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**OWEN-1 - A CODE FOR ANALYSIS OF TWO-PHASE COMPRESSIBLE
FLOW TRANSIENTS IN A SINGLE CHANNEL**

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W. J. TURNER

May 1975

ISBN 642 99688 1

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ABSTRACT

A computer code for simulating two-phase compressible flow transients including loss-of-coolant in a single water cooled channel is described. The momentum flux term is included and slip between the phases is allowed. A pipe model and detailed surface heat transfer is included. The conservation equations are solved by means of an implicit numerically stable finite difference scheme so that time step size can be chosen on accuracy considerations alone.

National Library of Australia Card number and ISBN 642 99688 1

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BOUNDARY CONDITIONS; EQUATIONS OF STATE; FINITE DIFFERENCE METHOD; HEAT TRANSFER; LOSS OF COOLANT; MATHEMATICAL MODELS; O CODES; PIPES; REACTOR CHANNELS; SIMULATION; TRANSIENTS; TWO-PHASE FLOW; WATER COOLED REACTORS

CONTENTS

	Page
1. INTRODUCTION	1
2. CONSERVATION EQUATIONS	2
3. FINITE DIFFERENCE EQUATIONS	3
4. EQUATION OF STATE	6
5. SLIP MODEL	6
5.1 No Slip and Single Phase	6
5.2 Jones' Slip	7
6. FRICTION	8
7. PIPE MODEL	9
8. HEAT TRANSFER	9
9. BOUNDARY CONDITIONS	10
9.1 Moody Model	10
9.2 Rodflow Model	11
9.3 Fauske Model	11
10. SOLUTION OF FINITE DIFFERENCE EQUATIONS	12
11. INPUT DATA	13
11.1 Title Card	13
11.2 Control Array	13
11.3 Boundary Condition Table	16
11.4 Geometry Table	16
11.5 Restart Table	17
11.6 Output Table	17
12. COMPUTING TIME	17
13. OUTPUT	17
13.1 Error Messages	18
14. PLOTTING	19
15. ACKNOWLEDGEMENTS	20
16. REFERENCES	20

Table 1 Orifice Friction Specific Volume v_o

(continued)

CONTENTS (continued)

Table 2	Pipe Friction Specific Volume v_p
Table 3	Correlations for Surface Heat Transfer Coefficient
Table 4	Rodflow Choked Flow Model
Table 5	REAL*8 Arrays in Common Area DIM
Table 6	Equation of State Mesh
Table 7	REAL*8 Arrays in Common Area DATA
Appendix A	Input Data for Sample Problem
Appendix B	Output from Sample Source
Appendix C	Notation

1. INTRODUCTION

Many finite difference digital computer programs have been written to analyse two-phase compressible flow transients in water cooled nuclear power reactors. Most of these (Elliott 1968, Redfield *et al.* 1967, Rettig *et al.* 1970) are limited to time steps of order of the time taken for a sound wave to cross one space segment. This is also true of most programs written to analyse the loss-of-coolant accident in sodium cooled fast reactors (Sha & Hughes 1970, Carter *et al.* 1970). In many transients of interest, conditions are changing sufficiently slowly that consideration of accuracy alone would allow much larger time steps.

This problem is overcome in the family of codes, HYDNA (Currin *et al.* 1961), POISE (Carver 1967) and SLIP (Moxon 1968), derived from the momentum integral method (Meyer 1961). In this method dependence of steam and water densities on pressure is neglected thus eliminating sonic effects and permitting larger time steps. The codes are suitable only for transients in which small pressure changes occur; even then neglect of compressibility is an uncertainty in the results. The method is not suitable for loss-of-coolant analysis where larger pressure changes occur. Thus a finite difference scheme is desirable which includes compressibility and is stable for large time steps.

Such a scheme has been given by Turner & Wendroff (1964) for a perfect gas. In FLASH-4 (Porshing *et al.* 1969) a partially implicit method is used which is described by Porshing *et al.* (1971), but in practice this scheme is only conditionally stable.

The code described in this report uses a fully implicit scheme for two-phase compressible flow with slip between the steam and water phases. It is a further development of the method of Turner (1969) and includes the spatial acceleration or momentum flux term. The scheme has been found stable for time steps of any size. Running times are comparable with the momentum integral codes mentioned above where both methods are applicable.

Input data required are geometric and thermohydraulic descriptions of the channel from inlet to outlet, parameters to determine space and time step sizes, boundary conditions at the channel inlet and outlet, and specification of the initial state of the system. Usually this initial state is a steady state calculated from input data giving pressure, thermodynamic quality and flow rate at a specified point in the channel. However, a restart option is available in which the initial state at every node is taken from input data.

The calculation consists mainly of computation of pressure, quality, pipe temperature, etc., at each space node for each time step. Four types of transients can be computed, depending on the three boundary conditions specified. These are inlet enthalpy; inlet pressure or inlet flow; and the boundary condition at the outlet, either outlet pressure or one of the choked flow conditions. Three such conditions between pressure, enthalpy and flow at the outlet are incorporated in the code.

The code is so written that the equations of state for other coolants are easily added in place of those for light water. Similarly additional heat transfer relations can be included.

The principal limitation of the code is that flow reversal must not occur at either the channel inlet or outlet. Reverse flow within the channel is allowed.

2. CONSERVATION EQUATIONS

Consider a channel or pipe of cross section area a , which is a function of Z , the distance along the channel. The conservation equations are (Meyer 1960):

$$\text{mass} \quad a \frac{\partial \rho}{\partial t} + \frac{\partial W}{\partial Z} = 0$$

$$\text{momentum} \quad \frac{1}{a} \frac{\partial W}{\partial t} + \frac{1}{a} \frac{\partial}{\partial Z} \left(\frac{M}{a} \right) + \frac{\partial P}{\partial Z} + F + \rho g \frac{dy}{dZ} = 0$$

$$\text{energy} \quad \frac{\partial}{\partial t} \left[a \rho H - a P + \frac{M}{2a} \right] + \frac{\partial}{\partial Z} (E + aK) + Wg \frac{dy}{dZ} = a \phi$$

where the pressure P is considered constant over each plane normal to the Z axis; h , d and u vary across this plane; and averages over the plane are

$$\text{density} \quad \rho = \frac{1}{a} \int_a d \, da \quad ,$$

$$\text{static enthalpy} \quad H = \frac{1}{\rho a} \int_a h \, d \, da \quad ,$$

$$\text{mass flowrate} \quad W = \int_a u \, d \, da \quad ,$$

$$\text{momentum flowrate} \quad \frac{M}{a} = \int_a u^2 \, d \, da \quad ,$$

$$\text{energy flow} \quad E = \int_a u \, h \, d \, da \quad , \quad \text{and}$$

$$\text{kinetic energy flux} \quad K = \frac{1}{2a} \int_a d \, u^3 \, da \quad .$$

Flow enthalpy, as distinct from static enthalpy H given above, is E/W .

We assume thermodynamic equilibrium and take the equation of state in the form (Meyer 1961)

$$\rho = \rho(P, H)$$

$$\text{then } \frac{\partial \rho}{\partial t} = R_P \frac{\partial P}{\partial t} + R_H \frac{\partial H}{\partial t} .$$

We assume that M and E are not explicit functions of the channel geometry or surface heat flux, i.e. given W, P and H the quantities M and E can be calculated, excluding such effects as subcooled boiling and diameter effect on slip ratio. Then the conservation equations may be written

$$\text{mass } aR_P \frac{\partial P}{\partial t} + aR_H \frac{\partial H}{\partial t} + \frac{\partial W}{\partial Z} = 0$$

$$\text{momentum } \frac{1}{a} \frac{\partial W}{\partial t} + \frac{1}{a^2} \frac{\partial M}{\partial X} \cdot \frac{\partial X}{\partial Z} + \frac{M}{a} \frac{d}{dZ} \left(\frac{1}{a} \right) + \frac{\partial P}{\partial Z} + F = -\rho g \frac{dy}{dZ}$$

$$\begin{aligned} \text{energy } a\rho \frac{\partial H}{\partial t} - H \frac{\partial W}{\partial Z} - a \frac{\partial P}{\partial t} + \frac{\partial E}{\partial X} \cdot \frac{\partial X}{\partial Z} + \frac{1}{2a} \frac{\partial M}{\partial X} \cdot \frac{\partial X}{\partial t} + \frac{\partial}{\partial Z} (aK) = \\ = a\phi - Wg \frac{dy}{dZ} \end{aligned}$$

$$\text{where } \frac{\partial}{\partial X} = \begin{pmatrix} \partial/\partial W \\ \partial/\partial P \\ \partial/\partial H \end{pmatrix} .$$

This form has the advantage that, except for the time independent derivatives and $\frac{\partial}{\partial Z} (aK)$ (usually small), only components of X appear in time or space differentials which therefore can be forward averaged. If this is done local stability (Richtmyer 1967) of the finite difference scheme, neglecting the small $\frac{\partial}{\partial Z} (aK)$ term, can be proved (Turner 1972). The drawback is that the derivatives of M and E with respect to P and H may be discontinuous at the boundaries between the two-phase and single phase regions. This is discussed in Section 3.

3. FINITE DIFFERENCE EQUATIONS

Finite differencing sets up averages over the space step Z_1 to Z_2 and over the time step t to t' . The space derivative terms of the form $Q(Z, W, P, H) \partial R/\partial Z$ are approximated by

$$\frac{(Q_1 + Q_2)}{2} \left(\frac{R'_2 - R'_1}{Z_2 - Z_1} \right) \dots (1)$$

and the time derivative terms $Q(Z, W, P, H) \partial R / \partial t$ by

$$\frac{Q_1(R_1' - R_1) + Q_2(R_2' - R_2)}{2(t' - t)}$$

where $Q_1 = Q(Z_1, W_1, P_1, H_1)$ etc. The terms containing $\frac{dy}{dz}$ are treated as follows:

$$\rho g \frac{dy}{dz} = \frac{(\rho_1 + \rho_2)}{2} \left(\frac{Y_2 - Y_1}{Z_2 - Z_1} \right) g$$

The power term is approximated by

$$a\phi = \frac{a_1 \phi_1 + a_2 \phi_2}{2}$$

The term F in the momentum equation is made up of contributions of friction at the pipe wall and at an orifice between the nodes. These two friction terms are specially treated to give better accuracy and to improve numerical stability at high flow rates; both were assumed to vary as $W|W|$ rather than to remain constant over a time step. Then

$$\partial F / \partial W = 2F/W \quad \text{and} \quad \frac{F_1 + F_1'}{2} \approx \frac{F_1 W_1'}{W_1}$$

Thus F is represented by

$$F = \frac{W_1' |G_1| v_{p1}}{a_1 D_{e1}} + \frac{W_2' |G_2| v_{p2}}{a_2 D_{e2}} + \frac{K_o}{Z_2 - Z_1} \left(\frac{W_1' |G_1| v_{o1}}{a_1} + \frac{W_2' |G_2| v_{o2}}{a_2} \right)$$

where K_o and D_e are input quantities and the friction specific volumes v_o and v_p contain the friction factors. Definitions of v_o and v_p are given in Section 6.

