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**GYMEA — A NUCLIDE DEPLETION, SPACE INDEPENDENT, MULTIGROUP
NEUTRON DIFFUSION, DATA PREPARATION CODE**

by

**J.P. POLLARD
G.S. ROBINSON**

March 1966

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ABSTRACT

GYMEA is a multi-purpose neutronics code used extensively in the H.T.G.C.R. project of the A.A.E.C. The code has three main functions:

- (i) running of burnup surveys of near-homogeneous reactor systems,
- (ii) running of neutron diffusion calculations in nuclear data library correlations with measurements made in near-homogeneous subcritical assemblies carried out at the A.A.E.C., and
- (iii) preparation of suitably averaged cross section data from the code's 120-group (with resonance parameters), 70-nuclide main data library for direct input to space-dependent codes such as CRAM, 4-ZOOM (9-ZOOM translated to the IBM 7040), DSN and TDC (Carlson's FLOCO versions translated to the IBM 7040), and WDSN.

Because free input is used and a toy compiler is provided, new applications of the code can be developed by code users to extend its function. The source language is FORTRAN IV (95 per cent.) and MAP (5 per cent.) for the IBM 7040 (and later an IBM 360). 32K words of core storage are required together with 4-6 magnetic tapes and an on-line 1401 including reader, punch, and printer. Typical runs for data preparation take 3 minutes, diffusion calculations 10 minutes and burnup calculations 15 minutes.

A few illustrative examples are included in the report.

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1. INTRODUCTION

The maintenance of a nuclear data file, and its ready transformation to group data for neutronics codes, is essential for a reactor project such as the H.T.G.C.R. (High Temperature Gas-Cooled Reactor) project under study at the A.A.E.C. Research Establishment. The file should be processed with groups collapsed over a reasonable spectrum. For a homogeneous core the flux obtained from a space-independent, or space-averaged, calculation of neutron balance provides a reasonable spectrum for group collapsing. In a reactor study extending beyond clean core conditions the calculation of burnt-up core nuclide concentrations is a valuable asset for determining group data at different stages of reactor life. In addition no reactor study is satisfactory unless the data file correlates satisfactorily with selected experiments carried out on macroscopic properties of materials to be used in the reactor project. Out of these needs came the codes GIBBA (Pollard and Clancy 1964) and MULGA (Clancy et al. 1963); and out of experience obtained with these codes came GYMEA.

GYMEA is essentially three codes in one: it is (i) a zero dimensional multigroup (λ 120) neutron diffusion code, (ii) a multinuclide (λ 100) burnup code, and (iii) a data preparation code. It is written in FORTRAN IV (95%) and MAP (5%) for the IBM 7040 computer system. The programme computes neutron spectrum (optionally) and nuclide concentrations to enable spectrum-weighted, group-averaged cross section data and scattering transfer matrix data to be prepared for direct input to a variety of neutronics codes (CRAM, 4-ZOOM, DSN, WDSN, and GYMEA).

Although the code has essentially the three main functions mentioned previously, these functions are by no means independent and experience with the code has shown the advantages to be gained from having all three aspects in the one programme. Features of the code include:

- (1) criticality on concentration or group-dependent buckling;
- (2) output of reaction rates, microscopic and/or macroscopic cross sections collapsed over the spectrum for mixtures defined by the user;
- (3) saving up of the previously mentioned information, which may be processed further using the toy-FORTRAN provided as part of the code;
- (4) detailed resonance theory using an extension of the Hill and Schaefer method (McKay et al. 1965);
- (5) allowance for resonance overlap following Keane's (1966) observations on the work of Rowlands (1963), including gross flux corrections suggested by Keane and Pollard (1966);
- (6) allowance for resonance screening of small particles of fuel separated from the main bulk of core material, following Keane (1964a);
- (7) production of resonance shielding factors to save time when using the code for a burnup survey;
- (8) burnup calculated analytically— permitting large time steps;
- (9) parallel searching of a tape for a chosen scattering matrix using the IBM 1401 while the IBM 7040 is engaged in calculation (Richardson 1965).

The code uses two completely binary libraries— one a cross section (and resonance parameter) tape, the other a scattering matrix tape— and needs a scratch tape for intermediate storage of information. Input is read from the card reader and output is produced on the printer and punch (CRAM data, etc.). A condensed and/or mixed GYMEA library can also be produced on one or two other tape units.

The cross section library contains information for six neutron reactions:

- (1) (n, γ) , (n, α) , or (n, p) , depending on the nuclide,
- (2) (n, f) ,
- (3) $(n, \text{absorption})$,
- (4) $(n, \nu-f)$ or (n, n) for moderator nuclides,
- (5) $(n, \text{transport})$, and
- (6) $(n, 2n)$ or other reaction, depending on the nuclide.

In burnup calculations reactions (1) and (6) are used to produce other nuclides and reaction (2) is used to produce fission products for fuel of one of four different types. A general table is given at the beginning of this library which states the burnup mechanism of each nuclide. This feature adds considerably to the flexibility of the programme as burnup mechanisms are not built into the programme. Such a table may contain an entry that reads

U235,2,NO0,U236,NO0,0.,0.,0.,0.,0.

which states that U^{235} yields fission products based on the 2nd yield entry for each fission product nuclide, that decay is not to be considered. (NO0), that reaction (1) of the library produces U^{236} , that reaction (6) of the library is not to be considered as far as producing another nuclide is concerned, and that the decay constant is zero, as are also the four possible yields from fuels designated as fuels of type 1,2,3, or 4. A restriction does apply however to the ordering of the library nuclides. Further details are given in Section 18.

The scattering matrix library contains scattering transfer matrices for materials such as BeO at a variety of temperatures. In addition, the theory used in the preparation of the scattering matrix can be selected, for example BeO crystal scattering. Since the cross section library usually contains data for individual nuclides, the user must provide the mixing rules as part of the card input to the programme.

Such a rule can be punched in free format (22)

'BEO' = 1.BE9, 1.016, DENSITY 1.8'

After the appearance of this card the user can refer to BeO as though data for it were given as part of the cross section library information. The density is used only if nuclide ratios are to be normalized to unit volume or if the effect of dispersed fuel particles is to be included (5.3).

Further GYMEA libraries can be created during running of the programme. These libraries can have resonance parameters turned into shielding coefficients, or simply cross sections, and the groups can be collapsed and nuclides mixed to form a new library oriented to the particular needs of the user.

2. CODE PHILOSOPHY

The structure of an input deck for GYMEA is not unlike a set of source decks for compilation and execution under the FORTRAN monitor. The reason for this is that particular users will bias the running of GYMEA to suit their needs, which may include minor extensions of the code using the toy-FORTRAN compiler. A run may thus consist of a set of subroutines and perhaps a set of jobs. The jobs may well be carried out for different people, with completely independent jobs being run to take advantage of time-saving resulting from loading and unloading tapes.

Each GYMEA user deserves to be protected from a previous 'foul-up' just as each job of a monitor run is guaranteed this protection. In fact the philosophy of GYMEA on detecting an error is to call standard error routines which essentially call in the next job. It is possible however for any user to write his own version of error routines using the toy-FORTRAN provided as part of the input facility. This type of variable extent input is possible as the code has a standard state, and entry of data simply overwrites, or extends, this state to suit the needs of the user. It is felt that the structure and philosophy of this code with its free input, variable extent data will find wider application in scientific computing in the near future.

Those using computer programmes must surely be amazed at hardware development of computers over the past decade and must surely be disappointed with software developments over the same period. Perhaps writers of codes have available many more powerful tools now than ten years ago, but we should seriously ask, how have the code users fared? Most scientific computing in the field of nuclear reactor development is now initiated by persons other than code authors. Code authors, and machine manufacturers, must therefore seriously look at this sadly neglected area of computing. It is pleasing to read reports of a compiler language being developed by IBM (1965) which promises to FORGO (perhaps a suitable name) the old tradition in computing of terminating a user's run without his permission.

CODE MECHANISM 3

3. CODE MECHANISM

3.0 Summary

The code is a chain programme and consists of six links. Common subroutines form the bulk of the main programme and common data consumes a large fraction of the available storage. One array WH (10000) is used for storing different quantities at different stages of the calculation and an understanding of how this array is used throws some light on the structure of GYMEA.

3.1 Chain 1- Compiler and Error Routines

All data up to and including a CEASE card is read into the array WH. Since cards are originally read with FORMAT (12A6), 800 cards can be entered. These cards are then inspected, using a free FORMAT input subroutine (22), and certain cards are compiled, that is, overwritten by direct linkages to internal subroutines; for example

$$Z(8) = Z(10)*0.8234E-2$$

is compiled into a card with a \$ in column 1 followed by numbers which provide the necessary direct calculation of Z(8). The cards are then written at the beginning of tape 4 as a complete file.

In addition, errors call this chain, which rewinds all libraries and prints self-explanatory messages.

3.2 Chain 2- Input Data

The cards previously written on tape 4 are read into WH and the cards are inspected. Data is then entered into COMMON to enable the neutronics part of the programme to function. Control is transferred to chain 3 when a time card is detected; such a card may read

$$T = 0. \text{ DAYS}, 20 \text{ MAX.FLUX IT.}, 0 \text{ CRIT. IT.}$$

3.3 Chain 3- Burnup

Should burnup be required, a 100 x 100 (maximum) nuclide balance array is stored in WH. As the array must be nearly triangular (6) the original information is stored in the lower half of the array and the solution is produced in the upper half of the array. In practice the use of this array and the flexible burnup mechanism table mentioned in Section 1 makes analytic solution feasible and worthwhile.

3.4 Chain 4- Group Flux Calculation, Condensing and Mixing Functions

When group flux is required to be calculated from neutron balance, WH is used to store information as an array 83 x 120 (maximum).

Scattering from one group to another is thus restricted to 80 groups, as the first two entries in each column are 'length of scattering vector' and 'position of self-scatter term'.

Condensed and mixed cross section data is then generated in WH for six reactions per group per mixture. Generated data is thus restricted by the requirement:

$$6 \times (\text{number of condensed groups}) \times (\text{number of specified materials}) \leq 10000$$

Should the user wish to delay output of this information it may be written on tape 4 as a further file.

3.5 Chain 5- Output for Most Neutronics Codes and Printed Output

The array WH is not disturbed by this chain, which uses the information for output either

as cards for a variety of codes or as printed output. Scattering matrix information is stored in full in an array WSS (3000), hence the following restriction applies:

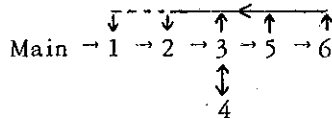
$$(\text{number of condensed groups})^2 \times (\text{number of specified materials with scattering matrices}) \leq 3000$$

3.6 Chain 6- Output for GYMEA Library

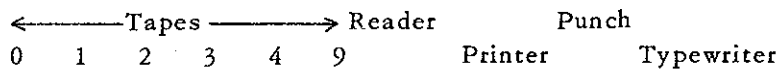
Again the array WH is left undisturbed as some of the information is needed for creating a new GYMEA library. Information not given in this array, such as resonance parameters or shielding coefficients, is obtained from the cross section library. Scattering matrices given as part of the scattering matrix library are simply condensed- mixing is ignored for this application. By repeated calling to this chain, it is possible to collapse the scattering matrices over different spectra for different temperatures and so build up a tape with the same layout as the original.

3.7 Programme Flow

Essentially each chain is called as follows:



3.8 IBM 7040 Operating System



Tape 0 = a possible library tape either used or created,

Tape 1 = " " " " " " " " " ,

Tape 2 = scattering matrix tape attached through on-line 1401,

Tape 3 = a possible library tape either used or created,

Tape 4 = scratch tape,

Tape 9 = normal system tape using \$EXECUTE GYMEA to call in programme (an ordinary chain tape may be used alternatively).

Reader = card data input unit,

Printer = printed output unit,

Punch = punched card output unit for CRAM data, etc,

Typewriter = unit used for operator messages.

3.9 Proposed IBM 360 Version

At the time of writing this report details have not been settled for the IBM 360 version of GYMEA. Magnetic disk storage will be used for the libraries, otherwise the programme will not be changed a great deal initially. In addition the output medium for other neutronics codes (CRAM, etc.) will be disk rather than the cards used for the present version. The programme should be operational in September, 1966.

Extensions of the code will follow after September 1966 which will enable more detailed treatment to be given to heterogeneous effects. In addition code-to-code communication will be improved using magnetic disk as a common data pool between programmes so that GYMEA may be readily used in a loop embracing other neutronics codes.

**A SIMPLE JOB AS 4
ILLUSTRATION**

4 A SIMPLE JOB AS AN ILLUSTRATION

4. A SIMPLE JOB AS ILLUSTRATION

\$IBSYS ...hints are given thus

\$ID

\$JOB 2210 J.POLLARD 'GYMEA'

\$* LOAD TAPES AS INDICATED

\$PAUSE 295 ON 1, 260 ON 2(VIA 1401),SCRATCH ON 4

\$EXECUTE GYMEA

 ...the above are normal 7040 monitor cards

*GYMEA ...this card must be first data card

 ...an * in col.1 indicates a comment card

,JOB1 A SIMPLE JOB AS ILLUSTRATION

 ...the first card of a job (here called JOB1)

 ...the input data is punched anywhere on a card

CALL STAN ...restores data to the standard state

USE LIB X-S=1 OF 100 NUCLIDES AND SCAT =2

 ...normally the first 3 characters identify the

 ... card and numbers are extracted from the

 ... remaining data.

