.74 486.20 479.34	477.92 477.07 484.63 489.96 491.40	485.74 472.33 477.20 477.92	467.82 476.12 476.79 471.86	471.86 475.56 467.19 467.82
e .			-	485.74 486.20 479.34 477.92	417.92 477.96 483.64 490.94 490.94	485.74 472.33 477.20 477.92	467.82 478.12 476.79 471.86	471.86 475.58 467.19 667.82
¥ CHOKE	002 002 002 002 002 002 002 002 003 003	50+08 40+08 50+08	RHO	511,24 665,17 857,43 647,19	640. 853.14 852.14 846.07 811.85	511,29 292,64 161,09 162,92	871.06 795.19 237.45	212.37 286.79 871.77 470.77
W AND GUALIT	-6.3968880- -6.3968880- -6.38583930- -6.4081710- -6.4687770- -6.46880- -6.57169290- -6.5716920- -6.5716920- -6.5716920- -6.7126380- -6.7126380- -6.7126380- -6.81726380- -6.7126380- -6.31726380- -6.31726380- -6.31726380-	ANT 5543 ANT 1.0375 ANT 0.0 ANT 0.0 ANT 4.5563 ANT -9.3759	ALPHA	0.403015 0.219221 -0.001520 0.249159	0.269157 0.348706 -0.042976 -0.022314 -0.005404	0.403015 0.667907 0.820820 0.818608	-0.013613 0.074873 0.731216 0.761550	0.761550 0.673496 -0.013413 -0.111444
NLET FOR FLO	00000000000000000000000000000000000000	POWER TO COOL POWER TO COOL POWER TO COOL POWER TO COOL POWER TO COOL	×	5 0.013676 5 0.005040 5 -0.002482 5 0.005141	0.005779 0.008748 -0.058234 -0.034963 -0.008626	0.013676 0.043639 0.144566	0.022832 0.001150 0.071618	0.081864 0.048046 -0.022552 -0.019393
UALITY IN	C C C C C C C C C C C C C C C C C C C	556350+08 F 758630+07 F 0 0 0 556350+08 F 375430+08 F	1	9.353770+09 9.211750+09 8.804060+09 8.837940+09	8.850200+05 8.8688600+05 9.2647400+05 9.312780+05	9.353770+05 9.33425n+05 1.14862D+06 1.14916D+06	8.285627.05 8.770747.05 1.006619.06 1.005690.06	1.005690+06 9.560330+05 8.257090+05 4.285620+05
SSUPE AND G	00000000000000000000000000000000000000	KATED -4. PATED 5. PATED 0. PATED 0.	Q.	.00847D+06 .0268BD+06 .80462D+06 .71619D+06	*716190+06 *685450+06 *162340+06 *775090+06 *355450+06	.00847D+06 .52863D+06 .691150+06	.70450D+06 .72351D+06 .47685D+06 .51364D+06	.513440+06 .635440+05 .674100+06 .654950+06
THROAT PRE	1.713580 7.725260 7.725260 7.725260 7.726830 7.66849790 7.66849790 7.66849790 7.66849790 7.66849790 7.66849790 7.66849790 7.66849790 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.6684990 7.668490 7.668	1 PONEER GENER 2 PONEER GENER 3 PONEER GENER 5 PONER GENER 5 PONER GENER 6 PONER GENER	3	033 A	334020+03 1 354140+03 1 354140+03 2 350860+03 2 243490+03 2	.59672D+00 2. .60312D+00 1. .00000P+01 1.	3.287760+03 1. 3.274910+03 1. 3.222480+03 1. 3.246840+03 1.	3.246860+03 1. 3.271270+03 1. 3.298347+03 1. 3.24745+03 1.
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CHUKED	L-	
INLET FOR FLOW AND QUALITY CHUKED	4358.14 -6.823383D-02	
HROAT PRESSURF AND QUALITY	1.528626D+06 4.363929D+02	61.661 KG
THROAT PRESSUR	1.5286280+06	CONTAINMENT
14 A T T	0000*52	TOTAL STEAM IN CONTAINMENT

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* INPUT FILE STRUCTURE IS:-
                                                    POSITION(S)
          VARTABLE
              TIME
                                                           2-23
              FLOWS
                                                          24-45
           PRESSURES
                                                          46-67
           QUALITIES
                                                          68-73
            P()WEHS
                                                             74
75
           TOT POWER
           TIME STEP
     DATA INPUT PHASE FOR COARSE CALCULATION
 FILE 1 75 ENTRIES
GET 2 1 13
* PLOT PHASE
PAPE 0 -8 -6 1 1 -1 2
LEGO 4 5
LEG 2 "COARSE CALCH."
HEAD "FLOW AT BREAK"
NAMX "TIME (SECONDS)"
NAMY "FLOW (KG/S)"
* RETURN TO INPUT PHASE TO ACCESS OUTPUT OF FINE CALCN
 DATA
 FILE 2 75
  4 1 13 35 75
* THIS GUES ON PREVIOUS PLOT
LINE 1 2 2
LEG 2 "FINE CALCH."
* NOW COMMENCE NEW PLOT
PAPE 8 -6 1.F-4 25 0 10
* SCALE PASCALS DOWN TO MPA
* SCALE PASCALS DOWN TO MPA
CONS 0 1.E-6
LINE 1 3 1
HEAD "PHESSURE AT HREAK (FINE CALCULATION)"
NAMX "TIME (SECONDS)"
NAMY "PRESSURE (MFGAPASCALS)"
PAPE 8 6 1.E-4 25 1.E-5 1
LINE 1 4 1
HEAD "TIME STEP (FINE CALCULATION)"
NAMX "TIME (SECONDS)"
NAMY "TIME STEP (SECONDS)"
 END OF FILE ON UNIT 1
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EXECUTION TEHMINATED

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