The coefficients of the derivatives of E and M are also given special treatment. Consider the expression

$$\left[w \left(\frac{\partial E}{\partial X} \right)_1 + (1 - w) \left(\frac{\partial E}{\partial X} \right)_2 \right] \cdot \left(\frac{X_2' - X_1'}{Z_2 - Z_1} \right)$$

For $w = \frac{1}{2}$ this expression is the form of finite difference approximation given in Equation 1. The truncation error (Richtmyer 1967) in this approximation is of order

$$\left[\left(\frac{\partial E}{\partial X} \right)_1 - \left(\frac{\partial E}{\partial X} \right)_2 \right] \cdot \frac{\partial X}{\partial Z}$$

If $\partial E / \partial X$ is continuous this is of order ∂Z . If $\partial E / \partial X$ is not continuous the

truncation error on a segment spanning the discontinuity is not small, and its effect on the solution is not reduced by choosing ∂Z small. Also in such cases a steady state solution to the finite difference equations for such a segment, may not exist. Both these problems are overcome by choosing

$$w = e \equiv \frac{E_2 - E_1 - \left(\frac{\partial E}{\partial X}\right)_2 \cdot (X_2 - X_1)}{\left[\left(\frac{\partial E}{\partial X}\right)_1 - \left(\frac{\partial E}{\partial X}\right)_2\right] \cdot (X_2 - X_1)} \quad \text{for } 0 \leq e \leq 1 \quad \dots (2)$$

$$w = 0 \text{ for } e < 0$$

$$w = 1 \text{ for } e > 1 \quad .$$

Then, provided ∂Z is not too large a steady state solution always exists and the truncation error is of order ∂Z plus order ∂t .

For single phase flow or if zero slip between the phases is assumed

$$E = H W$$

$$\frac{\partial E}{\partial X} = \begin{pmatrix} H \\ 0 \\ W \end{pmatrix}$$

$$\text{and } w = \frac{1}{2} .$$

Similarly, an interpolation coefficient $0 \leq S \leq 1$ is defined for the term $a^{-2} \frac{\partial M}{\partial Z}$ such that

$$\begin{aligned} & \frac{1}{2} \left(\frac{1}{a_1^2} + \frac{1}{a_2^2} \right) (M_2 - M_1) \\ & = \left[\frac{S}{a_1^2} \left(\frac{\partial M}{\partial X} \right)_1 + \frac{(1-S)}{a_2^2} \left(\frac{\partial M}{\partial X} \right)_2 \right] \cdot (X_2 - X_1) \quad \dots (3) \end{aligned}$$

Then $\frac{1}{a^2} \frac{\partial M}{\partial Z}$ is approximated by

$$\left[\frac{S}{a_1^2} \left(\frac{\partial M}{\partial X} \right)_1 + \frac{(1-S)}{a_2^2} \left(\frac{\partial M}{\partial X} \right)_2 \right] \cdot \frac{(X_2' - X_1')}{(Z_2 - Z_1)} .$$

The resulting three finite difference equations are such that mass flows from one segment appear unchanged in the next, and acceleration (including the effects of area change), friction and gravitational pressure drops in steady state agree with the usual formulations.

The coefficients R_p , R_H , ρ and other properties needed e.g. void fraction and temperature, are evaluated in the equation of state routine described in Section 4. The coefficients $\frac{\partial M}{\partial X}$, $\frac{\partial E}{\partial X}$, K and other quantities

such as quality and E, are evaluated in the slip routine described in Section 5. Evaluation of the friction terms is described in Section 6 and of the power-to-coolant in Section 7.

4. EQUATION OF STATE

The assumption of thermodynamic equilibrium allows the static quantities required for the difference equation coefficients and slip model to be determined from pressure and static enthalpy. The quantities are obtained by interpolation of tables read in from disk at the start of the calculation. The tables have been generated from the state equations for water (Meyer 1967) by means of the ASTEM code (Moore 1971). The pressure and temperature mesh used in the current tables is given in Table 6. Tables with a different mesh can easily be generated if required. The void fraction is always calculated from

$$\alpha = \frac{\rho - \rho_L}{\rho_V - \rho_L}$$

5. SLIP MODEL

In the slip routine quantities P, H and those calculated in the equation of state routine are used to find flow quality

$$x = \frac{H_L - E/W}{H_L - H_V}$$

E, M, K and derivatives $\partial E/\partial X$ and $\partial M/\partial X$. Both of the two slip models available make the approximation that u_v and u_L do not vary across the pipe cross section.

5.1 No Slip and Single Phase

If all the coolant is assumed to move at the same speed, then

$$E = HW$$

$$\frac{\partial E}{\partial X} = \begin{pmatrix} H \\ 0 \\ W \end{pmatrix}$$

$$M = W^2/\rho$$

$$\frac{\partial M}{\partial Z} = M \begin{pmatrix} 2/W \\ -R_P/\rho \\ -R_H/\rho \end{pmatrix}$$

$$K = \frac{G^3}{2\rho^2}$$

and
$$x = \frac{H - H_L}{H_V - H_L} .$$

In the single phase regions the value of x indicates the degree of subcooling or superheating.

5.2 Jones' Slip (Jones & Dight 1962)

In this case

$$u_V = \frac{\alpha u_V + (1-\alpha)u_L}{\gamma}$$

where $\gamma = (0.71 + 0.29P/P_c) (1-\alpha^r) + \alpha^r$,

$$r = 3.33 + (0.18 + 0.46 \times 10^{-3}P)P/1000$$

and P_c , the pressure at the critical point, and P are in psia. Now

$$\alpha = \frac{\rho - \rho_L}{\rho_V - \rho_L}$$

is a function of P and H only. Also

$$G = \alpha \rho_V u_V + (1-\alpha) \rho_L u_L .$$

Therefore $x = \frac{\alpha \rho_V}{\gamma \rho_L - \alpha(\rho_L - \rho_V)}$ is also a function of P and H only.

Thus
$$M = a^2 [\alpha \rho_V u_V^2 + (1-\alpha) \rho_L u_L^2]$$

$$= W^2 \frac{[\alpha \rho_V + \rho_L \frac{(\gamma-\alpha)^2}{(1-\alpha)}]}{[\alpha \rho_V + \rho_L (\gamma-\alpha)]^2} ,$$

and
$$E = a[(1-\alpha) \rho_L u_L H_L + \alpha \rho_V u_V H_V]$$

$$= \frac{W[(\gamma-\alpha) \rho_L H_L + \alpha \rho_V H_V]}{[\alpha \rho_V + (\gamma-\alpha) \rho_L]}$$

are functions W , P and H only. As the derivatives of H_V , H_L , ρ_V , ρ_L are obtained in the equation of state routine, the derivatives of M and E with respect to W , P and H , may be simply obtained. The kinetic energy flux is given by

$$K = \frac{1}{2} [(1-\alpha) \rho_L u_L^3 + \alpha \rho_V u_V^3] .$$

6. FRICTION

The friction F depends on the two friction specific volumes, v_p (a function of the relative pipe roughness ϵ/D) and v_o (a function of the orifice coefficient used for orifices and fittings). The latter depends on quality and the value of S_T (an input quantity specifying the friction model) as shown in Table 1. In single phase regions the orifice pressure drop is assumed proportional to G^2/ρ . In the two-phase region two options are available: if $S_T = 0$ friction corresponding to homogeneous flow is used; if $S_T = 1$, a separated flow model is used, except at high quality where the homogeneous model is used. To remove a discontinuity between separated and homogeneous flow regions an interpolation region is inserted.

The ϕ in Tables 1 and 2 is the product of approximations from Cowking & Enderby (1966) to the Martinelli-Nelson (1948) two-phase pressure drop multiplier and Collier (1964) mass velocity correction term. At low qualities this product becomes less than 1. If this occurs ϕ is taken to be 1.

To specify v_p we first define a Fanning friction factor f as a function of μ , a viscosity. The Reynolds number is given by

$$Re = \frac{GD}{\mu} \quad .$$

If $Re < 2000$, laminar flow is assumed and

$$f(\mu) = \frac{16}{Re} \quad .$$

If $Re \geq 4000$, the Colebrook equation for the transition and turbulent flow regions is used.

$$\frac{1}{\sqrt{f}} = -1.73716 \ln \left(\frac{\epsilon}{3.7D} + \frac{1.26}{Re\sqrt{f}} \right) \quad .$$

Often a Reynolds number of 2000 is taken as the boundary between the laminar flow equation and the Colebrook equation. To avoid the discontinuity that such an assumption would introduce, the following interpolation relation, which gives reasonable agreement with Nikuradse (1933) was introduced for Reynolds numbers between 2000 and 4000.

$$f(\mu) = \frac{16}{Re} + [f(Re=4000) - 0.004] \frac{(Re-2000)}{2000} \quad .$$

The function $f(\mu)$ is used in Table 2 to define v_p . The input quantity S_T also determines the option used for v_p in the same way as for v_o .

7. PIPE MODEL

A single region heat storage model was used with forward averaging of the fuel temperature to improve stability. The finite difference equation is

$$\frac{\lambda(T' - T)}{\delta t} = q - h_c(T' - T_c) \quad \dots(4)$$

where

$$h_c = \frac{h_F h_s}{h_F + h_s} ;$$

then $a\phi = ph_c(T' - T_c)$.

The pipe surface temperature needed to estimate h_s for the next time step is obtained by assuming the heat flux from pipe to pipe surface is equal to that from pipe surface to coolant:

$$h_F(T - T_s) = h_s(T_s - T_c) .$$

8. HEAT TRANSFER

Steady state heat transfer correlations expressed as a function of local conditions (coolant temperature, pressure, quality and pipe surface temperature) were assumed applicable to transient situations. This is the local condition hypothesis of Hassid & Rychlicki (1971) which predicts dryout somewhat sooner than it occurred in experiments.

Heat transfer correlations included to cover all situations are summarised in Table 3. Condensing heat transfer correlations are needed when steam or two-phase coolant heated in a heated pipe section moves to sections in which the pipework is considerably cooler. They may be required also if the pipe model is used to simulate a constant temperature heat sink (e.g. heat exchanger) by giving the pipe a very large heat capacity. Akers et al. (1959) gives a Reynolds number of 50 000 for the boundary between their two condensing heat transfer correlations. However, the correlations actually intersect at 76 228 and this value is used here. Also Akers correlations only cover the two-phase region. McAdams' (p.351 1954) recommendation that heat flux is almost independent of superheating of the vapour, is followed in the superheat region. In the boiling non-dryout region, the heat transfer coefficient corresponding to three correlations is calculated and the largest taken. This ensures continuity of surface heat transfer rate in this region as a function of coolant conditions and fuel surface temperature.

At each heat transfer coefficient calculation, pipe surface temperature and coolant state are known from the previous time step. The heat transfer coefficient is evaluated assuming no dryout, then if appropriate, tested to determine if the dryout limit is exceeded. If it is, the appropriate dryout heat transfer coefficient is calculated. In the steady state calculations, heat flux is known and surface temperature must be determined. Here an iterative procedure is used. In those steady state calculations where dryout and non-dryout solutions exist, the latter solution is used.

An indicator THT is available to show which heat transfer correlations are used in evaluating the heat transfer coefficient.

$$\text{THT} = \ell + m + n$$

where ℓ and m indicate the correlation used to evaluate the coefficient assuming no dryout, m indicates the dryout correlation used to determine if the dryout limit is exceeded, and n the post-dryout correlation used. In some situations, a dryout correlation is not appropriate, e.g. during condensation; zero values of m and n indicate these. Possible values of ℓ , m and n and their meanings are given in Table 3. For example, $\text{THT} = 313$ would indicate that:

- . the heat transfer coefficient assuming no dryout, was evaluated using the nucleate boiling correlation;
- . the dryout limit on the heat transfer coefficient was determined from Bernath's correlation;
- . dryout was found to be present; and
- . Beattie's post-dryout correlation was used to evaluate the actual heat transfer coefficient used.