'BEO'=1.BE9,1.O16,DENSITY =1.8' '

 ...the mixture BeO is defined in terms of

 ...constituents of the X-S library:' ' means end of

 ...definition

SCAT 900. 'BEO' CRY ' '

 ...900.K scattering matrix of BeO crystal is

 ...required

N(1,100)=0., 'BEO'=4.-2,U233=4.-5' '

 ...the nuclide concentrations are set

OUTPUT 3 PRINTX PRINTS CRAM

 ...the type of output required is set

MATS 'BEO', U233, PU239' '

 ...the resonance data of Pu²³⁹ will of course be

 ...obtained for infinite dilution

GROUPS 5 1 40 50 71 90 120

T=0.DAYS, 20 MAX FLUX ITERATIONS, 0 CRITICALITY

 ...this sends control to chain 3

END ...the last card of a job

CEASE ...the last card of a set of jobs (a load)

\$IBSYS ...note:the order of the data may be varied

 ...provided the quantities defined by the user precede use

5. CHAIN 2- INPUT DATA

5.0 Alphabetic List of Card Input Types

An example of each card input type is included in the alphabetic list below. Only the underlined characters are necessary to determine the type of the card. Most of the data is punched in free format (22) – the exceptions are double underlined and b(=blank) is used in these cases to indicate the setting out. Normally the data is punched as for FORTRAN coding. Data, other than FORTRAN-like cards, can be extended beyond one card (column 1 to 72); continuation cards must then be punched in column 1 with a non-descending integer, which is tested but not read as part of the data. Brief comments, which are not to be punched, are given thus:

...comments, page(s) giving details, main page
...first.

bbbb7bbb13bbb...comment card (in column 1), 15

=

/ 99 ...numbered card to assist logic (any card except END or
= ...CEASE), 59

=GYMEA NDXD NDSC ...library check card, 15

,GR01 ...beginning of a job, 14

,DATA GR01 ...beginning of data read by user's programme, 61

,DISCARD ...beginning of other than the first load, 14

,RETAIN ...alternative to above, 14

'H2O'=2.H1,1.O16,1.' ' ...a mixing rule, 19

ACCURACY = 1.E-5 ...group flux accuracy parameter, 32, 45

ALL N = 2*.04,68*0 ...concentrations for all nuclides, 21

AVX LIMIT=0. ...sets exit limit for B2AV, 33

BURN 2*-1,1,97*0 ...modifies burnup mechanism, 33

CALL NS1 ...call to subroutine, 62, 14 ; NS1 shielding
...subroutine, 23

CEASE ...terminating card of a load, 14

CONTINUE ...end of DO loop, 60

COPY 3 NDS1 1 5*0 4*1 ...sets parameters which determine
...structure of a scattering matrix
...library to be copied, 17

CREATE 3 NDX1 6*-1 64*0 ...sets parameters which determine
...structure of a cross section
...library to be created, 16

DO 2 -1=1,5,1 ...start of a DO loop, 60

DT= 100.,20,0,2 ...forces neutronics calculation, 28, 35

END ...final card of a job or subroutine, 14, 60

ERROR 1,16150,0,0 ...permits voluminous dump of programme common on
...detection of an error, 72

FIFA PU239 PU241 ' ' ...distinguishing card for a burnup user, 24, 36

FLUX = 1.+14 ...sets required flux level, 30

FORMAT(1X1P10E12.3) ...user's required format for output, 61

GB2 = 120*0.003 ...sets group-dependent buckling, 32

GFLUX = 1, 1.5, 2, 2.1, 3, 3, 4, 3.8, 5, 6.1, 6.2, 6.18,
6.2, 6.29, 8.1, 7, 80*7.2, 10, 11, 12, 12.2, 18*2, 0.1,
1.-2 ...sets flux spectrum (not level)- note continuation cards, 32

GO TO 3 ...control is transferred to statement numbered 3, 60

GROUPS 3 1 10 72 120 ...required condensed groups, 26

GSOURCE 1. 119*0. ...overwrites library fission source, 32

HEADb1 SAMPLE JOB ...user's choice of heading information, 17

IF (X(0)) 1,2,2 ...logical test, 60

ITERATE 5, .05 ...sets outer iteration control for equilibrium search, 71

J(1)=2 ...arithmetic card, 59

KEFF LIMIT = 1. ...sets exit limit for k_{eff} , 33

KINF LIMIT = 1.05 ...sets exit limit for k_{∞} , 33

MATS 'BEO' U233 (23,70)' ' ...required materials for output, 26

N 'BEO'=.04,(23,70)=1.-10' ' ...required concentrations, 21

OUTPUT 1 PRINTX ...printed group cross sections sought, 25

PART .01 'PU' ' ' ...particle size (for resonance theory), 22

PAUSE ...used with tape changing, 60

POWER = 11. WATTS/CC ...only a burnup user may use this, 30

PUNCHbOUTPUT NO 1 ...columns 13-72 punched as a card 1-60, 61

PUT 3 30X1 3 30S1 ...used with GYMEA library preparation, 17, 55

PX 51,121,1 ...punches output from X array, 61

REQD 10 ...required overlap and transport cross section, 27

RETURN ...returns from a user's subroutine, 62, 14

RX 671,790,1 ...reads data cards after CEASE, 61

SAVE 4 6*1 3*1 ...condensed data is saved in core, 27

SCAT 900. 'BEO' CRYbbb' ' ...scattering temperature and matrices, 17

SEARCH B10 .1, 1.2, 10., 1., 0.001 ...sets criticality search conditions, 29

SHUT DOWN ...burnup option, 31

START UP ...burnup option, 31

SUBROUTINE DAT ...first card of a subroutine, 62

TEMP=850 ...temperature used in resonance theory (if different from SCAT
...card temperature), 18

TYPEbbLOAD TAPE 260 ON 2 ...operator messages, 61

T= 0., 20, 0 ...forces neutronics calculation, 28, 35

UNLOAD 2 ...physical tape unload, 61

USE 1,70,2 ...libraries to be used, 15

VOID .4 ...rescales nuclide concentrations by .6, 22

WRITE10 TYPICAL BEO SYSTEM ...writes this message on the printer, 61

WX 551,670,1 ...outputs X variables on the printer with FORMAT specified, 61

X(9)=4.9 ...arithmetic card, 59

Z(0)=1.0 ...arithmetic card, 59

\$IBSYS ...7040 monitor card, example, 9

5.1 Layout of a Run

Some entities which help to describe the structure of a data deck must first be defined. In Section 4 a simple job was given as illustration. A job consists of all data cards for a problem, which starts with the problem identification

,NAME AND ANY COMMENTS

which has other control cards to force burnup calculations (detailed later), and which terminates

END

These cards are typical of input data which is punched in free format (22) anywhere on a card or set of cards. The job name is given by the 4 characters immediately following the comma of the leading card. This name is used to label output and serves to identify a job. A job must retain an identity since, if errors arise in a job, the job is terminated (see Section 10 however) and the next job is called. If the next job is dependent on the first, then it will also be terminated following an error in the first job, unless the standard data state (11) is restored thus

CALL STAN

before a neutronics calculation is attempted. Since data generated in one job is available in subsequent jobs the choice of boundary between jobs is somewhat arbitrary.

In addition to jobs, which are carried out in the order of loading, subroutines may be written by the user. These subroutines are common to all the jobs making up a load. A load thus consists of initial subroutines, jobs, and perhaps subroutines interspersed between jobs, loaded into the core for all data given prior to a card which is punched

CEASE

The selection of initial subroutines is based on the fact that the subroutines can be retained or discarded for the next load. Subsequent loads must be started with one of the cards

.DISCARD

,RETAIN

Between loads, data is always restored to the standard state as the compiler uses the COMMON data area for its own needs.

The compiler and subroutines are discussed in Section 10, but suffice it to say here that subroutines start with

SUBROUTINE EG1

they usually contain

RETURN

and they must conclude

END

Subroutines may call other subroutines, either written in the same fashion or provided internally, and the manner of calling is simply

CALL EG1

Unlike FORTRAN subroutines, arguments may not be appended to the calling statement.

The choice of a load is also somewhat arbitrary, although it should be borne in mind that a load may not contain more than 800 data cards and may not consist of more than 10 jobs.

It is possible for the user to test whether a previous job has been terminated owing to an error condition, even if the job is not in the same load, as one variable available to the user (12) is set thus:

X(0) = -ve if an error exit,

or X(0) = 0. otherwise. A typical test at the beginning of a job might be

IF(X(0)) 1,2,2

etc. Section 10 contains further details of the compiler language.

A full set of data for GYMEA is termed a run. A run starts with the necessary monitor cards, as given in the simple example of Section 4, and terminates with the monitor card

\$IBSYS

The first data card of a run is overwritten by the SEARCH routine (Richardson 1965) when this routine is transferred from 7040 core to the 1401 core. For this reason a spurious first card should be provided, punched say

*GYMEA

In general cards punched with an * in column 1 are comment cards designed to assist data checking as they are ignored by GYMEA.

5.2 Library Information

5.2.1 Library request

Before any reference is made to data in a library, a library tape or tapes must be specified thus

USE X-S 1 , 21 NUCLIDES, SCAT 2

Only the underlined are the important indicatives. The first number is the unit on which the cross section library is mounted (LIB), the second number (NN) is the number of the last nuclide beyond which the cross section nuclear data is ignored, and the third number is the unit on which the scattering transfer matrix library is mounted (LIBC). A cross section library must always be provided before a neutronics calculation is attempted; on the other hand a scattering matrix library is only necessary if group flux is to be calculated from neutron balance of a homogeneous system (7.2), and the actual unit must be 2, connected through the on-line 1401. Since this card forces the programme to check that the specified tapes are in fact GYMEA libraries LIBC should be punched as a negative number if a scattering matrix tape is not required. The number NN is mainly of interest to burnup users as they may save machine time by eliminating consideration of the nuclides at the end of the library. Other users may just as well use NN=100 (this is equivalent to the full nuclide set if the library does not contain 100 nuclides) as no time-saving results from using a smaller number.

If the libraries referenced in a previous USE card, not necessarily immediately before, are required to be checked to make sure they are specific libraries then the following type of card should be used

=GYMEA NDXD NDSC

The labels NDXD and NDSC must agree with the libraries mounted, cross section library first, otherwise the subroutine LIB is called. The user may write this subroutine himself if required (10). If LIBC < 0 then any non-blank label can be used for the second item. In addition to this checking feature, the =GYMEA card prints the library heading record of each library to provide the user with a record for later reference.

5.2.2 Library copies

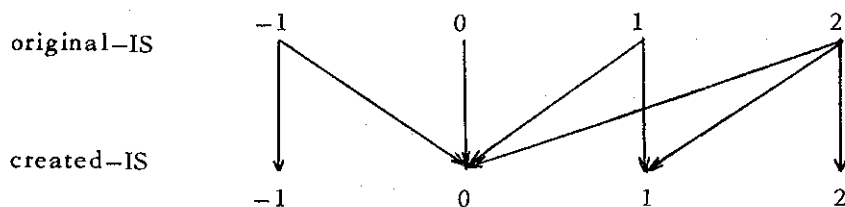
During the loading of information for chain 4 flux calculation, the user can create an additional library for each entry to this part of the programme. This created tape is then used in calculating reaction rates and collapsed and mixed cross section data rather than the original library— this need for reading the cross section library twice is explained more fully in Section 7. The saving of machine time that may result depends on the problem being undertaken and the actual original library. An example is given in JOB2, Section 15. For a problem not requiring burnup (5.3.9), only those nuclides specified for output (5.4.2) are created on the specified tape. The library may thus be considerably reduced in bulk with a consequent saving of machine time during later calculation of collapsed and mixed cross section data. This created library is a complete library in every detail, and may be used later in a run following a new USE card, or may be dismantled and saved for later use. The card that sets this option is, for example,

CREATE TAPE 3 NDX9 THUS 100*0

Again the underlined are the important indicatives and the example illustrates the use of repeated entry of a data word, 100*0 (22). The significant entries on the card have the following meanings:

- (i) the tape unit on which the additional library is to be created, or a negative number if create mode is no longer required,
- (ii) the identification label to be given to the library,
- (iii) the numbers (IS) that identify the type of nuclear data to be created, given for each nuclide of the original library (actually NN data words); -1 = light scatterer (scattering cross section is given rather than $\nu\sigma_f$), 0 = nuclide with only cross sections given, 1 = nuclide with shielding coefficients given (5.3.8) for the library-designated resonance groups (18.1.1), and ≥ 2 = nuclide with resonance parameters given for the resonance groups. If any entry is less than -1 then the data is created with the same structure as the original library.

A restriction obviously applies to the type of data that may be created, depending on the type of data given in the original library...



The additional information required for the production of shielding coefficients (7.4.3) is detailed later (5.3.8). It is important to note that data may be punched on more than one card following each control word, for example CREATE. To prevent the free input routine from reading through control words when insufficient data is given by the user, all continuation cards must be punched with a non-descending integer in column 1, which is tested, but otherwise ignored. This feature must be used with all data extending beyond one card; see however Section 10.1.4.

All libraries produced by GYMEA, and all printed output, are labelled with heading information which may be set by the user thus

HEAD SAMPLE RUN FOR I.P. AND G.R.

The alphabetic information punched in columns 13 to 54 is used as heading. The character given in column 12 is used as page control for printed output produced by chain 5 (8.3). In addition to the heading information given above, the date held in the 7040 core is always used as part of the heading record.

It has been shown that machine time may be saved by creating a reduced bulk copy of an original cross section library. It is not envisaged that machine time will be seriously affected by the bulk of a scattering matrix library, as the first required scattering matrix for a specific temperature is located by the 1401 working in parallel operation with the 7040 (Richardson 1965). Nevertheless a feature is available which permits the bulk of a library to be reduced. As shown later in this section, it is possible to group-collapse a scattering matrix library, and the newly-created scattering matrix tape must have the same structure as the original. This forms a strong motive for having available a feature for deleting information from an original library.

The card used to cause this copy to be made is, for example,

COPY TAPE 2 NDS8 THUS 1, 5*0, 1, 3*0

The underlined items have similar meanings to the items on the CREATE card with the exception that the fixed point numbers that follow the label have the following meaning

0 = don't copy

1 = copy

and numbers must be given first for the material set of the original scattering matrix library (18), 6 in the example, and then for the temperature set (18), 4 in the example. This may be understood if it is noted that the structure of the scattering matrix library is such that data for all materials is given for every temperature of the library temperature set.

Again the information on the HEAD card and the date is used to form a heading record. A point that was not previously emphasized is that a scattering matrix may be mounted or created on any unit, but if the tape is to be used in neutronics calculations it must be mounted as unit 2 and must be connected through the on-line 1401. Unlike the CREATE feature, which is delayed until the neutronics calculation is under way, COPY is carried out on detecting the data card.

5.2.3 Condensed and mixed GYMEA library

Most of the comments relating to preparation of condensed and mixed rather than simply copied or data-deleted libraries are better discussed in the section dealing with this type of preparation (9).

A card that sets the cross section library unit and prepared library name, and the scattering matrix unit and prepared library name, may be punched, for example:

PUT X-S 3 XS01 SCAT 0 SC01

The programme includes provision for the preparation of both libraries on the same unit; the cross section library precedes the scattering matrix library.

5.2.4 Scattering matrices request

As discussed earlier in this section, it is generally necessary to nominate a tape unit for the scattering matrix library (USE), and the programme can be made to check that the correct library is mounted (=GYMEA). As it is possible to have available scattering matrices for one material for different thermal scattering models, each scattering material to be included in the study must be nominated, as well as its method of preparation. The card used for this may read, for example:

SCAT 900., 'BEO' CRY ' '

It is necessary of course to know the temperatures, materials, and mode of preparation of available scatterers in the scattering matrix library. In the example above the first item is the required temperature in degrees Kelvin, which must agree with an entry in the library temperature set table to within one degree. The given temperature also sets the temperature to be used in resonance theory (7.4); this may, however, be changed as indicated below. The next entries are alternately material and mode of preparation, given for all scatterers of interest in the problem in hand, terminating with inverted commas ' '. In general, chemical names enclosed in commas, for example 'BEO', describe a material that is not necessarily given in the cross section library (5.3). The correspondence is determined by data cards which specify the mixing rules. The mode of preparation must always be given as 6 characters, for example,

CRYbbb (b = blank)

Abbreviations such as CRYbbb for crystal scattering and GASbbb for monatomic gas scattering are convenient for indicating the required scattering law; these conventions are, however, the conventions of authors of the scattering matrix library and have no explicit meaning to the code GYMEA.