9. BOUNDARY CONDITIONS

Three boundary conditions are needed for solution of the 3J-3 difference equations containing the 3J unknowns, W_i , P_i , H_i , $i = 1$ to J :

- . inlet enthalpy;
- . inlet flow or inlet pressure; and
- . pressure or a choked flow model $W_J = W_J(P_J, H_J)$ at the outlet.

These boundary conditions are given in the input data as a function of time.

9.1 Moody Model

The choked or critical mass velocity is contained in tables (Redfield et al. 1967) as a function of stagnation pressure and enthalpy. Pressure at the last node is taken as stagnation pressure, and stagnation enthalpy

is given by

$$H_o = \frac{E_J + a_J K_J}{W_J}$$

Mass flow at node J is then the product of choked mass velocity and break-area given in the input data.

9.2 Rodflow Model

This method, taken from Elliott (1968), uses a function fit to Moody results to give choked mass velocity at the end of a pipe from a large reservoir, as a function of stagnation pressure, stagnation enthalpy and fL/D of the pipe. An fL/D of zero gives results close to those obtained using the Moody table of 9.1 above. The relation is

$$G_c = I \left(\frac{P_J}{H_o - 300} \right)^Y$$

where I and Y are obtained from Table 4 and the units of G_c , H_o and P_J are $\text{lb ft}^{-2} \text{ s}^{-1}$, Btu lb^{-1} and psi , respectively. The stagnation quantities are obtained as for the Moody model.

9.3 Fauske Model

Discharge is assumed to occur through a pipe of diameter and length specified in the input data. Conditions at the end of the pipe are determined from conditions at node J by assuming the following:

- . steady state;
- . adiabatic flow;
- . thermodynamic equilibrium between phases;
- . Jones' slip correlation;
- . annular friction with zero roughness (Beattie 1973); and
- . choked flow as described below at the pipe exit.

These assumptions have been found to give reasonable agreement with Fauske's measured pressure profiles (Turner & Trimble 1974).

From conditions at the throat, choked flow is determined as follows. In the quality range 0.0025 to 0.99 the Fauske model, as specified by Hsu (1972), is used. In the single phase regions, flow at sonic speed is assumed. As problems in which a transfer between the single and two-phase region occurs can only be handled if choked flow is a continuous function of its parameters (here quality and pressure), the following joining functions are used:

$$0 < x < 0.0025, \quad G_c = G_F^{400x} G_L^{(1-400x)}$$

where G_F = Fauske choked mass velocity for quality x , and G_L = liquid choked mass velocity for $x = 0$ at the same pressure;

$$0.99 < x < 1, \quad G_c = 100 [(1-x)G_F + (x-0.99)G_s]$$

where G_s = steam choked mass velocity for $x = 1$ at the same pressure. Liquid and steam mass velocities are both calculated from the usual single phase speed of sound formulae

$$\left(\frac{\rho}{G}\right)^2 = \left(\frac{\partial \rho}{\partial P}\right)_{\text{entropy}} = R_p + \frac{R_H}{\rho}$$

10. SOLUTION OF FINITE DIFFERENCE EQUATIONS

The fuel temperature T' at the end of the time step is first eliminated from the finite difference conservation of energy equation (Section 3) using the pipe model equation (Section 7). The resulting three finite difference equations are linear in the values of W , P and H at the end of the time step. These equations may be written

$$A_j X_j' + B_j X_{j+1}' = C_j, \quad j = 1, 2, \dots, J-1 \quad \dots (5)$$

where the three-component vector $X_j' = \begin{pmatrix} W_j' \\ P_j' \\ H_j' \end{pmatrix}$,

A and B are 3×3 matrices and C is a three-component vector. A , B and C are not functions of W_j' , P_j' or H_j' . The three boundary condition equations involving components of X_1' and X_J' , and the $(J-1)$ matrix equations are solved by an elimination method as follows.

Define matrices D and vectors V such that

$$X_j' + D_{j-1} X_1' = V_{j-1}, \quad j = 2, 3, \dots, J \quad \dots (6)$$

Eliminating X_j' from Equations 5 and 6 we get

$$X_{j+1}' - B_j^{-1} A_j D_{j-1} X_1' = B_j^{-1} (C_j - A_j V_{j-1})$$

Therefore

$$D_j = -B_j^{-1} A_j D_{j-1} \quad \dots (7)$$

and $V_j = B_j^{-1} (C_j - A_j V_{j-1})$ for $j = 2, 3, \dots, J-1$.

Also from Equations 5 and 6

$$D_1 = B_1^{-1} A_1$$

$$\text{and } V_1 = B_1^{-1} C_1 \quad \dots (8)$$

Using Equations 7 and 8, D_{J-1} and V_{J-1} can be found. Equation 6 with $j = J$ and the three boundary conditions are solved for X'_1 and X'_J . Then Equation 6 is used to find all X'_j . $2J$ matrix multiplications and $(J-1)$ inversions of 3×3 matrices are the main computations required. The quantity

$$R_{\text{off}} = \text{MAX} \left| \frac{i^{\text{th}} \text{ component of } (A_j X'_j + B_j X'_{j+1} - C_j)}{i^{\text{th}} \text{ component of } C_j} \right|$$

for all i, j such that the denominator is non-zero, is a measure of accuracy of the solution of the difference equations. A message (Section 13.2) is printed if it exceeds 10^{-7} .

This solution procedure assumes that the determinate of B_j is non-zero. In fact, if this determinate is treated as a function of the time step, it has 3 zeros. These zeros have been evaluated for limited situations and are approximated by

$$\frac{\delta Z}{2\delta t} = \frac{G}{\rho}, \quad -\frac{G}{\rho} \pm c$$

where c is the thermodynamic equilibrium speed of sound. The only case where the presence of these zeros affected calculation was when a number of segments of equal length contained fluid in nearly the same state, so that the zeros for each B -determinate almost coincided. In these cases, for some time steps close to one of the zeros, the solution procedure failed and high values of R_{off} resulted. However, changing the finite difference mesh to segments of unequal length removed this difficulty.

11. INPUT DATA

Input data for a sample problem are shown in Appendix A. Each problem requires a title card, a control array and a number of tables.

11.1 Title Card - 1 card

Any 80 characters.

11.2 Control Array - 5 cards

Six numbers per card in format 6D12.3, containing:

CC(1) Not used.

CC(2) Initial pressure in psia at node CC(8). This is used only if an initial steady state calculation is specified, i.e. CC(8) > 0.

- CC(3) Initial value of the time step(s). This time step is only used if it satisfies the time step controls CC(13) to CC(16).
- CC(4) Inlet boundary condition indicator: if CC(4)=0 then inlet flow is a boundary condition; if CC(4)=1 then inlet pressure is a boundary condition. All boundary conditions are given in the boundary condition table (Section 11.3).
- CC(5) Finish time(s).
- CC(6) Number of space nodes J , $2 < J < 50$.
- CC(7) Relative roughness of channel walls ϵ/D_e .
- CC(8) Initial condition indicator:
- . <0 if conditions at each space node stored in the computer from a case immediately before are used. (Such a case must not have fewer nodes than the present one.);
 - . $=0$ if initial conditions at each space node are read in from the restart table (Section 11.5). (Cards for a restart are punched out on completion of every case.);
 - or
 - . >0 if initial steady state conditions are calculated. Four input data specify this steady state: mass velocity for time zero in the boundary condition table; pressure P and thermodynamic quality x at a specified node given as CC(2) and CC(10) respectively; CC(8) is the number of the specified node.
- CC(9) Not used.
- CC(10) Initial thermodynamic quality x at node CC(8). This is only if initial steady state calculation is specified, i.e. CC(8) >0 .
- CC(11) Number of time entries one to a card in the boundary condition table; ≤ 24 .
- CC(12) Output indicator: 0, if the full printout consists of only the first row of quantities at each node shown in Appendix B; -2, if all quantities used in evaluating the difference equation coefficients are included in each full printout as shown in Appendix B. (CC(20) specifies how often a full printout occurs.)
- CC(13) Time step controls. Values of the absolute fractional changes to
- CC(16) in pressure P and static enthalpy H are calculated for each node. If B is the largest of these, then for $B < CC(13)$ the

time step is multiplied by CC(14); alternatively if $B > CC(15)$ the time step is multiplied by CC(16) and repeated. Messages indicate changes in time step. Values of 0.01, 2.0, 0.05, 0.5 usually give reasonable accuracy.

- CC(17) Not used.
- CC(18) Maximum number of time steps for this case.
- CC(19) Not used.
- CC(20) At the first, the last and every CC(20)th time step the full printout as specified by CC(13) is produced. At other time steps the output specified by CC(21) and the output table (Section 11.6) is given.
- CC(21) The number of quantities in addition to time, to be printed at every time step. Any quantity stored in the main common area DIM (Table 5) may be so printed. The quantities are specified in the output table (Section 11.6). Any function of DIM quantities and desired input data can be specified in a special output routine, SOUT, and stored in the spare AZ array in DIM. The location of input data in common area DATA is given in Table 7. SOUT is called by the OUTPUT routine at every time step.
- CC(22) Choked flow indicators. If CC(23) > 0 the boundary condition at the channel outlet is one of three choked flow models described in Section 9:
- . ≥ 0 , Rodflow ($CC(22) = fL/D \leq 20$);
 - . -1, Moody; or
 - . -2, Fauske.
- CC(23) Choked flow indicator: > 0 if the boundary condition at the channel outlet is a choked flow relation with a break area of CC(23) ft^2 .
- CC(24) Break opening time(s). The break area is assumed to increase linearly from 0 to CC(23) ft^2 over CC(24) seconds and then remain constant.
- CC(25) Not used.
- CC(26) Hydraulic diameter of the discharge pipe (ft). Only used if the Fauske choked flow model is specified, i.e. CC(22) = -2.
- CC(27) Length of the discharge pipe (ft). Only used if the Fauske choked flow model is specified, i.e. CC(22) = -2.

CC(28) Minimum value of the time step. A message is printed if this limit is applied.

CC(29) Slip indicator: 0 if no slip, 1 if Jones' slip correlation. The slip models are described in Section 5.

11.3 Boundary Condition Table

CC(11) cards in format 6D12.3. On each card (enter zero if not required):

- (a) time(s);
- (b) inlet mass velocity G ($\text{lb s}^{-1} \text{ft}^{-2}$) (only required if $\text{CC}(4)=0$ or $\text{CC}(8)>0$);
- (c) change in inlet pressure from the initial value (psia) (only required if $\text{CC}(4)=1$);
- (d) change in inlet enthalpy from the initial value (Btu lb^{-1});
- (e) change in outlet pressure from the initial value (psia) (only required if $\text{CC}(23) \leq 0$); and
- (f) total power (MW).