As a rule the materials given in the scattering matrix library set are materials which must be defined by the user to obtain the correspondence with nuclides in the cross section library (19). This layout of scattering matrix library is not essential, but it does provide the user with the means of adding proportions of scattering matrices to allow for, say, the scattering of neutrons in O^{26} associated with heavy nuclides when they are present in a matrix of a dominant scatterer, say BeO. In Section 5.3.1 such an application is given as an example. The complication associated with materials stems from the fact that fundamentally the code, since it is a burnup code, must deal only with nuclides. The code never has available concentrations of materials, and these are always produced when required from the nest of mixing rules given by the user and the nuclide concentrations.

5.2.5 Resonance temperature

It is possible to change the temperature specified on the SCAT card as far as it determines the temperature used in resonance theory. A data card that will cause this change may be punched, for example:

TEMP 850. K

5.3 Nuclide Information

5.3.1 Material specifications

Three possible representations of material specifications are given below:

- (i) Nuclide labels, for example BE9 (or BE 9), are used to identify specific nuclides available in the cross section library.
- (ii) Material labels, for example 'BEO' (note the inverted commas enclosing at most four characters), are used to identify specific materials either available in the cross section library or defined using mixing rules detailed below.
- (iii) Nuclide sets, for example (1,5) (note the brackets), are used to identify nuclides of the cross section library, 1 to 5 in the example cited.

In addition, the blank material label ' ' is used to indicate the end of nuclide information, and commas or blanks may be used to separate nuclide information and assist readability. The distinction between nuclides and materials which are given in the library is unimportant except to authors of cross section libraries, who must be able to describe burnup mechanisms adequately (18). The distinction between materials which are given in the library and materials defined by the user is only important if an application required by the user requires material densities. Such densities are available only for defined materials (5.3.5 and 5.3.7).

5.3.2 Mixing rules

A mixing rule is given by the general statement:

material = weighted sum of nuclides or materials, either given in the library or defined previously, with density (g/cm^3) given immediately before the terminating blank material label ' '.

For example,

'BEO' = 1.BE9, 1.O16, 1.8' '

is read as 'for every one molecule of BeO we have 1 atom of Be⁹ and 1 atom of O¹⁶, the density of BeO being 1.8 g/cm^3 '. Alternatively this can be written

'BEO' = 1.BE9, 1.O16, DENSITY 1.8' '

or 'BEO' = 1.(1,2), 1.8' ' IF BE AND O 1 AND 2 OF LIB

In the second example DENSITY is not interpreted by the free input routine as it is seeking a number. This possibility of insertion of check words is typical of the application of the free input routine to GYMEA. However, check words used to assist data checking should be avoided when the routine is seeking alphabetic information. Further details are given in Section 22. For those people not familiar with the inverted comma, it is punched $\frac{4}{8}$ in the one card column (lower case =) and is frequently written as @ for punching detail to avoid confusion with other characters.

It is shown in Section 10 that elements of a Z array can be calculated or set. This feature can be used with mixing rules, since a negative weight is replaced by the implied indirectly referenced element of the Z array. For example,

Z(20) = 0.76

Z(21) = 0.17

Z(22) = 0.07

'PU' = -20 PU239, -21 PU240, -22 PU241, 0' '

defines 'PU' to consist of 0.76 parts of Pu²³⁹, 0.17 parts of Pu²⁴⁰, 0.07 parts of Pu²⁴¹; the density is zero, which is satisfactory for applications not requiring material density. The elements of the Z array may be defined later than the appearance of the card defining 'PU', as the interpretation is not made until immediately before actual use of the information given on the card, for example, to define a material concentration of 'PU'. Density may also be specified in this indirect fashion by giving a negative value. Again a feature of the free input routine has been used, namely, the decimal point was not punched for the floating point data, as the routine provides an imagined decimal point to the extreme right of the number (22). The material 'PU' may be used, if required, in mixing rules that follow, for example

'PUO2' = 1.'PU', 2.O16, 11.' '

Normally mixing rules must be associated with materials given in the scattering matrix library and selected using a SCAT card (5.2.4), for example,

'BE9' = 1.BE9, 0' '

'O16' = 1.O16, 0' '

As a more involved application, indicating the reason why the scattering matrices are usually given as materials which must be defined by the user (19), consider an example in which

$$N_{\text{BeO}} = 4 \times 10^{-24} \text{ nuclides/cm}^3$$

$$N_{\text{U}^{235}\text{O}_2} = 1 \times 10^{-4} \text{ nuclides/cm}^3$$

and the following SCAT card is used:

This interpretation may not be of universal application, but it is a simple matter to change the table using the library edit programme, EDITOR (Ford 1966- see brief details Section 19).

5.3.4 Concentration definition

Nuclide or material concentrations may be given simply as ratios (5.3.5), or as nuclides/ 10^{-24} cm³. Normally these concentrations are entered using a card which starts N and which terminates ' ', although it is possible to calculate concentrations and have the result available to GYMEA through the user's COMMON data area (12). An example of the latter method is

$$X(51) = 1.E-4$$

and an example of the former method is

$$N 'BEO' = 1.-4 ' '$$

Note that the condensed exponent notation is available for floating point data (E absent) using the free input routine (22). Remembering that only concentrations of library nuclides are stored by the code (5.2.4), one essential difference exists between the two methods - the first enters 1.E-4 as the concentration of Be⁹, provided it is the first library nuclide, and the second adds 1.E-4 to the previous concentrations of Be⁹ and O¹⁶. This last feature is a result of the requirement, for defined materials, of adding say O¹⁶ from BeO to O¹⁶ from ThO₂ to give the correct O¹⁶ content of a mixture of BeO and ThO₂. To avoid confusion with materials possibly available as part of the library data, this feature of adding concentrations is extended to cover all single nuclides and materials. When the concentration of a set of library nuclides is given, however, the concentration specified overwrites the previous data. This, plus the fact that more than one concentration may be specified on a card, enables concentrations to be easily set; for example,

$$N(1,100)=0., 'BEO' = 1.-4 'THO2' = 1.-6 ' '$$

will set

$$N_{Be^9} = 1. \times 10^{-4}$$

$$N_{O^{16}} = 1.02 \times 10^{-4}$$

and $N_{Th^{232}} = 1. \times 10^{-6}$

provided the materials 'BEO' and 'THO2' have been defined prior to the use of this card.

To enable nuclide concentrations to be readily punched using a computer programme, including GYMEA itself, concentrations of all nuclides (NN) may be given in library order following the control word ALL N, for example,

ALL N 1.00E-04 1.02E-04 ...ETC...

5.3.5 Normalization of concentrations to unit volume

A subroutine NUC is available which transforms nuclide or material ratios to nuclides/ 10^{-24} cm³. Only defined materials should be used in detailing material ratios, as the density of each material is used to calculate nuclide concentrations assuming that there is no change of volume on mixing. The concentration of any other nuclides specified will be set zero without an error condition resulting. An example of the use of this feature is

$$'BEO' = 1.BE9, 1.O16, 1.8 ' '$$

$$'THO2' = 1.TH232, 2.O16, -5 ' ' I.E.Z(5) = DENSITY$$

$$N(1,100) = 0., 'BEO' = 100, 'THO2' = 1 ' '$$

CALL NUC

On entry to the subroutine the array AFT.N (X(51)-X(150)) contains nuclide ratios derived from the material ratios given on the data card, and on exit the array AFT.N contains these concentrations normalized to nuclides/ 10^{-24} cm³.

It is important for this application that only materials required in the calculation should have (non-zero) densities given and that the materials with non-zero densities should each contain a unique dominant nuclide (5.3.3). Materials with zero densities given are ignored. The sample calculation, JOB2 of Section 15, illustrates these remarks.

5.3.6 Renormalization of concentrations to allow for void

When a data card, punched, for example,

VOID 0.05 I.E. 5 PERCENT

is detected in the input, the concentrations held in the array AFT.N are renormalized by the factor (1-VOID).

5.3.7 Effect of resonance self shielding of separated fuel particles

The procedure given by Keane (1964a) for modifying the potential scattering cross section of moderator nuclides, σ_s , when fuel particles are separated from the bulk of moderator, gives

$$\sigma_s^* = \sigma_s / (1 + \bar{\ell} \sum_p V_m / V_p)$$

which is applied in resonance theory as

$$\sigma_p^* = N_{\text{mod}} \sigma_s^* / N_{\text{res}}$$

for all library moderator nuclides (18), where

$\bar{\ell}$ = modified mean chord length for the fuel particles volume V, surface S

$$= \frac{3a}{2} \left(\frac{4V}{S} \right)$$

= aD of Keane and he recommends using $a = \frac{1}{2}$,

D = mean diameter of particles in cm, and

V_m / V_p = ratio volume of homogeneous mixture to volume of separated particles.

The potential scattering cross sections for moderator nuclides, originally extracted from the specified cross section library tape, are adjusted automatically when a card such as

PART 0.01 'TH' ' '

is detected. In this example $\bar{\ell} = 0.01$ cm is required (200 micron particles) for separated fuel particles composed of all the Th originally given as a homogeneous mixture, that is, the concentrations for the homogeneous system must be given before this card. The materials given on the PART card must be defined by the user as the densities are required in this application. A further feature which helps if $\bar{\ell}$ is to be calculated using the compiler, is that a negative data entry refers to an element of the Z array (12); for example if $\bar{\ell} = -1$, then Z(1) is used instead. New PART cards may be used as frequently as required - note that a card must be used if the initial composition is changed.

The total potential scattering cross section, σ_{pt} , used in GYMEA resonance theory (7.4.2) then amounts to the following:

$$\sigma_{pt} = \sigma_{sr} + \text{sum}(\sigma_p^*)$$

where σ_{sr} = the potential scattering cross section of the resonance nuclide being studied, and the sum is taken over all moderating nuclides. The expression for σ_{pt} given above is of the same form as the expression used in calculating heterogeneous effects of a more general nature than dispersed fuel particles. In such applications, $a = 2/3$ is generally used, and then $\bar{l} = 4V/S$.

The user may wish to adjust the potential scattering cross sections of the nuclides himself using the compiler and the fact that the cross sections are held in the user's COMMON (12) as X(351) - X(450) for nuclides 1-100. It should be noted that the resonance theory requires the identity of each nuclide to be retained as the theory determines a factor λ for each potential scattering cross section. The feature mentioned should prove useful although it must be used with some care.

5.3.8 Concentration definition for shielding theory

It is shown in Section 7.4.3 that a shielded cross section at one temperature is given by

$$\sigma_x = \sigma_\infty / (1 + \frac{A}{\sigma_p} + \frac{B}{\sigma_p^2})^{1/2} \quad (x = \gamma, f, \text{ and } \nu f),$$

where

σ_p = total moderator potential scattering cross section per absorber atom, and

σ_∞ , A, and B are the shielding parameters to be determined for each reaction type x.

The three parameters are determined from three forced calculations of the shielded resonance cross section using resonance theory (7.4.2). The forced calculations are carried out for three different specified nuclide concentrations contained in three arrays designated \tilde{N}_1 , \tilde{N}_2 , and \tilde{N}_3 . The nuclide concentrations may be entered into these arrays using three subroutines. These subroutines require the concentrations to be defined using the nuclide concentration card (that is stored in the array called AFT.N), for example,

N(1,100) = 0., BE9 = 4.-2, TH232 = 4.-4^{1 1}

The subroutine

CALL NS1

then transfers selected elements of AFT.N to the array \tilde{N}_1 . Two rules are observed in this selection:

- (i) if the library does not contain any shielding factors (IS=1) then all concentrations stored in AFT.N are transferred to \tilde{N}_1 , or
- (ii) if the library contains at least one shielded nuclide (IS=1) then only the concentrations of those nuclides that have resonance parameters given (IS \geq 2) are transferred to \tilde{N}_1 .

The subroutines

CALL NS2

CALL NS3

play a similar role but relate to the arrays \tilde{N}_2 and \tilde{N}_3 respectively. The reason for the selection rules cited above is that entry of data for the calculation of shielding parameters for a specified nuclide should not destroy previous data possibly read from the library tape (18).

where

N = nuclide concentration AFT.N

σ = collapsed cross section (6 reactions given)

ϕ = flux

S = source strength, hence $N\sigma\phi/S$ is a reaction rate per source neutron

N^* = nuclide concentration AVE.N

(N and N^* are the same for a non-burnup user)

If two successive digits are the same, the output terminates on the second digit of the pair. Using the compiler (10), X(9) can be set as part of the input data; thus

X(9) = 4.809

would result in two tables being produced, the first for σ and the second for $N\sigma\phi/S$, with the latter starting on a fresh page. The table for σ would follow a standard table, starting on a new page, which lists k_{eff} , k_{∞} , etc. (6.5). Only the first 7 digits may be used for this application.

If X(9) is negative, no output is produced unless $X(9) \leq -10$, in which case the complete set of condensed and mixed cross sections is produced as a further file (the first file is the compiled input data) on tape 4. The user may wish to prevent, or save up, output during an iterative procedure carried out under control of the compiler. When output is saved up on tape 4, it can be used later with the aid of the subroutine OUTPUT which is obtained thus:

CALL OUTPUT

Of course it is not possible to change the number of groups or number of materials between generating the information and using it (see also Section 13.3).

Output from chains 5 and 6 is requested using a card

OUTPUT 2 CRAM PRINTX

which informs the programme that 2 types of output are required, namely CRAMbb (b=blank) and PRINTX. The types of output that may be requested are:

PRINTX	(printed cross section data given one material at a time)
CRAM	(FORTRAN version of Hassit's code(1962))
4-ZOOM	(FORTRAN version of 9-ZOOM (Stone et al. 1959))
DSN	(FORTRAN version of Carlson's code (Carlson et al. 1960))
WDSN	(WINFRITH DSN code (Francescon (1963))
PRINTS	(printed scattering matrix data)
GYMEA	(a PUT card must also be used (5.2.3); Section 9 gives details of preparation of new condensed and mixed cross section and scattering matrix libraries for GYMEA).

From time to time other codes will be added to this list as they become available on the computer used at the A.A.E.C. For the user wishing to extend the source code, Section 20 contains brief details of the programme source structure.

A single fixed point number is given on a card, for example,

REQD 11

where the tens digit has the interpretation

0 = no overlap

1 = overlap

and the units digit has the interpretation

0 = method (i)

1 = method (ii), normal diffusion code requirement.

The standard state of the code (11) sets this variable to 11. During print out the variable is referenced

XS-OT= 11

and, if the option required cannot be accommodated by the current calculation of the code (7.4.5), the option used instead is printed with a minus added as a warning indication.

Entries on the MATS card are restricted to not more than 40 words of indicative information, based on a nuclide or material producing 1 such word or a material set producing 3 words for each data item on the card.

5.4.3 Output for later computation

The user can set the condensed and mixed cross section subroutine to save selected, weighted, condensed and mixed cross sections in the user's COMMON (12). Using this feature it is possible to calculate quantities by manipulation with the compiler (10). Entries given on the appropriate SAVE card relate to a previously used, although not necessarily immediately preceding, MATS cards, for example,

MATS U233, 'BEO', (3,8) ' '

SAVE AS 4 2*1, 4*0, 100 0 -1

The fixed point numbers supply the following information:

- (i) The first number is a digit 0 to 7 which has the same interpretation as digits of X(9) and X(6) (5.4.1), that is σ is to be saved for the example.
- (ii) The second to seventh numbers are six digits 0 or 1 which relate in order to individual neutron reactions to be saved or not;

0 = do not save

1 = save,

and reaction 1 = (n, γ) usually

2 = (n,f)

3 = (n,abs)

4 = (n, ν f) or (n,n) for moderator nuclides

5 = (n,transport)

6 = (n,2n) usually ,

(see library details, Section 18.1). Thus σ_γ^1 and σ_f are required for the example given.

- (iii) The remaining numbers, comprising one for every entry on the MATS card- (3,8) is one entry- have an interpretation which is best explained in terms of the example:

100 = save first entry for U^{233} in Z(100), second entry in Z(101), etc.,

0 = do not save information for BeO

-1 = use index register J(1) (see Section 10) and save first entry for nuclide 3 in Z(J(1)), second entry for nuclide 3 in Z(J(1) +1), etc.

If 3 condensed groups were requested with a previous card, then entries saved for U^{233} in the example would result in the Z array containing the following information:

Z(100) = σ_γ (group 1),

Z(101) = σ_γ (group 2),

Z(102) = σ_γ (group 3),

Z(103) = σ_f (group 1),

Z(104) = σ_f (group 2),

and Z(105) = σ_f (group 3).

If X(8) \neq 0, then, instead of information being directly stored in the Z array as indicated above, X(8) weighted entries would be added to existing Z elements thus:

Z(100) = Z(100) + X(8) * σ_γ (group 1), etc.

Using this feature it is possible to obtain time-averaged quantities, for example.

5.5 Forcing a Neutronics Calculation

5.5.1 Flux from neutron balance

Two types of time card cause the programme to call the neutronics and burnup chain, chain 3. These cards are discussed at length in Section 6.1. However, a user not intent on burnup, but requiring group flux to be calculated from neutron balance of a homogeneous mixture of materials with specified concentrations (7), will probably use a card punched

T = 0, 20, 0

or DT = 0, 20, 0, 1

which simply sets reactor time to zero, that is, clean core status, sets an upper limit to the number of iterations required in the flux calculation, and sets the maximum number of criticality iterations to zero. This latter number is only positive when a parameter is to be adjusted to achieve a specified state of criticality.

5.5.2 Criticality

If a parameter is required to be adjusted to achieve criticality, then, before a time card punched, say,

T = 0., 20, 6

can be used, an eigenvalue search card must be given. Such cards may be punched, for example,

A card punched with a negative number of iterations sends control to the subroutine that calculates reaction rates and condensed and mixed cross section data (7.3) without first entering the flux balance subroutine. Zero iterations also has an interpretation, which is that both the subroutines mentioned are bypassed. Of course a calculation of the type mentioned above must have previously been carried out— this feature is really only of interest to a burnup user with non-changing reaction rates.

Westcott flux (13.2) with a specified r value (ratio of epithermal to thermal flux) is obtained using three cards punched, for example,

```
FLUX = 1.+13
X(0) = 0.02
CALL WEST
```

The first sets the flux level (5.6), the second gives the r value, and the third forces the calculation to proceed. The neutronics calculation can then be carried out using a time card as the subroutine WEST will have generated the required elements of group flux.

The reader may wonder at this stage why a FLUX card is not always required. The reason is that the programme sets the flux level appropriate to neutron balance from unit source if no flux (or FLUX=0.) is given (5.6). If the flux is to be calculated in this fashion using WEST, data may be punched thus:

```
FLUX = 1.
X(0) = 0.08
CALL WEST
FLUX = 0.
T = 0., 1, 0
```

It should be noted that group flux is set to the level given on the FLUX card not when this card is used, but rather when WEST, or the subroutine that calculates reaction rates and mixed and condensed cross section data, is called.

5.6 Reactor Control Information

It was shown above that the total flux level can be set with a card

```
FLUX = 1.+13 N/CM**2/SEC
```

and it was noted that this does not immediately set the group-dependent flux to the specified level. Other possible data of this section may only be used by burnup users, that is, those users who have indicated their intention with a FIFA card (5.3.9). Firstly it should be noted that the use of a FLUX card also tells the programme that during burnup the FLUX level is to remain at the level specified. If in addition the group-dependent flux does not change, then the one analytic solution of the burnup equations (6.2) holds for all time, and advantage may be taken of this feature by using time steps chosen simply to provide output information. However, if the group-dependent flux, the spectrum, does change, time steps must be chosen consistent with this, using, say,

```
DT = 100., 20, 0, 1
```

(see also Section 6.1).

Burnup users may set the power level thus:

```
POWER = 10. WATTS/CM**3
```

which tells the programme to adjust the level of previously calculated reaction rates (7.3.1) to this level before proceeding with burnup. The status of constant flux level burnup can be retained using two cards, for example,

POWER = 10.

FLUX = 0.

which will overwrite zero FLUX with that calculated to give the specified power. From then on this calculated flux level will be retained. Similarly the status of step-wise constant power level burnup can be retained using

FLUX = 1.+14

POWER = 0.

If the flux is to be calculated at a level consistent with unit source then the feature previously indicated in Section 5.5.3 can be used, namely,

POWER = 0.

FLUX = 0.

The last mentioned cards appear to have shut the reactor down, but this is not the interpretation of the programme. Instead the SHUT DOWN option is set using a card punched

SHUT DOWN

When this card is used the previously defined status is retained, so that later the option for START UP can be set using a card punched

START UP

Since the SHUT DOWN card sets reaction rates to zero, a calculation of fresh reaction rates must be forced before proceeding. This can be done using, for example,

DT=0.,20,0,1

These features are of use to those interested in, say, the reactivity saving resulting from allowing certain nuclides, for example Pa²³³, time to decay, and are best clarified with an example.

Example

An infinite homogeneous reactor composed initially of Be⁹(4×10^{-2} nuc./ 10^{-24} cm³), U²³³(1×10^{-5} nuc./ 10^{-24} cm³), and Th²³²(1×10^{-4} nuc./ 10^{-24} cm³) is to be burnt for 500 days at a constant power level of 0.1 watt/cm³ and a temperature of 300 °K. It is then to be shut down for 10 days and started up again. What is k_{eff} after a further 90 days burnup?

Excluding monitor cards and cards associated with voluminous neutronics output (5.4), the data for this problem may read:

*GYMEA

,EX1 ILLUSTRATES BURN UP
 USE 1,61,2
 SCAT 300. BE9 GAS ' '
 N(1,100)=0., BE9=4.-2, U233=1.-5, TH232=1.-4 ' '
 POWER = 0.1
 FIFA U233 ' '
 T =0.,20,0
 DT =125.,20,0,4
 SHUT DOWN
 DT =10.,0,0,1

*NOTE...NO FLUX ITERATIONS

START UP
 DT=0.,20,0,1
 T=600.,20,0
 END
 CEASE

5.7 Group Information

Three types of group information may be entered as data, using cards punched, for example,

GFLUX = 1. 1.5 1.8 2. 5. 6.

1 8. 10. , ETC...

GSOURCE = 0.54, 0.46, 118*0

GB2 = 120*1.-4 1/CM**2

Data must be given for the library number of groups (NG) and since this data frequently extends beyond one card a non-descending integer must be punched in column 1 of continuation cards as explained in Section 5.4.2.

The first type illustrated above defines the spectrum of group flux but not the level (5.6). The second defines the spectrum of group source, which will be used rather than the library spectrum. The level is adjusted to a total source of 1 n/cm³.sec. The third defines a group-dependent buckling, which is normally defined with all elements the same. However it is possible roughly to simulate a reflected core for a burnup study, using a typical calculation carried out with a space-dependent code which calculates

$$B^2 = - \left[\frac{\int \nabla^2 \phi dV}{\int \phi dV} \right]_{\text{core}}$$

for each neutron energy group.

All of the above quantities may also be entered via the X array of the user's COMMON (12).

5.8 Flux Accuracy and Burnup Mechanism Information

The user may change the accuracy parameter used as a convergence test in the calculation of group flux from neutron balance (7.2). For example

ACCURACY = 1.-5

will set a parameter which will stop further iterations of group flux when roughly speaking the total absorption and leakage rate from unit source changes by less than 10^{-5} from one iteration to the next.

The standard state (11) consists of an ACCURACY parameter of 10^{-4} . It is not anticipated that this feature will be used except by people studying the effect of small changes on a system.

Frequently it is difficult to assess the relative importance of an individual nuclide on the burnup mechanism. It is possible for the user to change the burnup mechanism slightly with a card punched, for example,

BURN 2*-1, 1, 97*0

which should contain entries for each nuclide being USED (NN nuclides in all). Each entry of a fixed point number sets an option for each library nuclide in turn. The options are:

- (i) -1 do not burn up (still produced however)
- (ii) 0 follow library rules
- (iii) 1 burn up (does not produce further library nuclides however).

The example given in Section 6.2 better explains the function of this type of card.

5.9 Other Data

It is possible for the user to specify lower limits for k_{eff} , k_{∞} , B2AV (6) using the cards

KEFF LIMIT = 1.

KINF LIMIT = 1.1

AVX LIMIT = 1.-4

When a quantity falls below a level set in this manner the next job is called (SEND). The standard state (11) for each of these limits is 0., 0., and -1. respectively.

The user may make changes to data that is set by the subroutine STAN (11), or that is read from the heading file of the library (18), using the compiler discussed at length in Section 10. For example it may be desired to change the potential scattering cross section of the first library nuclide to 5.4 barns. This would be achieved thus:

X(351) = 5.4

6. CHAIN 3- BURNUP

6.0 Summary

This chain controls the neutronics calculation and burnup. Burnup is carried out analytically based on the burnup mechanisms given as part of the cross section library. Time-dependent quantities, such as FIFA, are also calculated analytically using average concentrations for a time step. The analytic solution procedure requires the 1-group microscopic reaction rates to remain constant, hence it may still be necessary for the user to specify 'integration' time steps consistent with the needs of his problem - say, 100 days. Chain 4 (7) carries out the calculation of group flux and cross sections and chain 3 uses the results to calculate k_{eff} , etc. The concentration of any nuclide, mixture, or level of group-dependent buckling may be adjusted to achieve a specified value of k_{eff} if required.

6.1 Data Options

Control is transferred to chain 3 from the input chain 2 when a 'time' card is detected. The data supplied with the card controls the functioning of the chain. Two types of 'time' cards are permitted, for example,

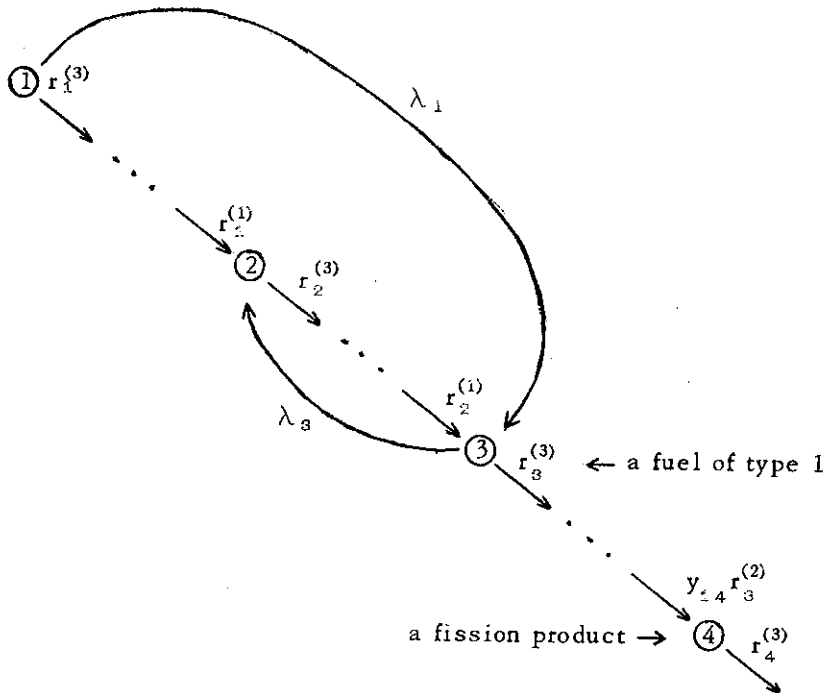
T=0.,20,0

DT=100.,20,0,1

but both determine the same quantities. The data entries given in order have the following meanings:

- (1) determines the time (in days) after first start up ($T=0.,...$) at which output is required for the behaviour of the reactor system being studied (if the specified time is negative then X(41) is used instead),
- (2) determines the option required with the flux spectrum calculation (LSID),
- (3) determines a limit to the number of permitted iterations to achieve criticality on, say, buckling (5.5.2), and
- (4) which is only used with the DT card, is the number of such time steps required at one entry to chain 3 - control is not passed to chain 5 until the number of repeats is satisfied, hence CRAM output, for example, would not be produced for the intermediate steps.

The manner in which the chain works can best be understood by referring to the flow diagram which follows.



Here $r_l^{(k)} = \sum_{i=1}^{NG} \sigma_{li}^{(k)} \phi_i$ for nuclide l ,

where subscript i indicates the neutron energy group,

superscript (k) indicates reaction (k) of the library,

namely (1) = reaction (1), usually (n, γ) ,

(2) = fission,

(3) = absorption (7.3),

and λ_l is the decay constant (10^{-24} /sec) of nuclide l .

Section 7 and Section 18 elaborate further on the quantities introduced here.

The quantity y_{34} is the yield of nuclide 4 from fission of fuel of type 1— say nuclide 3 was designated as being a fuel of this type. Let N_l = the nuclide concentration of nuclide l (nuc./ 10^{-24} cm³) at time t (10^{24} sec) after the previously considered time step of the reactor system composed of a homogeneous mixture of the nuclides designated 1 to 4. The following nuclide balance then applies for each nuclide;

Time rate of change of N_l

- = loss due to decay of nuclide l
- + loss due to n-absorption by nuclide l
- + gain due to decay of nuclides which produce nuclide l
- + gain due to n-capture reactions which produce nuclide l
- + yield from fission of nuclide l from the four possible fuel types.