11.4 Geometry Table

Two cards per node in format 6D12.3. For each node:

- (a) area a (ft^2);
- (b) length to the next node δZ (ft);
- (c) increase in height to the next node δy (ft);
- (d) equivalent diameter D_e (ft), used in the friction and heat transfer calculation;
- (e) power into the pipe per unit heated wetted surface area q - in any units since within the code these numbers are renormalised to give the total power of the boundary condition table;
- (f) orifice coefficient K_o for pressure losses in fittings, bends or orifices between this and the next node. (Note that the effect of accelerations due to area changes specified in (a) above must not be included here. See Section 3.);
- (g) heat capacity of fuel per unit heated wetted surface λ ($\text{Btu ft}^{-2} \text{F}^{-1}$). (If 0 then no pipe model calculation is needed and the power if any, is put directly into the coolant.)
- (h) heated wetted perimeter p (ft);
- (i) pipe surface to pipe centre heat transfer coefficient h_F ($\text{Btu s}^{-1} \text{ft}^{-2} \text{F}^{-1}$) - see Section 7. (Not used if $\lambda = 0$.);
- (j) friction model indicator S_T : 0 if the homogeneous model, 1 if Martinelli-Nelson.

Note that entries for δZ , δy and K_o for the last node are not used as there is no 'next node'.

11.5 Restart Table

Contains $2*CC(6)+1$ cards in format 3D24.10.

The first card is

time(s), power (MW);

then two cards per node containing

flow (lb s^{-1}),

pressure (psia)

enthalpy (Btu lb^{-1}), and

pipe temperature (F) at the node.

11.6 Output Table - format 6(A8,I4)

The values of quantities specified in this table and the time are printed at every time step - the first 9 on FORTRAN Unit 10, the next 9 on Unit 11, the next 9 on Unit 12, etc. Note that the required FORTRAN units must be specified by control cards, e.g. the sample problem in Appendix A. The quantities are also written on Unit 9 in unformatted records suitable for input to the PLOT program. The first case is put on the first file (FT09F001), and the second on the next (FT09F002), etc.

Output quantities are specified by their position in the common area DIM (Table 5). For each quantity there must be an 8-character heading to appear in the output and an integer giving the position of the quantity in DIM. Six quantities are specified per card.

12. COMPUTING TIME

The computing time on an IBM360/50 for problems not involving a choked flow model is approximately

$$(2.5 + N/10)J \text{ seconds}$$

where N is the number of time steps. Choked flow problems may take up to 50 per cent more time. The sample problem in the appendices took 75 s for a 12 s transient using 14 nodes.

A full printout of the last time step is always given since the remaining time before a system abend is monitored and no more time steps are computed when there is insufficient time.

13. OUTPUT

Output from the sample problem is shown in Appendix B. A listing of the control array appears first followed by a table giving the channel geometry, the same as in the geometry table of input data, except that power distribution data are normalised to $\text{Btu ft}^{-2} \text{ s}^{-1}$ per MW total power. The output table is given unchanged. The boundary condition table follows, giving values of actual boundary conditions, rather than changes from the initial values as

was specified in most of the input data boundary condition table. The remainder of the output consists of quantities specified in the output table (Section 11.6) at every time step, and a full printout when specified. Some columns in this full printout are headed by symbols which are explained in Table 5. If $CC(12) = 0$ these columns are not printed.

13.1 Error Messages

The code user specifies the finite difference mesh, *i.e.* the geometry and time step size criteria to be used in the calculation. If this mesh is too coarse the results are inaccurate. The effect of changing the mesh should be evaluated. If the mesh is very coarse the finite difference scheme breaks down and extreme pressures or enthalpies may be calculated. This is the most common cause of code failure. Error messages are:

(a) **TIME STEP LIMIT** Although the maximum fractional change in pressure and enthalpy is greater than $CC(13)$, the time step reduction $CC(14)$, is not made because of the minimum time step restriction, $CC(28)$. The calculation continues. The effect of this restriction can be evaluated by running cases with different values of the restriction $CC(28)$.

(b) ****WARNING-ROFF** d R_{off} is a measure of the error in the solution of the difference equations (Section 10). This message appears if the value of R_{off} exceeds 10^{-7} ; d is its actual value. Calculation continues. If d is of order 1 the solution is not acceptable. Changing the mesh to avoid the singularities discussed in Section 10 has always removed this difficulty.

(c) **RANGE OF TABLE TOO SMALL** The time interval covered by the boundary condition table is insufficient to perform the calculation specified in the input data. The calculation stops.

(d) **NODE i ONLY CONVERGED TO x** In the steady state calculation pressure and enthalpy at each node are obtained by a convergence procedure. At node i , x is the fractional difference between the pressures or enthalpies from the last two iterations. If $|x| > 10^{-14}$ the message appears and calculation continues with the unconverged value.

(e) **LEAKAGE FLOW NOT CONVERGED** x y Application of the choked flow correlation involves an iteration procedure. If the Fauske model is being used then x is the sum of the squares of the fractional errors in stagnation enthalpy and pressure. The message appears if $x > 10^{-4}$, although convergence is generally much better than this. For other critical flow nodes x should be small compared to the critical flow y . Then the message appears if $\left|\frac{x}{y}\right| > 10^{-5}$. Calculation continues.

(f) FUEL TEMPERATURE AT NODE i FAILED TO CONVERGE d, T, h THT

In the steady state calculation fuel surface temperature is calculated by an iteration procedure. If this fails at node i this message is printed with values from the last iteration: d is the difference in fuel temperature between the last two iterations, T is the fuel surface temperature, h is the fuel surface to coolant heat transfer coefficient and THT is the heat transfer indicator (Table 3).

(g) FRICTION FACTOR AT NODE i NOT CONVERGED x y The Colebrook equation is solved by an iteration procedure. If the procedure fails at node i the last two values of $f^{-0.5}$ are printed. Calculation continues.

(h) CRITICAL FLOW FAILED x y For meaningful pressure P_J as calculated from the finite difference equations, W_J must be between the limits x, y. No such choked flowrate was found. Calculation stops.

(i) **H x OUT OF RANGE - EXTRAPOLATING An enthalpy of x Btu lb⁻¹, which was outside the range of the table, was passed to the equation of state routine. The tables were extrapolated and calculation continued.

(j) **P x OUT OF RANGE - EXTRAPOLATING As for i, but with pressure (psi).

14. PLOTTING

Any quantity stored in the main common area DIM (Table 5) may be plotted against time. One or more quantities may be plotted per graph in combination with points or curves from input cards. Calculated quantities to be plotted must be specified in the output table (Section 11.6). Appendix A contains an example of input data to the PLOT routine. Input data are as follows:

(a) First set of cards (6I12)

NC number of quantities to be plotted on this plot with scales and axes as specified below. Must be >0.

KQ(I), I=1, NC Positions in the output table (Section 11.6) of the quantities to be plotted.

(b) Next card (2F12.3, I12, F12.3, I12, F12.3)

YMIN minimum value on y axis. This is compared with the lowest value of the quantity to be plotted and the lower used.

YMAX maximum value on y axis. This is compared with the highest value of the quantity to be plotted and the higher used.

NE number of curves to be plotted from cards on this plot with the scales and axes as specified.

XMAX the maximum value on the x or time axis is the larger of XMAX and the time at the end of the calculation.

NU if positive the quantities are plotted on a new plot; if negative they are plotted on the most recent plot, however the new x and y scales are not drawn. |NU| is the number of the FORTRAN Unit containing the necessary OWEN-1 output (Section 11.6).

CF conversion factor. All y values including those from cards are multiplied by CF before plotting. YMAX and YMIN are not changed.

(c) Next card (2F12.2)

XS >0, for log scale on x axis, <0 for a linear scale. |XS| is the length of the x axis in inches.

YS same as XS but for the y axis.

(d) Input data curves (I12,/(6F12.3))

Only required if NE >0.

One set of cards for each of the NE input data curves, containing

NP >0, if lines joining the points are drawn; <0 if only points are plotted.

Y(J),X(J),J=1,NP coordinates of points to be plotted.

15. ACKNOWLEDGEMENTS

The author is grateful to Mr. A.G. Chapman for information on dryout limits and condensing heat transfer; to Mr. D.R.H. Beattie for advice on post dryout heat transfer and two-phase friction; to Mr. V. Ilic for helpful discussions of choked flow; and to Mr. G.D. Trimble for improving the solution of the difference equations and coding the equation of state routine.

16. REFERENCES

- Akers, W.W., Deans, H.A. & Crosser, O.K. (1959) - Chem. Eng. Program Symp. Ser.No. 29 (55) 171.
- Beattie, D.R.H. (1972) - AAEC private communication.
- Beattie, D.R.H. (1973) - Nucl. Eng. Des. 25 : 395.
- Biasi, L., Clerici, G.C., Gerribba, S., Sala, R., Tozzi, A. (1967) - Energia Nucleare, 14 (9) 530.
- Carter, J.C., Fischer, G.J., Heames, T.J., MacFarlane, D.R., et al. (1970) - SASIA. ANL7607.
- Carver, M.B. (1967) - ASME Paper 67-HT-66 presented at the ASME-AICHE National Heat Transfer Conference, Seattle, August.

- Collier, J.G. (1964) - AERE Report RS/L207.
- Currin, H.B., Hunin, C.M., Rivlin, L. & Tong, L.S. (1961) - CVNA-77.
- Dittus, F.W. & Boelter, L.N.K. (1930) - Univ. of Calif. Pub. Eng., 2 : 443.
- Elliott, J.N. (1968) - RODFLOW. A Program for Studying Transients in a Power Reactor Coolant Circuit. AECL Power Projects. TDAI-11.
- Fauske, H.K. (1962) - ANL-6633.
- Hassid A. & Rychlicki, R. (1971) - Energia Nucleare 18 (6) 333.
- Heineman, J.B. (1960) - ANL-6213.
- Hsu, Y.Y. (1972) - NASA TN D-6814.
- Jones, A.B. & Dight, D.G. (1962) - KAPL-2208.
- McAdams, W.H. (1954) - Heat Transmission. 3rd Ed. McGraw-Hill, N.Y.
- Martinelli, R.C. & Nelson, D.B. (1948) - Trans. ASME (August).
- Meyer, C.A. et al. (1967) - ASME Steam Tables. ASME, N.Y.
- Meyer, J.E. (1960) - WAPD-BT-20.
- Meyer, J.E. (1961) - Nucl. Sci. Eng. 10 : 269.
- Moore, K.V. (1971) - ANCR-1026.
- Moxon, D. (1968) - AEEW-R448.
- Nikuradse, J. (1933) - VDI Forschungshft. 361.
- Porsching, T.A., Murphy, J.H., Redfield, J.A., et al. (1969) - FLASH-4. WAPD-TM-840.
- Porsching, T.A., Murphy, J.H. & Redfield, J.A. (1971) - Nucl. Sci. Eng. 43 : 218.
- Premoli, A. (1969) - Energia Nucleare, 16 (10) 625.
- Redfield, J.A., Murphy, J.H. & Davis, V.C. (1967) - FLASH-2. WAPD-TM-666.
- Rettig, W.H., Jayne, G.A., Moore, K.V., Slater, C.E. & Uptmoor, M.L. (1970) - RELAP 3. IN-1321.
- Richtmyer, R.D. & Morton, K.W. (1967) - Difference Methods for Initial-value Problems. 2nd Ed. Interscience Publishers, N.Y.
- Sha, W.T. & Hughes, T.H. (1970) - VENUS. ANL-7701.
- Shrock, V.E. & Grossman, L.M. (1959) - TID-14632.
- Tong, L.S. (1965) - Boiling Heat Transfer and Two-phase Flow. Wiley, N.Y.
- Turner, J. & Wendroff, B. (1964) - LA-3007.
- Turner, W.J. (1969) - CREST Specialist Meeting, Frankfurt.
- Turner, W.J. (1972) - Energia Nucleare, 19 (5) 317.
- Turner, W.J. & Trimble, G.T. (1974) - Fifth Australasian Conference of Hydraulics and Fluid Mechanics, Christchurch, NZ, 9-13 December.
- Wright, R.M. (1961) - USAEC Report UCRL-9744.