Using a dot to denote differentiation with respect to the variable t , the set of balance equations may be written:

$$\begin{pmatrix} \dot{N}_1 \\ \dot{N}_2 \\ \dot{N}_3 \\ \dot{N}_4 \end{pmatrix} = \begin{pmatrix} -r_1^{(3)} - \lambda_1 & 0 & 0 & 0 \\ r_1^{(1)} & -r_2^{(3)} & \lambda_3 & 0 \\ \lambda_1 & r_2^{(1)} & -r_3^{(3)} - \lambda_3 & 0 \\ 0 & 0 & y_{14} r_3^{(2)} & -r_4^{(3)} \end{pmatrix} \begin{pmatrix} N_1 \\ N_2 \\ N_3 \\ N_4 \end{pmatrix}$$

or $\dot{\tilde{N}} = \tilde{A} \tilde{N}$. -(1)

Notice that A is a lower triangular matrix except for the term λ_3 - this form is typical for a complete library set of coupled equations with only a few isolated just off-triangular elements.

Before the solution of Equation 1 is considered, a few additional features will be discussed. In a general set of equations, reaction (6) (usually $(n, 2n)$) may produce a nuclide belonging to the library set, hence production terms $r^{(6)}$ would appear in the matrix A .

Use of a card such as

BURN 0 1 0 -1

(5.8) for the example being studied would force nuclide 2 to burn out, rather than to burn to nuclide 3; hence the production term $r_2^{(1)}$ would not appear in A , and, since nuclide 4 would be prevented from burning out, the loss term $-r_4^{(3)}$ would not appear in A . This type of card forms a convenient way to investigate the effect of chosen nuclides on the overall burnup mechanism. It should be borne in mind that the BURN card does not alter decay or yield from fission.

Returning to the solution of the set of equations (1), the first step is to carry out transformations on \tilde{A} to yield a lower triangular matrix. In the example chosen this may be achieved using the matrix

$$\tilde{I}_3(x) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & x & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \tag{2}$$

Now $\tilde{I}_3(x) \tilde{I}_3(y) = \tilde{I}_3(x+y)$,

$$\therefore \tilde{I}_3(x) \tilde{I}_3(-x) = \tilde{I}_3(0) = \tilde{I}$$

$$\therefore \tilde{I}_3(-x) = \tilde{I}_3^{-1}(x)$$

Then, from Equation 1,

$$(\tilde{I}_3(x) \dot{\tilde{N}}) = (\tilde{I}_3(x) \tilde{A} \tilde{I}_3^{-1}(x)) (\tilde{I}_3(x) \tilde{N})$$

which is of the same form as the original Equation 1, namely

$$\dot{\tilde{N}}^* = \tilde{A}^* \tilde{N}^* , \tag{3}$$

except that \underline{A}^* may be made lower triangular by a suitable choice of x . If x is chosen so that

$$x^2 a_{32} + x (a_{22} - a_{33}) - a_{23} = 0 \tag{4}$$

then the element a_{23}^* is zero in the transformed matrix and no other upper triangular element appears, as any off-triangular elements are always separated from one another by at least one zero element—a requirement of the ordering of the library nuclides. The roots of Equation 4 are always real as the off-diagonal terms are always positive, and the actual choice of root is immaterial as the new diagonal terms a_{22}^* and a_{33}^* simply interchange roles if one root is used rather than the other. For the more general solution of Equation 1 applied to other library nuclide sets, a series of transformations $I_{i_1}(x_1) I_{i_2}(x_2) \dots$ may be necessary; however a set of equations (3) is eventually obtained for which \underline{A}^* is lower triangular.

Assuming that all leading diagonal elements a_{ii}^* are different from each other, excluding possible zero elements, then the eigenvalues of \underline{A}^* are simply the diagonal terms and a solution of the Equations 3 of the form

$$\underline{N}^* = \underline{B} \underline{e} \tag{5}$$

can be assumed, where \underline{B} is the required solution matrix and is lower triangular,

$$\text{and } \underline{e} = \begin{pmatrix} \exp(a_{11}^*t) \\ \exp(a_{22}^*t) \\ \exp(a_{33}^*t) \\ \exp(a_{44}^*t) \end{pmatrix}$$

In practice the requirement that the eigenvalues are distinct is always met and will only be violated if an attempt is made to use a library containing the same nuclide in several different positions. Differentiation of Equations 5 and comparison with Equations 3 then gives

$$b_{ij} = \frac{\sum_{k=j}^{i-1} a_{ik}^* b_{kj}}{a_{jj}^* - a_{ii}^*} \quad \text{if } i \neq j \tag{6}$$

for the off-diagonal terms, and the initial condition $\underline{N}^* = \underline{N}^*(t=0)$ leads to the equation

$$b_{ii} = N_i^*(t=0) - \sum_{k=1}^{i-1} b_{ik} \tag{7}$$

for the diagonal terms. The result for $i=1$ may be imagined contained in Equations 6 and 7 if we set backward sums to zero, $\sum_{k=1}^0 \dots = 0$. Equations 5 and 6 express a recurrence relation which enables

all the elements of the solution matrix \underline{B} to be determined provided the indices are ranged as follows:

$$(1,1); (2,1), (2,2); (3,1), (3,2), (3,3); \text{ etc.}$$

The set of nuclide concentrations \underline{N}^* may then be obtained from Equation 5 and the original library nuclide set may be obtained from the simple transformation:

$$\tilde{N} = I_3(-x) \tilde{N}^* = \begin{pmatrix} N_1^* \\ N_2^* - x N_3^* \\ N_3^* \\ N_4^* \end{pmatrix} \quad (8)$$

One objection that may be raised concerning the analytic solution technique is that it appears to be prone to round-off error, especially for the production of He³ from the Be (n,α) reaction. However the experiences of the authors indicate otherwise. Stated differently, the analytic procedure enables long (say 100-day) time steps to be taken and hence the small (say 1-day) time steps, which would produce the greatest round-off error, are unnecessary.

6.3 Analytic Calculation of Average Nuclide Concentrations

Since the analytic solution of the burnup equations is available in the form of Equations 5 and 8, time integrals of power and FIFA are best given in terms of average nuclide concentrations for the time step. The average concentrations for the time step are simply

$$\hat{\tilde{N}} = I_3(-x) B \hat{\tilde{e}} \quad (9)$$

where $\hat{\tilde{e}}$ contains elements of the form

$$\begin{aligned} &= \frac{1}{t} \int_0^t \exp(a_{ii}^* t') dt' \\ &= \frac{1}{a_{ii}^* t} [\exp(a_{ii}^* t) - 1] \end{aligned} \quad (10)$$

For small argument $a_{ii}^* t$, Equation 10 may be expanded as a series in t and a few terms retained, thus avoiding excessive round-off error.

6.4 Time Integrals

Some clarification is necessary on the statement made early in Section 6.2, namely that t (10²⁴ sec) is the time after the previously considered time step. For a burnup carried out so that power is continually adjusted to the same specified level given on the POWER card, t is simply the time between steps specified explicitly by the user, perhaps using a card DT=... For a burnup carried out so that flux and reaction rates remain fixed (LSID=0), the first entry to chain 3 for a given DT card will force a solution of the burnup equations. However, further repeats, as specified on the card by the user, will use the previous analytic solution, and hence t is the time from the last step of previous entry to chain 3. Since the user has presumably requested intermediate steps only to provide output for his own information, rather than to assist the solution of the problem, this possible point of confusion should be borne in mind.

The various time integrals are given simply by the equations below.

$$\begin{aligned} \phi - T(t) &= \phi - T(0) + \phi t(\text{sec}), \\ XX(t) &= XX(0) + [\sum_{\ell} \hat{N}_{\ell} (r_{\ell}^{(4)} - r_{\ell}^{(3)})] t(\text{days}) / D\phi(t), \\ B2AV(t) &= XX(t) / \text{total } t(\text{days}) \\ P - T(t) &= P - T(0) + (\sum_{\ell} \hat{N}_{\ell} f_{\ell} r_{\ell}^{(2)}) t(\text{days}), \\ FIS(t) &= FIS(0) + (\sum_{\ell} \hat{N}_{\ell} r_{\ell}^{(2)}) t(\text{sec}), \end{aligned}$$

FIFA(t) = FIS(t)/ initial fissile atom concentration,

where

ϕ -T = flux time in n/cm²,

B_{2AV} = average estimated buckling to give k_{eff}=1 in cm⁻²,

D = 1-group diffusion coefficient based on \bar{N} in cm
 $= \frac{\sum_{i=1}^{NG} D_i B_i^2 \phi_i}{B_1^2 \sum_{i=1}^{NG} \phi_i}$

P-T = power-time in watt days/cm³,

f_ℓ = energy release/fission of nuclide ℓ in joule/fission,

FIS = fissions/cm³, and

FIFA = fissions/initial fissile atom.

-(11)

6.5 Multiplication Constant, etc.

One-group collapsed macroscopic cross sections for fission emission ($\nu \Sigma_f$), total absorption (Σ_a), and specified leakage (DB^2) are all obtained from the following type of rule

$$\nu \Sigma_f = \frac{\sum_{i=1}^{NG} (\nu \Sigma_f)_i \phi_i}{\sum_{i=1}^{NG} \phi_i} \quad (1)$$

The effective multiplication constant of a homogeneous system is then given by

$$k_{eff} = \nu \Sigma_f / (\Sigma_a + DB^2). \quad (2)$$

In addition the infinite multiplication constant is simply

$$k_{\infty} = \nu \Sigma_f / \Sigma_a \quad (3)$$

For a large system an estimate of leakage necessary to make k_{eff}=1 may be made from Equation 2, namely

$$B_{\infty}^2 = (\nu \Sigma_f - \Sigma_a) / D \quad (4)$$

and the code produces this as output B2 (INF). Similarly Equation 2 can be used to obtain an additional leakage term DB_{eff}^2 which makes k_{eff}=1 when leakage DB^2 has already been included based on a specified group-dependent buckling (5.7), that is

$$B_{eff}^2 = (\nu \Sigma_f - \Sigma_a - DB^2) / D. \quad (5)$$

Since it is possible to obtain group-collapsed transport cross sections, weighted by DB^2 for each group (5.4.2), then, when buckling has not been specified, the code uses $B_i^2 = 10^{-15}$ (i = 1, 2, ..., NG) to prevent the weight from being zero.

For a system which is not large, a change in leakage produces a change in the flux which in turn results in a change in the group-collapsed cross sections, hence the buckling obtained from Equation 4, B_{∞}^2 , may not be particularly accurate. The code, however, has provision for adjusting B^2 to achieve criticality (5.5.2). After each adjustment a further flux calculation is carried out until the specified multiplication is achieved.

Two other quantities produced by the code are extrapolation distance, d, and radius of an equivalent critical sphere, R, obtained by applying Equation 4. The equations used are

$$d = 2.13 \frac{\sum_{i=1}^{NG} (D_i B_i^2)^2 \phi_i / B_1^2}{\sum_{i=1}^{NG} D_i B_i^2 \phi_i} \quad (6)$$

and

$$R = (\pi/B_\infty) - d \quad (B_\infty^2 > 0)$$

$$= 0 \quad (B_\infty^2 \leq 0) \quad (7)$$

For a system with group-independent buckling B_i^2 ($i=1,2,\dots,NG$) adjusted to make $k_{eff} = 1$, then $B_1^2 = B_\infty^2$.

The macroscopic cross sections which are used in the above equations are derived from the cross sections calculated in the first library reading (group flux calculation subroutine, Section 7.2) for a non-burnup user (no FIFA card). For a burnup user the cross sections used are calculated in the second library reading (reaction rates, condensed and mixed cross section calculation subroutine, Section 7.3). The reason is quite simple - a non-burnup user may not want output for all materials that contribute to k_{eff} (5.4.2), hence all the information may not be available from the second library reading. The only possible difference is that the cross sections calculated during the second library reading are corrected for resonance overlap and resonance gross flux correction (7.4.5) whereas the first are not.

A burnup user obtaining output for the full library group structure can obtain an idea of the magnitude of the overlap and gross flux correction from 'total neutron loss rate/source neutron' printed by chain 5 (8), since differences from unity result from the condition that 'total neutron loss rate' is based on the first library reading (no correction) whereas 'source neutron' is based on the second reading (includes correction). When taken together, the two corrections mentioned above to some extent compensate each other as they are corrections of the same form but with different signs.

After burnup, the power density, in watts/cm³, is given by

$$P = \sum_{\ell} N_{\ell} f_{\ell} r_{\ell}^{(2)} \quad (8)$$

using the notation of Section 6.4. Before burnup is calculated an equation similar to (8) above is used to readjust the level of the reaction rates $r_{\ell}^{(k)}$ if the 'constant' power option is set (5.6). Only a burnup user may specify this 'constant' power option.

**CHAIN 4 – GROUP FLUX CALCULATION,
CONDENSING AND MIXING FUNCTIONS**

7. CHAIN 4-- GROUP FLUX CALCULATION, CONDENSING AND MIXING FUNCTIONS

7.0 Summary

A group flux calculation may be carried out for a homogeneous system containing a uniformly distributed neutron source emitting 1 neutron per cm^3 per second. The calculation requires a first reading of the cross section library during which stage a smaller bulk library may be prepared for the second reading if required (5.2.2). Scattering matrices specified (5.2.4) are used in the calculation. Using this group flux, or the flux obtained from other options (5.5.3), group-condensed and nuclide mixed cross sections and 1-group microscopic reaction rates are obtained. This also requires a reading of a cross section library.

7.1 Required Option

As discussed in Sections 5.5 and 6.1, chain 4 is called when the second entry on the time card (LSID) is non-zero, for example,

$$T=0., 20, 0$$

If LSID is positive then the main cross section library is read and a copy is created (5.2.2 and 5.5.3) if required. If LSID is also greater than 1 the flux balance subroutine is entered. Following this the reaction rate, condensing and mixing subroutine is called. On the other hand if LSID is negative, only the reaction rate, condensing and mixing subroutine is called.

7.2 Group Flux Calculation

The specified cross section library (5.2) is read to obtain $\Sigma_a, \nu\Sigma_f, \Sigma_{tr}$ for each group (1 to NG) for the homogeneous mixture to be studied (5.3.4), and the specified scattering matrix library is read to obtain group-to-group transfers $\Sigma_{i \rightarrow j}$ based on the chosen scattering matrices. As mentioned in Section 5.2 it is possible to CREATE a cross section library while reading the main library perhaps to change resonance parameters to shielded cross sections for use with the collapsing procedure (7.3). This procedure is probably only worthwhile for a person requiring burnup using a cross section library containing many resonance parameters. Resonance theories are discussed in Section 7.4.

If the group flux is to be determined from neutron conservation ($LSID > 1$), the equation solved is

$$\underline{W} \underline{\phi} = \underline{X}$$

where \underline{X} is the group fission spectrum (18 and 5.7) normalized so that

$$\sum_{j=1}^{NG} X_j = 1$$

$\underline{\phi}$ is the group flux, and

\underline{W} has elements given by

$$W_{ji} = (AA)_i \delta_{ij} - \sum_{\ell} N_{\ell} \sigma_{i \rightarrow j}^{\ell}$$

where $(AA)_i = \frac{B_1^2}{3(\Sigma_{tr})_i} + (\Sigma_a)_i$,

$\sigma_{i \rightarrow j}^{\ell}$ is the transfer from group i to j for nuclide ℓ , except that

$$\sigma_{i \rightarrow i}^{\ell} = - \sum_{j \neq i} \sigma_{i \rightarrow j}^{\ell}$$

and ℓ is summed over nuclides with scattering matrices given on the SCAT card (5.2.4).

The method used is based on the Gauss-Seidel iteration technique as given by Ralston and Wilf (1960) but advantage is taken of the main features of the flux equation.

For the $K-1$ groups into which no neutrons are scattered from lower energies, that is groups j such that

$$W_{ji} = 0 \quad \text{for } i > j, \quad j = 1, 2, \dots, K-1,$$

the matrix is lower triangular and the solution is immediate. For the remaining groups the equations become

$$\sum_{i=K}^{NG} W_{ji} \phi_i = X_j', \quad j = K, \dots, NG,$$

where X_j' is now a slowing down source.

The iteration scheme used for the n^{th} time applied to solving the preceding equation for $\phi_j^{(n)}$ amounts simply to solving the three equations given below for $\phi_j^{(n)}$. The solutions $\phi_j^{(n)}$ and $\phi_{j+1}^{(n)}$ from the same equations are not obtained, as they are not required at this stage. The three equations are

$$\begin{pmatrix} W_{j,j} & W_{j,j+1} & W_{j,j+2} \\ W_{j+1,j} & W_{j+1,j+1} & W_{j+1,j+2} \\ W_{j+2,j} & W_{j+2,j+1} & W_{j+2,j+2} \end{pmatrix} \begin{pmatrix} \phi_j^{(n)} \\ \phi_{j+1}^{(n)} \\ \phi_{j+2}^{(n)} \end{pmatrix} = \begin{pmatrix} X_j' \\ X_{j+1}' \\ X_{j+2}' \end{pmatrix} - \begin{pmatrix} W_{j,K} & \dots & W_{j,j-1} \\ W_{j+1,K} & \dots & W_{j+1,j-1} \\ W_{j+2,K} & \dots & W_{j+2,j-1} \end{pmatrix} \begin{pmatrix} \phi_K^{(n)} \\ \vdots \\ \phi_{j-1}^{(n)} \end{pmatrix} - \begin{pmatrix} W_{j,j+3} & \dots & W_{j,NG} \\ W_{j+1,j+3} & \dots & W_{j+1,NG} \\ W_{j+2,j+3} & \dots & W_{j+2,NG} \end{pmatrix} \begin{pmatrix} \phi_{j+3}^{(n-1)} \\ \vdots \\ \phi_{NG}^{(n-1)} \end{pmatrix}.$$

A division into 2 regions is made after the first iteration by the requirement that i should be as large as possible, say $i = L-1$, such that

$$\left| \sum_{j=1}^{i-1} W_{ji} \right| < 0.1 W_{ii};$$

that is, the number of upscatters from group $L-1$ is to be just less than one tenth of the total removals from that group. If $NG-L$ is greater than 5, these groups are collapsed into 3 sections such that

$$\sum_i (AA)_i \phi_i^{(1)} \text{ is approximately the same in each section.}$$

Using these three sections, the three equations

$$\sum_k \left(\sum_{i,j}^{(k',k)} W_{ji} \phi_i^{(n)} \right) \lambda^{(k)} = \sum_j^{(k')} X_j',$$

where k and k' have values from 1 to 3 are solved for $\lambda^{(k)}$ after each major iteration to give

$$\phi_i^{(n')} = \lambda^{(k)} \phi_i^{(n)}.$$

The initial estimate $\phi^{(0)}$ is either a Westcott flux (13.2) or the group flux from a previous calculation.

The accuracy criterion used for convergence is

$$\sum_{j=1}^{NG} (AA)_j | \phi_j^{(n)} - \phi_j^{(n-1)} | \leq \epsilon ,$$

where ϵ , usually 10^{-4} , may be set in the input data (5.8). The number of iterations required depends mainly on the ratio of absorption to scatters in each group, and $\phi_j^{(0)}$, and usually ranges between 4 and 7 for the basic 120-group libraries (19).

7.3 Condensing and Mixing Functions

7.3.1 Reaction rates

The burnup of a nuclide depends on microscopic reaction rates (6.2). For this reason if burnup is required (5.3.9), reaction rates for all library nuclides (NN) are calculated. If burnup is not required, only reaction rates of those nuclides requested for output (5.4.2) are obtained. The six reactions treated are those obtained from the cross section library (18), namely

- (1) (n,capture 1)- usually (n, γ)
- (2) (n,fission)
- (3) (n,absorption)
- (4) (n, ν fission emission)
- (5) (n, transport)
- (6) (n,capture 2)- usually (n,2n).

Using a superscript to indicate the particular reaction and subscripts to indicate in order nuclide and group, we have

$$r_{\ell}^{(k)} = \sum_{i=1}^{NG} \sigma_{\ell i}^{(k)} \phi_i \quad (k=1,2,3,4,6), \quad (1)$$

$$r_{\ell}^{(5)} = \left(\sum_{i=1}^{NG} \sigma_{\ell i}^{(5)} D_i B_i^2 \phi_i \right) \left(\sum_{i=1}^{NG} \phi_i / \sum_{i=1}^{NG} D_i B_i^2 \phi_i \right) , \quad (2)$$

I = 1, standard option,

or
$$r_{\ell}^{(5)} = \sum_{i=1}^{NG} \sigma_{\ell i}^{(5)} \phi_i , \quad I = 0 ,$$

where ϕ_i is the group flux, which is calculated from neutron balance (7.2), read into the programme (5.7), or calculated from a Westcott flux (13.2),

D_i is the group diffusion coefficient $\approx 1/3 (\sum_{tr})_i$,

$(\sum_{tr})_i$ is the group transport cross section for the specified mixture,

B_i^2 is the group buckling specified by the user (5.7), and

I is the option required with transport averaging - the second digit on REQD card (5.4.2).

$$\sigma_a = (1-p) \bar{\xi} \sigma_p / p \delta u \quad (\text{group resonance cross section}), \quad (1)$$

where δu = width of resonance group in lethargy units,

$$p = e^{-\rho I / \bar{\xi}} \sigma_p \quad (\text{resonance escape probability}), \quad (2)$$

$$\rho = 2/(1+p), \quad (3)$$

$$I = \frac{\Gamma_a \sigma_o}{E_r} a_\lambda J(\theta, a_\lambda) \quad (\text{resonance integral}), \quad (4)$$

$$\lambda_\ell = 1 - \alpha_\ell \frac{(C_1 + C_\lambda)}{x_\ell} \arctan \left(\frac{x_\ell}{C_1 + \alpha_\ell C_\lambda} \right), \quad (5)$$

$$x_\ell = 2E_r (1 - \alpha_\ell) / \Gamma, \quad (6)$$

$$\sigma_{p\lambda} = \sum_{\ell=1}^k \lambda_\ell \sigma_{p\ell} \quad \left(\begin{array}{l} k = \text{resonance nuclide} \\ \ell = 1 \text{ to } k-1 = \text{light nuclides} \end{array} \right) \quad (7)$$

$$\sigma_{p\ell} = N_\ell \sigma_{s\ell} / N_k \quad (8)$$

$$a_\lambda = \frac{\sigma_{p\lambda}}{\sigma_o} \frac{\Gamma}{\Gamma_a + \lambda_k \Gamma_n}, \quad (9)$$

and $C_\lambda = (1 + 1/a_\lambda)^{1/2}. \quad (10)$

The function $J(\theta, \beta)$ used in GYMEA is an approximation proposed by Doherty (1963).

Equations 5 to 10 and 4 and 3 are simply solved using one iteration, and numerical accuracy is retained in Equation 1 using an expanded form based on Equation 2 for $(1-p) < 10^{-4}$. Other cross sections are obtained from σ_a of Equation 1 using ratios of resonance widths, for example,

$$\sigma_{n,\gamma} = \frac{\Gamma_\gamma \sigma_a}{\Gamma_a} + \sigma_{n,\gamma}^* \quad \text{resonance removed}.$$

The entry held in the library (18) as cross section (3) – normally absorption – is not used in the resonance region for nuclides with resonance parameters or shielding coefficients given, as cross sections (1) and (2) are summed to give the required cross section (3).

7.4.3 Cross sections from shielding coefficients

Following Keane (1964b) it can be seen that for zero temperature the equations of Section 7.4.2 give approximately (for one resonance)

$$\sigma_a = \sigma_I / (1 + \sigma_o / \sigma_p)^{1/2}, \quad (1)$$

where $\sigma_I = \pi \Gamma_a \sigma_o / 2E_r \delta u \quad (2)$

= infinitely dilute cross section.

Values of the concentration variation (σ_p variation) of a cross section can therefore be anticipated for several resonances in a group, and not necessarily at zero temperature, using

$$\sigma_x = \sigma_\omega / (1 + \frac{A}{\sigma_p} + \frac{B}{\sigma_p^2})^{1/2} \quad (x = \gamma, f, \text{ and } \nu f). \quad (3)$$

Three parameters σ_∞ , A, and B have been introduced in Equation 3 and these are determined from three calculations of σ_x for three different σ_p . In all, nine parameters are determined for three reactions.

Within this context σ_∞ need not agree with σ_I , summed over resonances of the group. The three σ_p chosen in calculating σ_∞ , A, and B would be based on the anticipated range of concentration variation of the nuclide under study for either a burnup study or a survey. For variation over a range of σ_p , typically, 10^4 , 5×10^4 , 5×10^5 , Equation 3 has been found by the authors to be remarkably satisfactory. This point needs further verification however. It should be remembered from Section 5.3.8 that resonance parameters can be automatically turned into shielding coefficients using the CREATE feature (5.2.2).

The σ_p used in Equation 3, unlike its counterpart in resonance theory (7.4.2), does not include the potential scattering cross section of the resonance nuclide— this convention is used throughout shielding theory.

When using shielding coefficients the temperature is fixed to the value chosen when preparing these parameters.

7.4.4 Particle self shielding

Section 5.3.7 introduced a basis for allowing for particle self shielding following Keane (1964a). This method modifies the potential scattering cross section of light nuclides thus

$$\sigma_s^* = \sigma_s / (1 + \bar{\ell} \sum_p V_m / V_p) .$$

The modified potential scattering cross sections, σ_s^* , are used both in resonance theory (7.4.2) and shielding theory (7.4.3).

7.4.5 Resonance overlap

Following the option set with a REQD card (5.4.2), it is possible to allow for overlap of resonances using the method of Rowlands (1963) investigated by Keane (1966). In addition the resonance gross flux correction factors investigated by Keane and Pollard (1966) are combined with the overlap corrections, assuming these two factors to be independent. In GYMEA the method amounts to correcting cross sections (1) to (4) for all nuclides for the resonance groups using the factor

$$C_{\ell i} = 1 - \left((\sum_a)_i - \sum_{\ell}^{(3)} \right) h / \sum_p , \tag{1}$$

then

$$\sigma_{\ell i}^{(k)*} = \sigma_{\ell i}^{(k)} C_{\ell i} \quad (k = 1, 2, 3, 4), \tag{2}$$

where

$$(\sum_a)_i = \sum_{\ell=1}^{NN} \sum_{\ell i}^{(3)} \tag{3}$$

= total absorption cross section for the i^{th} group,

$\sum_{\ell i}^{(3)}$ = absorption cross section for the ℓ^{th} nuclide,

\sum_p = total potential scattering cross section, and

$$h = 1 - \delta u / 2\bar{\xi} .$$

Allowance for overlap and gross flux correction is only possible if the vector $(\sum_a)_i$ is available from a first reading of the library, or is set by the user in COMMON (12), otherwise the programme will change the option. Normally this means that

REQD 01

should precede

T=0.,-1,0

In addition, the storage assigned to $(\Sigma_a)_i$ is shared with information used in the burnup calculation; hence after a burnup calculation, $(\Sigma_a)_i$ should be regenerated if overlap and gross flux correction is required, using, say

DT=100., 20 (NOTE-POSITIVE), 0

**CHAIN 5 - OUTPUT FOR MOST
NEUTRONICS CODES AND
PRINTED OUTPUT**

8

8. CHAIN 5 - OUTPUT FOR MOST NEUTRONICS CODES AND PRINTED OUTPUT

8.0 Summary

Chain 5 produces output on cards suitable for input to the codes CRAM, 4-ZOOM, DSN, WDSN, as well as printed multigroup cross-section and scattering matrix data (PRINTX and PRINTS respectively). The required options are set by use of the card

ØUTPUT ...

Each of the code names in format (A6) appearing on this card, for example CRAMbb(b=blank), may be followed by a fixed point number. These numbers are used for PRINTX, CRAM, and DSN. The specific meaning of each is given below.

8.1 Condensing and Mixing

The condensed and mixed cross sections calculated in Chain 4 and stored in the array WH are used here. However the condensing and mixing of scattering matrices occur in this chain. The method of condensing the scattering matrices is given by the equation

$$\sigma_{m \rightarrow n} = \frac{\sum_i \sum_j \phi_i \sigma_{i \rightarrow j}}{\sum_i \phi_i}$$

where i is summed over the m^{th} set of groups and j over the n^{th} set of groups as defined by the GROUPS card.

For example the following equation applies for the 2 groups 1 to 70, 71 to 120:

$$\sigma_{1 \rightarrow 2} = \frac{\sum_{i=1}^{70} \sum_{j=71}^{120} \phi_i \sigma_{i \rightarrow j}}{\sum_{i=1}^{70} \phi_i}$$

The mixing of scattering matrices follows the rules specified by the user. However, one point needs to be stressed regarding the use of a compound library scattering matrix. For example, if

SCAT 900. 'BEO' CRY bbb' '

is used with 'BEO' = 1.BE9, 1.O16, 0.' '

then neither BE9 nor O16 has a scattering matrix associated with it. The same is true of the usual macroscopic material given by

'MIX' = 0. (1,70) 0.' '

The type of card which should be used in this case is

'MIX' = 0. 'BEO', 0. 'O16', 0. (3,70) 0.' '

where 'O16' is defined as some portion of O16 not included in 'BEO' - see Section 5.3.2 for further details.