TABLE 1
ORIFICE FRICTION SPECIFIC VOLUME v_o

Quality	$S_T = 0$	$S_T = 1$
$x \leq 0$	$\frac{1}{\rho}$	
$0 < x \leq 0.8$		$\frac{\phi}{\rho_L}$
$0.8 < x < 0.9$		$10 \left[\begin{array}{l} (0.9-x)\phi/\rho_L \\ + (x-0.8) \left(\frac{1-x}{\rho_L} + \frac{x}{\rho_v} \right) \end{array} \right]$
$0.9 \leq x < 1$		$\frac{1-x}{\rho_L} + \frac{x}{\rho_v}$
$x \geq 1$	$\frac{1}{\rho}$	

TABLE 2
PIPE FRICTION SPECIFIC VOLUME v_p

Quality	$S_T = 0$	$S_T = 1$
$x \leq 0$	$f(\mu_L)/\rho$	
$0 < x \leq 0.8$		$\phi f(\mu_L)/\rho_L$
$0.8 < x < 0.9$		$10 \left\{ \begin{array}{l} (0.9-x)\phi f(\mu_L)/\rho_L \\ + (x-0.8) \left(\frac{1-x}{\rho_L} + \frac{x}{\rho_V} \right) f \left[\frac{\mu_L \mu_V}{(1-x)\mu_V + x\mu_L} \right] \end{array} \right\}$
$0.9 \leq x < 1$		$\left(\frac{1-x}{\rho_L} + \frac{x}{\rho_V} \right) f \left[\frac{\mu_L \mu_V}{(1-x)\mu_V + x\mu_L} \right]$
$x \geq 1$	$f(\mu_V)/\rho$	

TABLE 3
CORRELATIONS FOR SURFACE HEAT TRANSFER COEFFICIENT

Heat Transfer Indicator THT = $\ell + m + n$

Surface Temperature T_s	Quality x		
	$x \leq 0$	$1 > x > 0$	$x \geq 1$
$T_s < T_{sat}$	<u>Liquid</u> Dittus, Boelter (1930) $\ell = 1, m = 0, n = 0$	<u>Condensing</u> Akers (1959) $R_e \leq 76228 \ell = 3$ $m = 0, n = 0$ $R_e^e > 76228 \ell = 4$	
$T_s \geq T_{sat}$	<u>Boiling</u>		<u>Superheat</u>
	<u>Non-dryout</u>		
	Largest of - subcooled liquid		
	Dittus, Boelter (1930)	$\ell = 5$	
	- nucleate boiling		
	Tong p.118 (1965)	$\ell = 3$	
	- forced convection boiling		
	Shrock & Grossman (1959) with coefficient from Wright (1961)	$\ell = 2$	
	<u>Dryout</u>		Heineman (1960)
	$x \leq 0$, Bernath (Tong p.156, 1965)	$m = 10$	$\ell = 2$
	$x > \rho_v / (\rho_v + \rho_L)$, Biasi (1967)		$m = 0$
	low quality	$m = 30$	$n = 0$
	high quality	$m = 40$	
	Other x , linear interpolation on x between Bernath & Biasi low quality		
		$m = 20$	
	<u>Post Dryout</u> - Dryout limit not exceeded		
		$n = 0$	
	$x \leq 0.05$, Tong p.129 (1965)		
	low speed flow	$n = 100$	
	high speed flow	$n = 200$	
	$x > 0.05, \frac{h_s D_e}{k_v} = 0.0192 \left(\frac{GD_e x}{\mu_v} \right)^{0.8} \left(\frac{C_{pv} \mu_v}{k_v} \right)^{0.4}$		
	evaluated at temperature T_s		
	Beattie (1972)	$n = 300$	

TABLE 4
RODFLOW CHOKED FLOW MODEL
(Section 9.2)

FL/D	I	Y
0	1850	1.073
1	1500	1.088
2	1140	1.175
3	920	1.227
4	788	1.261
5	674	1.290
10	442	1.373
20	306	1.359

TABLE 5
REAL*8 ARRAYS IN COMMON AREA DIM

Dimensioned Array	Location	Definition - i, j and k are the first, second and third subscripts respectively
A(3,3,50)	1	Coefficients in the difference equations $\sum_{j=1}^3 [A_{ijk} X_{jk} + B_{ijk} X_{j(k+1)}] = C_{ik}$ k = 1, 2 and 3 for the mass, momentum and energy equations respectively
B(3,3,50)	451	
C(3,50)	901	
X(3,50)	1051	For i = 1,2,3 this is W _j , P _j , H _j in lb s ⁻¹ , psia, Btu lb ⁻¹ respectively at the end of the time step
U(50)	1201	Internal energy Btu lb ⁻¹
R(50)	1251	Density ρ lb ft ⁻³
E(50)	1301	Product of flow enthalpy and flow rate (Section 2) E Btu s ⁻¹
DM(3,50)	1351	∂M _j /∂X _{ij} (Section 3)
RLP(50)	1501	dρ _L /dP lb ft ⁻³ psi ⁻¹
XQ(50)	1551	Static quality calculated from static enthalpy
T(50)	1601	Coolant temperature T _c (°F)
TS(50)	1651	Coolant saturation temperature (°F)
RH(50)	1701	R _H lb ² ft ⁻³ Btu ⁻¹
RP(50)	1751	R _p lb ft ⁻³ psi ⁻¹
F(50)	1801	Frictional force/unit volume of coolant F (poundal ft ⁻³)
D(3,3,50)	1851	Equation 6 of Section 10 is $X_{ik} + \sum_{j=1}^3 D_{ijk} X_{ij} = W_{ik}$
W(3,50)	2301	
G(50)	2451	Mass velocity G lb ft ⁻² s ⁻¹
XS(3,50)	2501	X at the start of a time step
XC(3,50)	2651	X at the start of a time step
RS(50)	2801	Not used
AL(50)	2851	Void fraction
HSL(50)	2901	Saturated liquid enthalpy H _L Btu lb ⁻¹
HSV(50)	2951	Saturated vapour enthalpy H _v Btu lb ⁻¹
RSL(50)	3001	Saturated liquid density ρ _L lb ft ⁻³
RSV(50)	3051	Saturated vapour density ρ _v lb ft ⁻³
XF(50)	3101	Quality x
RM(50)	3151	M lb ft ³ s ⁻²
RVP(50)	3201	dρ _v /dP lb ft ⁻³ psi ⁻¹
GS(50)	3251	2G G v _o (Section 6)
RITP(50)	3301	Momentum equation interpolation factor S (Section 3)

(continued)

TABLE 5 (continued)

Dimensioned Array	Location	Definition - i, j and k are the first, second and third subscripts respectively
HVP(50)	3351	dH_V/dP , Btu lb ⁻¹ psi ⁻¹
HLP(50)	3401	dH_L/dP , Btu lb ⁻¹ psi ⁻¹
DE(3,50)	3451	$\partial E_j / \partial X_{ij}$
EITP(50)	3601	Energy equation interpolation factor w. (Section 3)
TPR1	3651	Power at the start of the time step, MW
ROFF	3652	Fractional error in solution of difference equations (Section 10)
AZ(38)	3653	Not used, available to code user for output from SOUT
XO	3691	Quality at exit of discharge pipe if Fauske choked flow model used
HO	3692	Stagnation enthalpy in discharge pipe, Btu lb ⁻¹
P2	3693	Pressure at exit of discharge pipe if Fauske choked flow model used, psi
H2	3694	Static enthalpy at exit of discharge pipe if Fauske choked flow model used, Btu lb ⁻¹
DET(50)	3701	Determinate of matrix B for segment i
BI(3,3,50)	3750	Inverse of B for node k
EX(3,50)	4200	Change in X_{ij} over the time step
Q(3,50)	4350	Working area for solution of the difference equations
TF(50)	4501	Pipe temperature T_F at the end of the time step (°F)
TFC(50)	4551	Pipe temperature T_F at the start of the time step (°F)
HT(50)	4601	Heat transfer coefficient h_c (Section 7) Btu ft ⁻² s ⁻¹ °F ⁻¹
THT(50)	4651	Heat transfer indicator (Table 3)
QC(50)	4701	Power to coolant per unit length along channel, Btu ft ⁻¹
PC	4751	Total power to coolant, MW
TPR2	4752	Power at the end of the time step, MW
TFS(50)	4753	Fuel surface temperature, T_S °F
EK(50)	4803	Kinetic energy flux lb s ⁻³

TABLE 6
EQUATION OF STATE MESH

Pressures (MPa)

0.090 to 0.200 at intervals of 0.005

0.20 to 1.00 at intervals of 0.01

1.00 to 1.50 at intervals of 0.02

1.50 to 4.00 at intervals of 0.05

4.0 to 10.0 at intervals of 0.1

10.0 to 16.0 at intervals of 0.2

Temperatures (K)

at each pressure point 5 liquid and 6 vapour points

Liquid:	T_{sat}	Vapour:	T_{sat}
	$T_{\text{sat}-10}$		$T_{\text{sat}+10}$
	$T_{\text{sat}-20}$		$T_{\text{sat}+20}$
	$T_{\text{sat}-50}$		$T_{\text{sat}+50}$
	273.15		$T_{\text{sat}+100}$
			1073.15

TABLE 7

REAL*8 ARRAYS IN COMMON AREA DATA

Dimensioned Array	Location	Definition	Units
CC(50)	1	Control array	
TIME	76	t	s
		<u>Boundary conditions</u>	
TI(24)	127	Time	s
GII(24)	151	Inlet mass velocity	lb ft ⁻¹ s ⁻²
PII(24)	175	Inlet pressure	psia
HII(24)	199	Inlet enthalpy	Btu lb ⁻¹
PIN(24)	223	Outlet pressure	psia
POW(24)	247	Power	MW
TITLE(20)	271	Title card in A format	
		<u>Geometry</u>	
AR(50)	293	Area at node i	ft ²
SL(50)	343	Segment length L node i to i+1	ft
SH(50)	393	Segment height y node i to i+1	ft
SD(50)	443	Hydraulic diameter D _e	ft
PD(50)	493	Power distribution	$\frac{\text{Btu ft}^{-2} \text{ s}^{-1}}{\text{MW total power}}$
CORIF(50)	549	Orifice coefficient K _O for segment	
CP(50)	699	Pipe heat capacity λ	Btu ft ⁻² °F ⁻¹
HP(50)	749	Heated perimeter p	ft
HF(50)	799	Heat transfer coefficient pipe to pipe surface h _p	Btu ft ⁻² s ⁻¹ °F ⁻¹
ST(50)	849	Friction model indicator S _T	

APPENDIX A
INPUT DATA FOR SAMPLE PROBLEM

In this problem the flow path from inlet to outlet is as follows:

- (a) a constant displacement pump,
- (b) a hydraulic damper simulated by a section of large diameter pipe of correct volume,
- (c) a preheater,
- (d) feeder,
- (e) a throttling valve,
- (f) an inlet valve,
- (g) heated test section,
- (h) outlet valve,
- (i) riser,
- (j) discharge pipe simulated by the Moody choked flow model.