8.2 Special Materials

So that different output materials may be produced for different codes during the one calculation, the four material labels 'CRAM', 'ZOOM', 'DSN', and 'WDSN' are given a special meaning. Thus, if a material for output is labelled 'CRAM', it will not be produced as output for the other codes, and output for CRAM will consist only of this material.

8.3 Printed Output

Printed output of both PRINTX and PRINTS is weighted according to the value of X(6) as described in Section 5.4.1. The title given by the HEAD card (5.2.2) is printed at the beginning

of each block of data, so that it usually heads the data for each material. The character in column 12 of the HEAD card is used as the carriage control character when the heading is printed. The use of this feature may considerably improve the layout of the data.

8.3.1 PRINTX

PRINTX output consists of weighted multigroup cross sections for the output materials printed in a readable form. Normally the data is printed as group information for each nuclide, but, if the digit 1 follows PRINTX on the MATS card, material information for each group is produced. The output is always terminated by a table of group summary information.

8.3.2 PRINTS

PRINTS output consists of a print-out of the condensed scattering matrices. The matrices are printed as all outscatters from each group with the self-scatter term equal to the negative sum of the outscatters.

8.4 Data for Other Computer Codes

Condensed cross section data following the user's requirement (5.4.2) is produced on cards for the 4 computer codes CRAM, 4-ZOOM, DSN, and WDSN. To delimit a set of data for a code, 3 title cards are punched at the beginning and a distinctive card is punched at the end. Labelling of cards in columns 73 to 80 is used to a variable extent. No action is taken to bring the data into line with the particular restrictions of the code versions used at the A.A.E.C. For example, no truncation of scattering matrices for DSN is carried out. Only for 4-ZOOM was it found necessary to modify the data for the code.

8.4.1 CRAM

Output for CRAM consists of:

- (1) The fission spectrum for each group.
- (2) The following information for each material, where each cross section is given as a block for all groups i :

$$\sigma_{tr}, \sigma_r, \nu\sigma_f, (\sigma_{i \rightarrow j}, j = 1, \text{NGD})$$

and the optional information

$$\sigma_c, \sigma_f, \sigma_{n, 2n}$$

Here

$$\sigma_{i \rightarrow i} \text{ is set to zero}$$

$$\text{and } \sigma_r^i = \sigma_a^i + \sum_{j=1}^{\text{NGD}} \sigma_{i \rightarrow j}$$

Normally none of the optional cross sections are produced, but if the name CRAM is followed by a fixed point number ℓ , then extra cross sections are produced. Thus, if $\ell=1$, only σ_c is produced, if $\ell=2$, σ_c and σ_f are produced, etc. The format used is (1P5E13.5).

8.4.2 4-ZOOM

For a 4-ZOOM, $(\sigma_a - \nu\sigma_f)$ and $\nu\sigma_f$ are produced to be used with a ν value of 1.

Output for 4-ZOOM consists of:

- (1) The fission spectrum for each group
- (2) The average energy of each group
- (3) The following information for each material as in Section 8.4.1:

$$(\sigma_a - \nu\sigma_f), \nu\sigma_f, \sigma_{tr}, (\sigma_{tr} - \sigma_a), (\mu_{j \rightarrow i}, j = 1, \text{NGD})$$

where

$$\mu_{j \rightarrow i} = \begin{cases} \sigma_{j \rightarrow i} / \sigma_{ss}^j & i \neq j \\ (\sigma_{ss}^j - \sum_{k \neq j} \sigma_{j \rightarrow k}) / \sigma_{ss}^j & i = j \end{cases}$$

and

$$\sigma_{ss}^j = \sigma_{tr}^j - \sigma_a^j .$$

The format used is (1P6E12.4) except for the scattering matrices of heavy materials where the 4-ZOOM repeat notation is used.

8.4.3 DSN

For DSN a fixed point number of the form llkkn may follow the label DSN. Here ll is the table length for output, kk is the position of the self-scatter term in the table and n is the number of optional cross sections required. If ll, kk are not given, the programme calculates the minimum required values. n is set to zero if no value is given. The programme checks to see if the data will fit into the required table.

The output for DSN consists of the G00, K7, K6, V1, C0 vectors, with the K7 vector containing the fission spectrum, the velocities in V1 (cm/10⁻⁸sec) condensed inversely, and the K6 vector containing the group flux. The C0 vector contains for each material the following table for each group i :

$$\sigma_c^i, \sigma_f^i, \sigma_{n,2n}^i, \nu\sigma_f^i, \sigma_{tr}^i, \sigma_{i+k \rightarrow i}, \dots, \sigma_{i \rightarrow i}, \dots, \sigma_{(i-l) \rightarrow i}$$

where k and l have the maximum values required and

$$\sigma_{i \rightarrow i} = \sigma_{tr}^i - \sigma_a^i - \sum_{j \neq i} \sigma_{i \rightarrow j} .$$

The first three cross sections are optional and are deleted in the order $\sigma_{n,2n}, \sigma_f, \sigma_c$. The format used is (1P6E12.5).

8.4.4 WDSN

The output for WDSN must be specified by the user to be macroscopic and, for each material, consists of the table indicated below, with entries for each group i .

$$LV, PSS, 1., \nu\sigma_f^i, \sigma_a^i, (\sigma_{i \rightarrow j}, j = \ell, k)$$

where LV is the length of the table excluding the first 4 entries; PSS is the position of self scatter, l and k are the appropriate numbers for group i to give all non-zero values of $\sigma_{i \rightarrow j}$, and $\sigma_{i \rightarrow i}$ is as in 8.4.3. The format used is (2I6,1P2E12.4/(6E12.4)).

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9. CHAIN 6 - OUTPUT FOR GYMEA LIBRARY

9.0 Summary

The preparation of libraries with smaller bulk than the main libraries (19) is frequently desirable to save machine time when a burnup survey of several systems is being carried out. The bulk may be reduced by mixing fission products and condensing the cross section data. For this reason the code GYMEA has available facilities for preparing its own smaller bulk libraries. The cross section library application of deleting nuclides, with perhaps resonance mode changed, is best carried out using the CREATE feature (5.2.2). Additions to the libraries can be carried out using the library editing programmes (19.4).

9.1 Required Data

Two data cards which must be specified for the preparation of condensed and possibly mixed cross section GYMEA libraries are, for example,

```
PUT 3  XS01  0  SC01

OUTPUT 3  PRINTX  CRAM  GYMEA
```

The first specifies that tape 3 is to be created as a cross section library to be called XS01 and tape 0 is to be created as a scattering matrix library to be called SC01 (5.2.3). In addition to these cards, data must be given for

```
MATS    ...

GROUPS  ...

HEAD    ...
```

as with any other output request. Either entry for tape specification on the PUT card may be negative if only one library is to be prepared. Should both entries be the same, then, between preparation of the cross section library and the scattering matrix library, the first tape will be unloaded and a message typed advising the operator to mount a tape for creation of a GYMEA scattering matrix library. The user should type a message before chain 6 is called, using a compiler feature (10) to nominate the required tape, for example,

```
TYPE      WHEN ADVISED MOUNT TAPE 300 ON 3

T=0.,20,0
```

A fixed point number may follow GYMEA on the OUTPUT card. This number sets options detailed in Section 9.2, for example,

```
OUTPUT 1  GYMEA  1
```

9.2 Cross Section Library

9.2.1 Library structure

A condensed and/or mixed cross section library is prepared for nuclides in the order specified on the MATS card. Since this order is important in burnup studies (6.2) care should be used when preparing the MATS card. The data created is the same type as that given in the specified library, that is, if a nuclide on the main library has resonance parameters given then the created tape also has resonance parameters for the nuclide under discussion.

For the off-resonance region (18) condensed data produced by chain 3 (7.3) is used, whereas for the resonance region, resonance parameters or shielding coefficients are extracted from the specified library. For this reason mixing is not permitted in the preparation of a GYMEA library

except for nuclides which do not have resonance information as part of the main library (IS ≤ 0 Section 5.2.2). Resonance parameters are simply stored (λ 100 resonances) for a group (unless the group is not collapsed) and then created on tape. On the other hand shielding parameters are used to reconstruct three cross sections, using the arrays $\tilde{N}_1, \tilde{N}_2, \tilde{N}_3$ (5.3.8), which are collapsed in the usual fashion (7.3.2) and then reconverted to shielding parameters. In both cases the entries held in the library (resonance-removed cross sections) are also collapsed in the standard way.

Using the optional data that may be given after GYMEA on the output card it is possible to select different resonance group boundaries, provided they are 'inside' boundaries which would result using the specified library resonance groups, and to select a library resulting from CREATE (5.2.2) rather than the main library. Hence if the main library contains resonance parameters for a nuclide we can CREATE a library with shielding parameters, and this structure can be used in the condensed library preparation. The optional fixed point data word has the form

S L L L U U U B

the digits of which have the interpretation:

- S = sign - if negative no resonance information is created,
- L L L = lower resonance group boundary in terms of condensed group structure,
- U U U = upper resonance group boundary in terms of condensed group structure unless
- U U U = 0, in which case the structure of the main library determines the resonance group boundaries,
- B = 0 - the main library is used,
- B = 1 - the created library is used.

As an example, suppose the main library consists of 12 groups, of which 3 to 10 are designated in the library heading information as being resonance groups (18), and the card used (5.7) is:

GROUPS 6,1,3,4,6,8,9,12

that is R R R R R R R R

1 1 2 3 4 5 6 7 8 9 10 11 12

1 1 2 3 4 5 6

Then, in terms of the condensed group structure, groups 2 to 5 would be considered to be resonance groups since groups 1 and 6 contain only partial resonance information for the full range. If then

- (i) S L L L U U U B = 0 (this is implied if the data word is not specified), the resonance groups would be 2 to 5,
- (ii) S L L L U U U B = 30040
the resonance groups would be 3 to 4

or (iii) S L L L U U U B = 10051
an error condition would result as the first group is not a full resonance group.

9.2.2 Burnup table

The burnup table for the mixed nuclide library is obtained from the main library burnup table as follows.

- (i) For a library nuclide the table is simply copied.
- (ii) For a defined material the burnup and decay scheme of the dominant nuclide (the first nuclide of a definition - Section 5.3.3) are used except that the decay constant and fission yields are each calculated as an aggregate for all component materials.

The aggregate for decay constant is based on the mixing rules used in the preparation of cross sections for the material (5.3.2) whereas the aggregate for fission yields may be weighted in the same fashion or weighted differently using the following feature. If the mixing rule for a material is given in terms of elements of the Z array, for example,

$$'F2A' = -1 \text{ MO95}, -3 \text{ TC99}, -5 \text{ RU101}, 0'$$

then in the preparation of 'F2A' cross sections and decay constant a weight Z(1) would be associated with Mo⁹⁵, Z(3) with Tc⁹⁹, and Z(5) with Ru¹⁰¹, whereas in the preparation of fission yields a weight Z(2) would be associated with Mo⁹⁵, Z(4) with Tc⁹⁹, and Z(6) with Ru¹⁰¹.

For this application a mixing rule containing a zero coefficient, for example,

$$'MIX' = 0. (1,55), 0.'$$

is not given a special meaning as it is for cross section mixing (5.3.2), but is taken to mean zero proportion of the specified nuclides. In general it may be necessary to prepare a burnup table beyond the scope of these features, in which case the library may be corrected later using the library edit programme, EDITOR (Ford 1966 - Section 19.4).

9.3 Scattering Matrix Library

Mixing rules are ignored in this application as a full main library complement (18) is always involved. Each matrix of the library for the specified temperature is condensed to the required groups (5.4.2) and the condensed library is copied to the specified tape (9.1). If the specified temperature is the first of the library set, not necessarily the smallest temperature, then the tape is rewound and a heading record is created. For other than the first temperature, the tape is not rewound and the condensed scattering matrices are created on the tape without a further heading record. Using these features it is possible to create a condensed scattering matrix library with scattering matrices for each temperature averaged over an appropriate flux spectrum. If the full set of scattering matrices or full set of temperatures is not required, a scattering matrix COPY should first be produced (5.2.2).

10. CHAIN 1 - COMPILER AND ERROR ROUTINES

10.0 Summary

The ability to calculate quantities of direct interest in-situ in GYMEA is a useful feature which enables people not necessarily involved in writing the code to extend the code to carry out many different calculations. It has been found that the mixing of macro-data and compiler instructions greatly adds to the versatility of GYMEA.

10.1 The Compiler

This routine is used to read data cards into core, to 'compile' some of them, to write them out on unit 4 as one binary file for execution in CHAIN 2, and finally to set the core to standard state (11).

The process of compilation consists of:

- (1) All transfer points are set to absolute form.
- (2) Most of the FORTRAN-like cards are interpreted and changed into a set of words describing the function of the statement. These words are stored in place of the original card for later execution. For completeness a description of all the FORTRAN-like cards (GYMEA code) follows.

GYMEA code cards are all executable and may be mixed with the other input cards as desired. The data may be punched anywhere on a card but on one card only. Statement numbers are indicated by a slash (/) and are the first characters on a card. They may appear on any of the input cards except END and CEASE; for example,

/ 2 T = 0., 20, 0

10.1.1 Constants and variables

Any integer or floating point constant acceptable to FREE input may be used (22). There are only three variables available, all subscripted, namely, X(n), Z(n), J(n), dimensioned 2000, 860, and 9 respectively. As well there are the special variables X(o), Z(o), and J(o). X and Z are floating point variables and J is a fixed point variable. The subscript n may be any integer, positive or negative. A negative subscript has the meaning X(-n) = X(J(n)). Thus J(n) is referred to as -n when used as a subscript or a DO loop index.

10.1.2 Arithmetic statements

An expression consists of one or two variables or constants separated by one of the four operators: +, -, *, /. However an operation on two constants is invalid. An expression is either of fixed or floating point mode; mixed mode expressions are not allowed. A negative sign may precede an expression.

A statement is then of the form

a = b

where a is a variable and b is an expression. a and b may be of different mode, the expression being evaluated in its mode and converted to the other if required.

Examples are:

X(2) = Z(1) * 1.3
 Z(10) = -J(9) + J(-2)
 J(5) = 7
 Z(-5) = -Z(0)

10.1.3 Control statements

The main control statements are:

- GO TO The statement GO TO n , where n is a statement number, transfers control to the card with statement number n . The statement GO TO $J(0)$ means: transfer to the card which has the card number of the present card plus $J(0)$. This provides a feature equivalent to a computed GO TO statement.
- IF The statement IF (V) n_1, n_2, n_3 transfers control to the statement number n_1, n_2 , or n_3 , depending on whether the fixed or floating point variable V is negative, zero, or positive.
- DO The statement DO $n -i_1 = i_2, i_3, i_4$ where n is a statement number and i_1, \dots, i_4 are positive integers, executes repeatedly the statements up to that numbered n , with $J(i_1)$ having the value i_2 incremented in steps of i_4 until i_3 is reached. The integer i_4 must always be given and each DO loop must end with a separate CONTINUE statement.