A graph is required of the fraction of the initial mass still in the heated test section (calculated in SOUT and stored in AZ(2)). Some experimental results are read in and also plotted on the graph.

```

//WJT1P JOB (PIA03200,M1),W.J.TURNER,
// MSGLEVEL=(2,0),
// CLASS=B,
// TIME=2
//RR EXEC PROGEXEC,
// REGION,GO=220K,
// LIB='PHYS.WJTOWEN1'
//GO.FT02F001 DD DUMMY,DISP=OLD
//GO.FT04F001 DD DSN=GDT.BTABLE,DISP=SHR
//GO.FT09F001 DD DSN=PHYS.WJTP15,
// DISP=(,CATLG),
// SPACE=(TRK,(10,20),RLSE),
// DCB=(RECFM=VRS,BLKSIZE=1092),UNIT=SYSDA
//GO.FT10F001 DD SYSOUT=A,DCB=(RECFM=FBA,LRECL=133,BLKSIZE=1330)
//GO.SYSIN DD *
PRENOLI (1969) BURST AT RISEN OUTLET 59 MM FIXED INLET FLOW
      711. .01      8.      14.
.0001 7.      .03      2.      -2.
.04 2.      .5      9999999.
      32.      14.      -1.      .0000454      .04
      .001      1.      .189
#. .11037
1000. .11037 .189
2.6 3.3 0. 1.
      .1 1.
.0007 540. 6. .03 .100269
.332 .094 1.02 1.
.0007 1. 1. .03 .100269
.332 .094 1.02 1.
.0000454 35.67 8.623 .0328
.2319 .10307 2.71 1.
.0000454 1. .05 .0328 64.21
.2319 .10307 2.71 1.
.0000454 3.27 .164 .0328 2.2
.2319 .10307 2.71 1.
.0007 .12 .063 .0299
.2211 .09379 2.85 1.
.0007 2.64 2.64 .0299 2.
.2211 .09379 2.85 1.
.0007 5.28 5.28 .0299 2.
.2211 .09379 2.85 1.
.0007 5.28 5.28 .0299 2.
.2211 .09379 2.85 1.
.0007 .12 .064 .0299 2.
.2211 .09379 2.85 1.
.0007 21.478 1.345 .0299 2.2
.2211 .09379 2.85 1.
.0000454 5.318 .291 .0328
.2319 .10307 2.71
.0000454 .0328
.2319 .10307 2.71
P4 2661P1 2652CHASS 3653R/R0 3654ROFF 3652PC 4751
P7 2670P12 2605P9 2676P14 2691W1 2651W4 2660
W7 2669W12 2604
/*
//WJT12P EXEC RELEASE,COND=EVEN
//
//WJT12P JOB (PIA03200,M1),W.J.TURNER,
// MSGLEVEL=(2,0),
// TYPRUN=HOLD,

```

```

// CLASS=B,
// TIME=2
//AF EXEC RUFFPROG,PRG=AEPL0T
//TT EXEC PROGEXEC,
// MEM=PLOT,
// REGION,GO=180K,
// LIB='PHYS.WJTOWEN1'
//GO.FT09F001 DD DSN=PHYS.WJTP15,DISP=(OLD,DELETE)
//GO.SYSIN DD *
      1 4
0. 1. 1 9 1.
-0. -5.
-8 CISA EXPT 12
0. 1. .316 .793 .562 .75
.798 .665 1.048 .566 1.377 .521
1.746 .488 2.201 .462
      6 1 2 7 8 9
10
-8. -5. 9 1.
/*
//

```

APPENDIX B

OUTPUT FROM SAMPLE PROBLEM

PRCMBLI (1969) BURST AT RISER OUTLET 59 MM FIXED INLET FLOW

INPLT DATA ARRAY CC	2	3	4
1 0.0	711.0000	0.10000000-01	0.0
5 8.000000	14.0000	0.10000000-03	7.000000
9 0.0	10 -0.30000000-01	11 2.000000	12 -2.000000
13 0.40000000-01	14 2.000000	15 0.20000000	16 4.50000000
17 0.0	18 9999999.	19 0.0	20 32.000000
21 14.00000	22 -1.000000	23 0.84540000-03	24 0.40000000-01
25 0.0	26 0.0	27 0.0	28 0.10000000-02
29 1.000000			

28 FEB 74

PREMOLI (1969) BURST AT RISER OUTLET 59 MW FIXED INLET FLOW

TIME=	0.0	SECS	POWER(MW)*	0.189000	POWER TO COOLANT=	0.189000	VOID	FRACTION	TEMP	SAT TEMP	HEAT FLUX	
MODE	DENSITY	FLCM	QUALITY	FLOW	PSIA	STATIC ENTH	FUEL TEMP	VOID	FRACTION	TEMP	SAT TEMP	HEAT FLUX
	LB/FT3			LB/SEC		BTU/LB	DEGF			DEGF	DEGF	BTU/SEC FT
1	61.417	-0.7269	0.28696	0.28696	898.01	42.791	71.247	-0.31688	0.0	71.247	531.69	0.0
2	61.379	-0.71946	0.28696	0.28696	897.39	44.098	76.939	-0.31588	0.0	72.529	531.61	0.22743
3	49.745	-5.11854D-02	0.28696	0.28696	786.83	472.06	490.01	-3.70961D-02	0.0	486.49	516.30	0.22743
4	46.722	-5.05066D-02	0.28696	0.28696	786.43	472.45	486.82	-3.86107D-02	0.0	486.82	516.24	8.12543D-15
5	46.698	-4.84518D-02	0.28696	0.28696	778.77	472.44	486.81	-3.50775D-02	0.0	486.85	515.12	-4.06268D-15
6	46.697	-3.09353D-02	0.28696	0.28696	714.43	472.44	486.85	-2.23825D-02	0.0	486.85	505.36	4.06279D-15
7	46.697	-3.09000D-02	0.28696	0.28696	711.00	472.44	486.86	-2.17196D-02	0.0	486.86	504.82	0.0
8	46.654	-2.87492D-02	0.28696	0.28696	710.95	473.32	543.73	-2.07780D-02	0.0	487.60	504.48	4.1918
9	31.813	2.62735D-02	0.28696	0.28696	708.40	505.59	543.43	0.35803	0.0	504.48	504.48	4.1918
10	15.924	0.13074	0.28696	0.28696	701.65	539.30	536.50	0.69519	0.0	503.34	503.34	4.1918
11	11.133	0.26754	0.28696	0.28696	690.32	566.38	530.48	0.79664	0.0	501.53	501.53	4.1918
12	11.093	0.26882	0.28696	0.28696	690.03	566.65	501.48	0.79749	0.0	501.48	501.48	-9.07103D-14
13	10.267	0.26172	0.28696	0.28696	624.86	554.33	490.59	0.81467	0.0	490.59	490.59	2.49176D-14
14	10.186	0.26302	0.28696	0.28696	618.42	553.70	489.46	0.81634	0.0	489.46	489.46	-7.01180D-15
MODE	DEZ	DE3	RM	DM1	DM2	DM3	RP	RP	RP	RH	EITP	RITP
1	0.0	0.28696	1.36078D-03	9.34468D-03	-4.99222D-09	4.41151D-08	2.28678D-04	-2.02078D-03	0.59000	0.0	0.59000	0.0
2	0.0	0.28696	1.34162D-03	9.35948D-03	-5.01349D-09	4.71836D-08	2.29368D-04	-2.15864D-03	0.50000	0.0	0.50000	0.54757
3	0.0	0.28696	1.65540D-03	1.15374D-02	-1.69420D-08	1.53755D-06	5.09105D-04	-4.68041D-02	0.50000	0.0	0.50000	0.53739
4	0.0	0.28696	1.65603D-03	1.15418D-02	-1.69795D-08	1.56012D-06	5.09842D-04	-4.68456D-02	0.50000	0.0	0.50000	0.0
5	0.0	0.28696	1.65616D-03	1.15427D-02	-1.69907D-08	1.56081D-06	5.10101D-04	-4.68592D-02	0.50000	0.0	0.50000	1.0000
6	0.0	0.28696	1.65695D-03	1.15482D-02	-1.70554D-08	1.56546D-06	5.11554D-04	-4.69542D-02	0.50000	0.0	0.50000	1.0000
7	0.0	0.28696	1.65897D-03	1.15483D-02	-1.70571D-08	1.56563D-06	5.11594D-04	-4.69579D-02	0.50000	0.0	0.50000	0.0
8	0.0	0.28696	1.65894D-03	1.15584D-02	-1.71440D-08	1.57149D-06	5.13312D-04	-4.70522D-02	0.73816	0.62886	0.73816	0.62886
9	-4.10870D-02	0.28696	1.65894D-03	1.81239D-02	-1.67579D-05	9.70964D-05	9.17857	-0.89500	0.41420	0.57748	0.41420	0.57748
10	-3.10470	0.48778	5.57710D-03	3.88742E-02	-2.38574D-05	9.70907D-05	5.42040D-02	-0.22626	0.24943	0.45999	0.24943	0.45999
11	-0.13237	0.7269	8.06631D-03	6.20731D-02	-3.57406D-05	1.23678D-04	3.19926D-02	-0.11142	0.0	0.45999	0.0	0.45999
12	-0.13257	0.77255	8.06631D-03	6.23834E-02	-3.59032D-05	1.23989D-04	3.18319D-02	-0.11142	0.0	0.52689	0.52689	0.52689
13	-0.15256	0.87802	1.09984D-02	6.97541D-02	-4.51137D-05	1.43594D-04	3.22994D-02	-0.10416	0.0	0.32697	0.32697	0.32697
14	-0.15477	0.81170	1.01256D-02	7.05713D-02	-4.61954D-05	1.45801D-04	3.23559D-02	-0.10346	0.0	0.0	0.0	0.0
MODE	E	EK	F	TMT	HT	GS	HSL	HSL	HSL	HSV	RSL	RSV
1	12.279	1.78215D-07	3.99744D-05	1.75163D-75	1.00000D-01	3.96682D-04	526.38	526.38	1196.5	1196.5	47.118	1.9915
2	12.654	5143.4	903.32	1.0000	0.54985	5476.0	526.28	526.28	1196.5	1196.5	47.123	1.9901
3	135.46	13921.5	953.35	1.0000	0.68675	6756.7	507.49	507.49	1199.7	1199.7	48.027	1.7267
4	135.53	7908.5	608.16	1.0000	0.69343	4634.2	507.42	507.42	1199.8	1199.8	48.030	1.7257
5	135.57	7904.8	609.21	1.0000	0.69343	4634.6	506.06	506.06	1200.0	1200.0	48.035	1.7078
6	135.57	7917.3	608.48	1.0000	0.69344	4636.8	498.32	498.32	1201.5	1201.5	48.644	1.5585
7	135.57	13947.0	957.81	1.0000	0.79577	6763.1	493.68	493.68	1201.6	1201.6	48.674	1.5507
8	135.82	13971.0	958.32	15.000	0.79616	6769.0	493.67	493.67	1201.6	1201.6	48.674	1.5505
9	146.89	34536.0	2509.4	23.000	1.1474	17851.0	493.27	493.27	1201.6	1201.6	48.693	1.5456
10	159.02	1.78271E+05	6221.5	42.000	1.3478	44237.0	491.91	491.91	1201.8	1201.8	48.756	1.5292
11	191.14	5.12870E+05	10233.0	42.000	1.5435	72698.0	489.75	489.75	1202.0	1202.0	48.855	1.5033
12	191.36	5.18698E+05	10283.0	4.0000	1.3088	73060.0	489.70	489.70	1202.1	1202.1	48.858	1.5026
13	191.39	3.81447E+05	5529.0	42.000	1.4177	47954.0	475.83	475.83	1203.4	1203.4	49.443	1.3549
14	191.39	3.51860E+05	5606.8	4.0000	1.1968	48648.0	475.51	475.51	1203.4	1203.4	49.503	1.3404
TIME	P4	P1	CMASS	R/PO	ROFF	PC	P7	P7	P7	P12	P9	
0.0	785.427	808.011	97.9976	1.00000	0.0	0.189000	711.000	711.000	690.029	708.804		
1.00000D-02	786.629	898.011	98.1572	1.00165	3.539017E-14	0.189862	711.004	711.004	696.613	708.864		
2.00000D-02	785.501	848.011	98.5975	1.00612	1.236764E-13	0.218667	711.176	711.176	707.771	709.404		
3.00000D-02	787.057	908.011	99.2049	1.01232	8.463439E-14	0.225237	712.359	712.359	715.242	711.372		
TIME STEP MULTIPLIER BY	0.50000											
3.50000E-02	787.546	908.011	99.5177	1.01551	2.790835E-13	0.231907	713.932	713.932	717.032	713.060		

4.3777777777777777	782.588	898.010	1.01899	4.127511F-14	0.241777	715.510	719.472	715.189
4.5000000000000000	780.587	898.009	1.02283	5.986429E-13	0.252552	717.874	722.517	717.559
5.0000000000000000	791.899	898.008	1.02679	3.082959E-10	0.261034	720.273	722.998	719.994
TIME STEP MULTIPLIED BY								
2.00000								
5.0000000000000000	794.064	898.007	1.03019	1.734829E-11	0.264728	723.749	721.291	723.227
6.5000000000000000	800.010	898.004	1.03438	1.350335E-12	0.273388	730.936	714.151	729.825
7.5000000000000000	805.384	898.001	1.03583	6.928975E-13	0.279665	736.438	708.496	734.263
TIME STEP MULTIPLIED BY								
2.00000								
8.5000000000000000	803.811	897.999	1.03510	9.361269E-12	0.278799	739.516	705.299	736.619
0.105000	803.842	898.000	1.02991	1.571643E-13	0.282553	740.382	703.333	736.934
TIME STEP MULTIPLIED BY								
2.00000								
0.125000	801.720	898.006	1.02224	4.315870E-14	0.270052	731.759	697.499	729.125
0.135000	787.054	898.009	0.998526	1.015450E-13	0.271941	715.716	678.020	713.038
0.205000	774.566	898.012	0.971823	9.405893E-14	0.275693	699.833	659.478	696.969
0.245000	766.000	898.040	0.944581	1.698383E-13	0.274153	686.836	644.581	683.692
0.245000	757.016	898.044	0.919343	2.072948E-13	0.273557	671.809	630.655	668.658
0.345000	753.640	898.040	0.897217	3.334463E-13	0.265731	662.800	620.522	659.291
TIME STEP MULTIPLIED BY								
2.00000								
0.365000	746.496	898.031	0.879586	1.499197E-13	0.251537	648.411	608.019	645.008
0.445000	742.400	897.975	0.852690	3.930378E-13	0.241988	635.063	591.507	631.138
0.535000	732.740	897.883	0.831932	2.924748E-11	0.242429	611.861	568.503	607.855
TIME STEP MULTIPLIED BY								
2.00000								
0.635000	734.024	897.743	0.827953	1.501472E-10	0.220161	604.423	558.012	599.853
0.745000	729.832	897.513	0.822393	2.701158E-13	0.216691	588.783	539.369	583.563
0.945000	735.042	897.726	0.814569	1.526333E-12	0.206693	579.617	528.621	574.963
1.045000	728.630	895.991	0.829332	2.438474E-12	0.207487	568.940	518.498	563.301
1.245000	731.127	895.128	0.812507	6.525237E-12	0.201315	561.742	511.730	556.081
1.485000	731.463	894.152	0.818416	1.247116E-11	0.205601	553.110	503.995	547.581
1.745000	735.044	893.045	0.798689	1.034285E-11	0.201447	548.308	499.454	542.771
1.725000	730.389	891.874	0.803654	9.042408E-12	0.205330	541.728	493.461	536.283
1.885000	733.521	890.581	0.787920	7.329057E-12	0.202229	538.229	486.966	532.757
2.045000	729.379	889.193	0.792907	4.776293E-12	0.205715	533.130	485.188	527.719

28 FEB 74

PREMOLI (1969) BURST AT RISER OUTLET 59 MW FIXED INLET FLOW

TIME= 2.20500 SECS POWER(MW)= 0.18900C POWER TO COOLANT= 0.20285S

NODE	DENSITY LB/FT3	FLOW GAL/SEC	PSIA	STATIC ENTH BTU/LB	FUEL TEMP DEGF	VCIO FRACTION	TEMP DEGF	SAT TEMP DEGF	HEAT FLUX BTU/SEC FT
1	61.418	0.28696	897.71	42.791	71.281	-0.31432	71.286	530.33	0.0
2	61.995	0.29604	887.04	23.154	66.187	-0.32689	52.063	530.24	0.544413
3	48.730	0.38545	732.71	492.67	501.92	-5.18900D-03	503.94	508.20	-4.64570D-02
4	48.792	0.38683	732.30	491.36	498.96	-6.44955D-03	502.84	508.13	-0.21162
5	50.057	0.39417	721.27	464.62	480.60	-3.13148D-02	480.07	506.43	2.85234D-02
6	41.533	0.39471	544.56	463.48	479.46	0.17681	475.91	475.91	0.14089
7	36.123	0.39705	530.74	463.48	475.52	0.28895	473.20	473.20	0.19967
8	35.427	0.39713	530.58	463.92	512.92	0.30312	473.17	473.17	3.4698
9	22.977	0.40430	525.30	477.50	508.96	0.57481	472.12	472.12	4.3027
10	14.328	0.41570	509.24	493.96	498.24	0.73248	468.90	468.90	4.2826
11	10.572	0.41604	493.43	506.06	490.29	0.80852	463.55	463.55	4.2833
12	10.533	0.41606	482.75	506.11	484.69	0.80931	463.40	463.40	9.35239D-02
13	5.2603	0.31143	190.38	412.21	378.93	0.91052	377.69	377.69	0.10255
14	3.3894	0.35900	99.218	353.39	375.32	0.94364	327.25	327.25	0.56482

NODE	DF2	DE3	RM	DM1	DM2	DM3	RP	RH	EITP	RITP
1	0.0	0.28696	1.3476D-03	9.3445D-03	-4.9903D-09	4.4103D-08	2.29607D-04	-2.0204D-03	0.50000	0.57969
2	0.0	0.29504	1.4136D-03	9.5505D-03	-4.9759D-09	-1.2382D-09	2.1821D-04	5.4302D-05	0.50000	0.61126
3	0.0	0.38645	3.7648D-03	1.5861D-02	-3.4645D-08	3.0881D-06	5.0839D-04	-4.9100D-02	0.50000	0.61690
4	0.0	0.38683	3.0668D-03	1.5856D-02	-3.4457D-08	3.0772D-08	5.4821D-04	-4.8958D-02	0.50000	0.47897
5	0.0	0.39417	3.0039D-03	1.5749D-02	-3.0831D-08	2.8630D-06	4.9818D-04	-4.6172D-02	0.92208	0.90551
6	-4.6997D-02	0.54412	3.7421D-03	1.9028D-02	-3.9463D-05	1.7512D-04	0.43345	-1.9236	0.51332	0.51981
7	-6.3567D-02	0.67007	4.3741D-03	2.2033D-02	-4.2886D-05	1.8289D-04	0.34914	-1.4886	0.50588	0.57679
8	-6.6191D-02	0.68276	4.6339D-03	2.2478D-02	-4.5317D-05	2.1145D-04	0.33710	-0.4322	0.48128	0.55183
9	-0.16740	0.98797	7.5896D-03	3.7544D-02	-5.5317D-05	2.1145D-04	0.14687	-0.56110	0.45812	0.54077
10	-0.24118	1.2654	1.3372D-02	6.4395D-02	-8.7964D-05	2.9129D-04	7.3476D-02	-0.24287	0.42972	0.51995
11	-0.30064	1.3536	2.0279D-02	8.7494D-02	-1.3298D-04	3.8137D-04	4.8146D-02	-0.13841	0.23479	0.40245
12	-0.30152	1.3545	2.0389D-02	8.8011D-02	-1.3383D-04	3.8339D-04	4.7955D-02	-0.13756	0.71208	0.75255
13	-1.0929	1.9768	5.1518D-02	0.29457	-1.1377D-03	1.5269D-03	5.8033D-02	-7.7797D-02	0.64075	0.71752
14	-2.0758	2.6179	0.12389	0.61125	-4.6118D-03	3.6926D-03	7.1928D-02	-5.7453D-02	0.0	0.0

NODE	E	FK	F	THF	HT	GS	MSL	MSV	RSL	RSV
1	12.279	1.7820D-07	3.9370D-05	1.7516D-75	1.0000D-01	3.9667D-04	524.69	1196.8	47.200	1.9666
2	6.354	9840.6	954.48	1.0000	0.55129	5770.1	524.58	1196.8	47.206	1.9650
3	190.39	35431.	1670.4	1.0000	0.79822	12509.	497.72	1201.1	48.486	1.6007
4	15.007	26120.	1044.9	1.0000	0.83203	8582.0	497.64	1201.1	48.690	1.5997
5	133.14	20226.	1084.3	1.0000	0.82879	8686.0	495.60	1201.4	49.585	1.5783
6	183.54	28584.	2057.2	22.000	0.52436	16455.	459.75	1204.3	50.201	1.1762
7	185.33	70438.	3955.7	22.000	0.79353	29282.	456.63	1204.5	50.337	1.1455
8	185.73	73410.	4063.3	22.000	0.81633	30079.	456.59	1204.5	50.339	1.1451
9	201.17	2.1639D+05	7754.7	42.000	1.2749	57541.	455.39	1204.5	50.391	1.1335
10	229.10	7.47E+05	15126.	42.000	1.5943	1.1259D+05	451.69	1204.6	50.552	1.0983
11	252.42	1.9513D+06	24327.	42.000	1.7540	1.8073D+05	445.57	1204.7	50.814	1.0418
12	252.60	1.8754D+06	24478.	42.000	1.7557	1.8184D+05	445.41	1204.7	50.821	1.0404
13	256.77	1.3664D+07	40018.	42.000	1.8162	3.7081D+05	351.11	1197.7	54.546	0.41694
14	257.73	6.3699D+07	81750.	342.00	0.12373	7.8777D+05	297.95	1187.0	56.391	0.22401

TIME	P4	P1	CMASS	R/R0	PUFF	PC	P7	P12	P9
2.20500	734.206	887.712	76.3727	0.779332	6.737898E-12	0.202855	526.738	482.751	525.304
2.35500	723.344	886.144	76.7783	0.783471	1.346902E-11	0.205773	526.755	479.010	521.368
2.52500	757.951	884.401	75.5143	0.770573	1.655998E-11	0.203009	525.140	477.326	519.730
2.59500	727.643	882.740	75.7568	0.773048	1.018750E-11	0.207781	521.978	474.365	516.608

TIME STEP MULTIPLIED BY 2.0000

TIME	STEP MULTIPLIED BY	P4	P1	CMASS	R/R0	PUFF	PC	P7	P12	P9
2.34500	733.008	881.054	74.5466	0.757699	1.590572E-12	0.208341	523.168	473.407	515.768	

TIME STEP MULTIPLIED BY	2.0000								
3.16500	977.781	73.9227	0.754332	9.024290E-12	0.175764	521.107	472.983	515.719	
3.80500	870.442	67.5185	0.688981	6.328992E-12	0.190189	518.798	471.137	513.355	
4.44500	862.924	62.3418	0.636157	1.134656E-11	0.194422	505.479	458.413	500.108	
5.08500	855.307	57.8129	0.589942	8.829193E-12	0.209311	494.470	447.845	489.129	
TIME STEP MULTIPLIED BY	2.0000	54.8862	0.560078	7.392280E-12	0.199944	484.929	437.728	479.570	
5.72500	847.859								
TIME STEP MULTIPLIED BY	2.0000	50.4228	0.514531	3.469771E-10	0.200699	473.360	426.210	468.039	
7.70500	833.119								

28 FEB 74

PREMULI (1969) BURST AT RISER OUTLET 59 MW FIXED INLET FLOW

TIME= 5.56500 SECS POWER(MW)= 0.18900C POWER TO COOLANT= 0.20392D

NODE	DENSITY LB/FT3	FLOW LB/SEC	QUALITY	FLOW LB/SEC	PRESSURE PSIA	STATIC ENTH BTU/LB	FUEL TEMP DEGF	VOID FRACTION	TEMP DEGF	SAT TEMP DEGF	HEAT FLUX BTU/SEC FT
1	61.427	0.28696	-0.67912	0.28696	804.14	42.791	71.494	-0.29370	71.604	518.80	0.0
2	62.271	0.29805	-0.72160	0.29805	803.46	13.322	46.632	-0.31184	42.703	518.70	0.22279
3	46.786	0.27828	1.002680-03	0.27828	677.450	487.85	504.01	2.48893D-02	499.47	499.47	0.27391
4	46.595	0.27822	2.06257D-03	0.27822	677.16	488.35	504.01	5.00309D-02	499.39	499.39	3.08082D-02
5	41.538	0.27788	7.57226D-03	0.27788	668.65	489.52	499.39	0.15778	498.00	498.00	1.06717D-02
6	18.386	0.27824	7.62585D-02	0.27824	664.25	458.00	465.42	0.65238	459.43	459.43	5.14693D-02
7	17.376	0.28031	8.34548D-02	0.28031	664.96	464.96	461.59	0.67430	455.33	455.33	4.73393D-02
8	17.198	0.28038	8.46800D-02	0.28038	665.61	465.23	468.31	0.67687	455.30	455.30	4.2438
9	13.285	0.28160	8.13628	0.28160	440.47	492.13	492.13	0.75483	454.13	454.13	4.2380
10	9.2805	0.28323	0.28323	0.28323	424.78	494.46	484.57	0.83441	450.52	450.52	4.2388
11	6.9627	0.33624	0.33624	0.33624	400.10	509.17	476.93	0.88002	444.62	444.62	5.03284D-02
12	6.9371	0.33745	0.33745	0.33745	399.45	509.24	480.99	0.88051	444.46	444.46	4.69857D-02
13	3.4888	0.43746	0.43746	0.43746	413.85	366.46	366.46	0.94277	361.17	361.17	2.79235D-02
14	2.3519	0.28887	0.28887	0.28887	188.344	322.11	322.11	0.96191	318.97	318.97	

NODE	DE2	DM3	DM2	DM1	DM3	DM2	DM1	RP	RH	EITP	RITP
1	0.28596	1.34058D-03	8.34225D-03	4.51244D-09	4.51244D-08	4.51244D-08	4.51244D-08	2.28023D-04	-2.02182D-03	0.50000	0.59784
2	0.29805	1.42855D-03	9.57280D-03	4.87799D-09	-2.57446D-08	-2.57446D-08	-2.57446D-08	2.12932D-04	1.12379D-03	1.00000	0.88891
3	0.36615	1.62072D-03	1.16480D-02	-1.36121D-05	7.14871D-05	7.14871D-05	7.14871D-05	0.40005	-2.1011	0.49980	0.50127
4	0.37150	1.66153D-03	1.19439D-02	-1.36819D-05	7.15384D-05	7.15384D-05	7.15384D-05	0.38225	-1.9988	0.51381	0.54142
5	0.42150	1.86032D-03	1.33896D-02	-1.42412D-05	7.24474D-05	7.24474D-05	7.24474D-05	0.31572	-1.6063	0.63066	0.64232
6	0.78761	4.41529D-03	3.17372D-02	-3.57134D-05	1.20943D-04	1.20943D-04	1.20943D-04	0.12834	-0.43379	0.50812	0.53090
7	0.82950	4.82800D-03	3.42880D-02	-4.03588D-05	1.31331D-04	1.31331D-04	1.31331D-04	0.12300	-0.39923	0.50727	0.50954
8	0.82929	4.84774D-03	3.45803D-02	-4.06583D-05	1.32063D-04	1.32063D-04	1.32063D-04	0.12151	-0.39347	0.45237	0.51810
9	0.92929	6.76359D-03	4.80367D-02	-5.37171D-05	1.60325D-04	1.60325D-04	1.60325D-04	7.97867D-02	-0.23719	0.34788	0.50771
10	0.57111	1.13864D-02	8.04044D-02	-8.61550D-05	2.19598D-04	2.19598D-04	2.19598D-04	4.68150D-02	-0.11951	1.00000	0.50080
11	0.93066	1.79046D-02	0.12521	-1.28928D-04	2.75888D-04	2.75888D-04	2.75888D-04	3.24591D-02	-7.09080D-02	1.00000	0.83079
12	0.52872	1.79093D-02	0.12593	-1.29728D-04	2.76821D-04	2.76821D-04	2.76821D-04	3.23467D-02	-7.04802D-02	0.68715	0.73493
13	1.2018	5.74817D-02	0.39891	-1.04949D-03	1.01443D-03	1.01443D-03	1.01443D-03	4.07913D-02	-0.07724D-02	0.60624	0.67480
14	1.3444	0.10996	0.76133	-3.35355D-03	2.03919D-03	2.03919D-03	2.03919D-03	4.80060D-02	-3.06609D-02	0.0	0.0

NODE	F	FK	F	TMT	HT	GS	HSL	HSV	RSL	RSV
1	12.279	1.78160D-07	3.99378D-05	1.75163D-79	1.00000D-04	3.96621D-04	510.53	1199.3	47.883	1.7674
2	2.9707	5955.7	966.10	1.00000	0.54999	5822.6	510.41	1199.3	47.888	1.7658
3	135.81	13761.	1124.8	23.900	0.31887	7974.7	487.30	1202.3	48.968	1.4743
4	135.97	8212.0	846.48	22.000	0.22775	6405.3	487.22	1202.3	48.971	1.4733
5	136.44	10317.	1101.2	22.000	0.41345	8385.6	485.57	1202.5	49.047	1.4539
6	138.88	13251.	358.7	42.000	1.0540	26660.	440.89	1204.8	51.013	1.0001
7	140.26	13392D+05	6014.5	42.000	1.2313	41751.	436.25	1204.8	51.209	0.96005
8	140.55	1.36211D+05	6082.5	42.000	1.2365	42223.	436.20	1204.8	51.211	0.95968
9	152.01	2.93669D+05	9032.4	42.000	1.4051	62720.	434.89	1204.8	51.266	0.94857
10	173.97	9.50614D+05	15980.	42.000	1.5941	1.10901D+05	430.82	1204.7	51.437	0.91468
11	195.27	2.47221D+06	24772.	42.000	1.7068	1.71641D+05	424.19	1204.6	51.712	0.86157
12	195.52	2.50221D+06	24914.	42.000	1.7068	1.72615D+05	424.02	1204.6	51.719	0.86019
13	204.68	1.62152D+07	32631.	42.000	1.7348	2.98827D+05	393.54	1194.6	55.175	0.34296
14	207.71	5.97155D+07	60409.	32.000	1.8253	5.60359D+05	289.33	1184.9	56.672	0.20072

TIME 9.56500 P4 677.143 P5 904.138 P6 45.7678 P7 445.777 P8 399.453 P9 440.472

28 FEB 74

PPEMCLI (1969) BURST AT RISER OUTLET 59 MW FIXED INLET FLOW

TIME	P14	W1	W4	W7	W12
0.0	618.419	0.286962	0.286962	0.286962	0.286962
1.000000D-02	563.211	0.286962	0.286962	0.286959	0.251346
2.000000D-02	481.408	0.286962	0.286962	0.286962	0.188991
3.000000D-02	394.918	0.286962	0.286962	0.285976	0.751957
3.500000D-02	337.459	0.286962	0.286743	0.286869	0.146128
4.000000D-02	317.875	0.286962	0.286536	0.283935	0.133355
4.500000D-02	269.954	0.286962	0.286277	0.282538	0.117621
5.000000D-02	257.330	0.286962	0.285888	0.279436	0.116772
5.500000D-02	248.125	0.286962	0.285416	0.278654	0.132037
6.500000D-02	237.270	0.286962	0.283768	0.276721	0.191912
7.500000D-02	227.371	0.286962	0.281716	0.275781	0.248741
8.500000D-02	219.254	0.286962	0.279498	0.275856	0.303699
0.105000	208.907	0.286962	0.275217	0.274803	0.343809
0.125000	201.819	0.286962	0.272065	0.277180	0.364229
0.165000	195.617	0.286962	0.274927	0.279089	0.412882
0.205000	185.284	0.286962	0.282013	0.285025	0.434197
0.245000	175.598	0.286962	0.291684	0.294439	0.446416
0.295000	167.052	0.286962	0.303723	0.306687	0.444416
0.325000	160.057	0.286962	0.315160	0.316509	0.438285
0.365000	154.210	0.286962	0.327860	0.330412	0.425929
0.445000	145.394	0.286962	0.349909	0.350882	0.425100
0.525000	138.400	0.286962	0.371979	0.373915	0.401263
0.605000	133.049	0.286962	0.384258	0.384560	0.406205
0.745000	123.881	0.286962	0.403435	0.405016	0.417009
0.925000	117.122	0.286962	0.402643	0.413831	0.426063
1.08500	111.540	0.286962	0.409065	0.448033	0.429885
1.24500	108.071	0.286962	0.397840	0.410656	0.434632
1.40500	105.623	0.286962	0.400322	0.437367	0.430338
1.58500	101.864	0.286962	0.393737	0.402092	0.429958
1.72500	102.272	0.286962	0.396365	0.428536	0.422647
1.88500	101.185	0.286962	0.390792	0.399868	0.421918
2.04500	100.103	0.286962	0.392275	0.422173	0.416673
2.20500	99.2183	0.286962	0.386828	0.397047	0.416062
2.36500	98.3606	0.286962	0.387381	0.416587	0.411535
2.52500	97.6529	0.286962	0.381793	0.393755	0.411796
2.68500	96.5727	0.286962	0.381553	0.410664	0.407791
2.84500	96.4409	0.286962	0.375892	0.391172	0.408644
3.00500	95.7276	0.286962	0.369338	0.404827	0.409570
3.16500	95.6109	0.286962	0.340349	0.377417	0.402225
4.44500	97.1898	0.286962	0.325716	0.354336	0.373425
5.18500	96.3743	0.286962	0.314905	0.335015	0.351611
5.72500	93.5229	0.286962	0.306733	0.322498	0.332986
7.07500	91.8044	0.286962	0.298619	0.309180	0.317446
9.56500	89.3437	0.286962	0.278220	0.280308	0.284427