Other control statements are:

- CONTINUE a dummy statement used mainly to end a DO loop,
- PAUSE causes the machine to halt for operator action,
- END must be the last physical statement of each job or subroutine. Also terminates a job.

The following example is a sample portion of a programme illustrating control statements.

```
IF (X(-1)) 1,2,2
/ 1 DO 3 -2 = 1,10,1
    DO 4 -3 = 1,5,1
    J (4)   = J(2) + J(3)
    Z (-4)  = Z(-2) * X(-1)
/ 4 CONTINUE
/ 3 CONTINUE
    GO TO 5
/ 2 X(-1)   = -X(-1) + Z(10)
    GO TO 1
/ 5 CONTINUE
```

10.1.4 Input-output statements

Input-output of numeric (floating point only) data differs more from FORTRAN than other statements do. The first 2 or 3 characters indicate the type of statement and only the numeric information following is relevant. Here W,P,R stand for print, punch, and read respectively.

Where only 2 letters are given the statement is of DO loop form, for example,

WX (X(I), I = 1, 10, 1)

RZ 1 100 2

Where the third letter is S, the first number is the number of single variables following, for example,

WZS 3 ONLY Z(0), Z(5), Z(-3)

PXS 5 -1, -2, -3, -4, -5

A maximum of 9 variables can be processed this way. If the first number is zero, input-output of a card image columns 1-72 is achieved. Two card images can be read and written this way: one by using X, the other by using Z. Thus

RXS 0

WXS 0

would be used to read and print a card.

The standard FORMAT for output is (1P5E14.5). This can be changed by using a format card

FORMAT()

where the left bracket must be in column 13. The FORTRAN variable-format feature is used to effect this. The given FORMAT is then the one used until replaced.

When a read card is used, data is read from the card reader under the standard form of free input (22); no continuation integer is required in column 1 for this application. This data for any job must begin with the card

,DATA JOBN

(when JOBN is the job name) which follows the CEASE card. The read feature is useful to enter blocks of data for which no macro-card exists.

There are also four cards for writing messages:

WRITE1 comments

PUNCH "

TYPE "

WRITE2 "

The first 3 output columns 13-72 of the card on the specified device in columns 1-60 and the fourth writes columns 13-72 of the card and columns 1-72 of the next card on the printer in columns 1-132. Note that column 13 of these cards is used for printer control.

The statement UNLOAD n causes the tape on unit n to be physically unloaded.

10.1.5 Subroutines

For the purpose of GYMEA, there are three types of subroutine:

- (1) Internal— written in FORTRAN or MAP originally and part of the system.
- (2) External— written in GYMEA code by the user as input.
- (3) Replaceable— internal versions are available, but if the user enters his own version that will be used for the run.

An external subroutine is a non-argumented subroutine which may include any of the input cards of the code except another SUBROUTINE card. The return from a subroutine, provided by a RETURN card, is to the next input card after that being executed when the subroutine was called. The indices J(1),..., J(9) are saved on entry and restored on return. The last physical statement of a subroutine must be END.

External subroutines are normally loaded first, but they may be loaded between or after the jobs. If two subroutines of the same name are loaded, the last one is used. Subroutines may be nested up to ten deep.

Subroutines are called by

CALL NAME

where only the first 3 letters of NAME are significant. Subroutines which are part of the system and may be called are:

- EXP sets $Z(0) = \text{EXP} (X(0))$,
- LOG sets $Z(0) = \text{ALOG} (X(0))$, that is \log_e ,
- SEND terminates a job,
- SCEASE terminates all jobs in a load,
- EXIT returns control to the system monitor,
- CLOCK sets $Z(0)$ equal to the running time of the code in minutes,
- DAT called when a data error occurs, calls SEND
- LIB called when a library error occurs, calls SEND
- FINISH called on internal timer overflow, calls DUMP, then calls EXIT,
- STAN sets storage to standard state (11),
- DUMP prints out all non-zero J, X, Z variables,
- EQU carries out equilibrium search (13.1),
- WEST computes a Westcott group flux (13.2),
- NS1 transfers data to shielding array 1 (5.3.8),
- NS2 transfers data to shielding array 2,
- NS3 transfers data to shielding array 3,
- NUC normalizes defined material concentrations (5.3.5),

OUTPUT writes out the output set up previously, if this was retained prior to this card using X(9) ≤ -10.(5.4.1).

DAT, LIB, FINISH, and SEND are the subroutines used by the code that may be replaced.

Examples:

- 1. To set $Z(59) = e^{5.5}$:

```
X(0) = 5.5
CALL EXP
Z(59) = Z(0)
```

- 2. To print the machine time used:

```
SUBROUTINE TIME
CALL CLOCK
FORMAT (15 H0TIME TAKEN WAS F6.2,5H MINS)
WZS 1 Z(0)
RETURN
END
```

CALL TIME may then appear anywhere in a job or subroutine.

- 3. To obtain a dump before termination when the programme detects a data error:

```
SUBROUTINE DAT
CALL DUMP
CALL SEND
END
```

- 4. To check the tape on a library error:

```
SUBROUTINE LIB
TYPEbbbCHECK TAPE ON 1, MUST BE GYMEA NDXD
PAUSE
USE 1, 70, -1
RETURN
END
```

10.2 Errors

This chain is called when any error is detected by the programme. All libraries are rewound and the input cards read back into storage, and a self-explanatory message is then printed. The bulk of the 68 different error messages is one of the main reasons for the inclusion of all errors in this small chain.

If an ERROR card has been given as data (13) the appropriate sections of common storage are dumped out in G format (interpreted as E mode for numbers with a non-zero exponent and I mode otherwise). This should only be used for programme debugging. The programme next calls CHAIN 2 for termination or error recovery.

Errors detected during compilation are treated differently. Instead of the procedure described above, a message is written and the card is not compiled. It is not until an attempt is made to execute the card that the job is terminated.

In relation to termination, it should be remembered that by having subroutines DAT or LIB as input the user can obtain his own termination or recovery procedures. For instance, a check can be made on all the cards of a large job without performing the calculations if an error occurs by using

SUBROUTINE DAT

RETURN

END

11. STANDARD STATE (STAN)

The common area of storage is set to its standard state after each load is compiled, and whenever requested, by the card

CALL STAN

This should always be the first card of a job that is independent of previous jobs to give protection from an error in the previous jobs. It should also always precede a USE card if the requested libraries have a different structure from libraries used previously.

In the standard state the data areas (apart from WH, which contains the input cards, and a few variables given below) are set to zero. No information pertaining to any library, or any information previously calculated, is retained.

The variables not set to zero and the values they are given are:

LIB = -1)
)
 LIBC = -1) no libraries have been specified,
)
 LIBS = -1)
)
 LXG = -1)
)
 LSG = -1)

KTR = 11 overlap and transport options each set to 1,

STR = 1.E-4 accuracy of absorption required in group flux calculation,

SXA = -1. terminate the job if B2AV falls below -1,

IFP = -1 the programme is set to normalize the group flux to unit source when it is calculated,

IWT = -1 reaction rates must be calculated before burnup is attempted,

X(0) is unchanged so that the user can test for previous errors,

FORMAT for GYMEA code output is set to (1P5E14.5)

and HEAD for neutronics output is set to blanks.

12. USER'S COMMON

12.0 Summary

It is possible for the user of GYMEA to use or modify certain variables held in COMMON storage in areas readily accessible using the compiler variables J(0)–J(9), X(0)–X(2000), Z(0)–Z(860) (10.1.1). For example, we may modify the potential scattering cross section of the first library nuclide thus

$$X(351) = 3.0$$

A list follows which gives entries in the user's COMMON in terms of variables of the source language programme (21) and names associated with printed output.

12.1 Index Registers J(0) – J(9)

The variables J(0) to J(9) are not used in the source programme. These variables enable data to be processed using ordered arrays; for example, to set Z(1) to Z(100) to zero we may use

$$DO9 -1 = 1,100,1$$

$$Z(-1) = 0.$$

/ 9 CONTINUE

Here Z(-1) is interpreted as Z(J(1)).

12.2 Special Registers X(0) – X(10)

The variables X(0) to X(10) are used as subroutine arguments and they obey the same rules as do variables appended to a normal FORTRAN function or subroutine-calling statement, for example to calculate e^t ($t=X(11)$) we may write

$$X(0) = X(11)$$

CALL EXP

The result of the calculation is then available using the variable Z(0). In addition to this function these variables are used to control certain aspects of the programme indicated below.

X(0) =1. on error exit from a job, 0. normal exit,

X(1) to X(4) not used at present,

X(5) =1. programme uses average concentrations (AVE.N) in determining k_{eff} , etc.,

0. programme uses final concentrations (AFT.N),

X(6) sets type of printed multigroup condensed output (5.4.1),

X(7) used in conjunction with CALL EQU (13.1),

X(8) used as a weighted coefficient for data accumulated with SAVE (5.4.3),

X(9) sets type of printed one-group condensed output (5.4.1),

X(10) used by the programme for burnup control.

12.3 Result Registers X(11) - X(1999)

<u>X</u>	<u>Coding</u>	<u>Print Out</u>	<u>Remarks</u>
11	TA	T(DAYS)	time after burnup in days
12	DTD		difference in time between current and previous time step in days
13	AFL	FLUX	flux level in n/cm ² sec
14	APR	POWER,AFTER	power after burnup in watts/cm ³
15	RPR	POWER,SPEC.	specified power used as input
16	VPR	POWER,AVERAGE	average power for the time step
17	BPR	POWER,BEFORE	power at the beginning of the time step
18	FIN	FLUX-T	flux-time integral in n/cm ²
19	PIN	POWER-T	power-time integral in watt-days/cm ³
20	VIN		total fissions /cm ³
21	XIN		X-time integral in days/cm ² with X = ($\nu \Sigma_f - \Sigma_a$)/D
22	FIFA	FIFA	fissions/initial fissile atom = VIN / FIFC
23	AKI	K(INF)	k _∞
24	AKE	K(EFF)	k _{eff}
25	XI	B2(INF)	estimated buckling to give k _{eff} = 1 in cm ⁻²
26	XE	B2(EFF)	estimated buckling to give k _{eff} = 1 including DB ² term, from B ² defined by the user, in cm ⁻²
27	XAV	B2AV	time average buckling in cm ⁻²
28	CNA	A/S	neutrons absorbed/source neutron (should be close to 1. if the group flux is calculated from neutron balance)
29	SKI		KINF LIMIT for job exit (CALL SEND)
30	SKE		KEFF LIMIT for job exit
31	SXA		AVX LIMIT (X(27)) for job exit
32	STR		ACC IN ABS. RATE (5.9)

33	GL		lower eigenvalue limit
34	GG		second guess at eigenvalue
35	GU		upper eigenvalue limit
36	AER		required k_{eff} (or FIFA (13.1))
37	ACER		required accuracy in k_{eff} (or FIFA)
38	ACET		outer iteration accuracy in power (13.1)
39	FIFC		initial fissile atom concentration
40	TEMP		resonance temperature (degrees Kelvin)
41	TLAST		if T = -1,20,0 is used then X(41) replaces -1
42	BPRS		previous BPR (GYMEA modifies BPR)
43	XY(1)		$\Sigma_a + DB^2$
44	XY(2)		$\nu \Sigma_f$
45	XY(3)		DB^2
46	EPART		\bar{l} used in resonance shielding
47	VOID		void used to normalize nuclide concentrations
48	ETC1(1)		B2AV saved from group flux calculation EQU search (13.1)
49	ETC1(2)		number of file skips through tape 4 before reading or writing cross section data (13.3)
50	ETC1(3)		not used
51	AN	AFT.N	nuclide concentrations at the end of the time step (nuclides/ 10^{-24} cm ³)
150			
151	AV	AVE.N	average nuclide concentrations for the time step (the numbers printed above AFT.N and BEF.N refer to the extremes of time in days for the average concentrations)
250			
251	BN	BEF.N	for constant power runs this is the N's of the previous step and for constant flux runs this is the initial fuel concentrations
350			
351	SXS		potential scattering cross sections for each library nuclide
450			
451	FER		fission energy release for each nuclide (joule/fission)
550			
551	GFLUX		group flux
670			
671	GSCE		group source (the programme normalizes this to unit sum)
790			
791	GB2		group buckling
910			

911 to 1030	WU and AA	$\Sigma_a + DB^2$ for each group, or burnup matrix terms if group flux calculation does not follow burnup
1031 to 1150	WV	$\nu \Sigma_f$ for each group, or burnup matrix terms as mentioned above
1151 to 1270	DB2	DB^2 for each group
1271 to 1370	R(L,1)	microscopic reaction rates for capture 1 reaction for all nuclides ($r_0^{(1)}$ Section 7.3.1) in 10^{-24} /sec
1371 to 1470	R(L,2)	fission
1471 to 1570	R(L,3)	absorption
1571 to 1670	R(L,4)	ν -fission
1671 to 1770	R(L,5)	transport
1771 to 1870	R(L,6)	capture 2
1871 to 1970	DC	decay constants for each nuclide in 10^{-24} /sec
1971	ETC2 CRIT.R	radius of critical sphere in cm (extrapolation distance removed) based on B2(INF)
1972	EXTRAP.D	extrapolation distance in cm
1973 to 1999		spares

12.4 Additional Registers Z(0) - Z(860)

Z(0) ($\equiv X(2000)$) is used by subroutines as output argument and this variable is used to indicate which set of N's are to be used in forming macroscopic cross section information (5.3.2).

Z(1) to Z(500) are available for the user's calculations and are retained.

Z(501) to Z(860) are available for the user's calculations but are destroyed by the programme after a time card, T=0., 20,0, for example. The programme uses the array in terms of three arrays dimensioned WX(120), WY(120), and WZ(120).

12.5 Other Programme Data

It is possible to make available data beyond the user's COMMON by using the fact that the origin of blank COMMON is X(19491), and all programme variables in COMMON are listed relative to this origin in Section 21. Fixed point data may be changed using the origin J(19501).

13. SPECIAL FEATURES

13.0 Summary

Details are given of (1) a simple model used to represent ultimate performance of a recycling fuel system after reaching equilibrium, (2) a method of obtaining a soft spectrum essentially based on joining a Maxwellian-spectrum to a $1/E$ -spectrum, (3) a procedure for accumulating cross section data for later processing, and (4) a method used in the original source debugging.

13.1 Equilibrium Search

Following the ideas of Bicevskis (1963), a feature is available which allows the study of a homogeneous recirculating fuel reactor system operating under equilibrium conditions. The model implies infinite recirculation rate of fuel particles and does not fully account for resonance self screening by the fuel. Nevertheless the idea is a useful starting point for survey calculations biased towards calculating ultimate performance of the system after it has been operating for some time. For this model the composition of the core is the average of all states of burnup of a recirculating fuel particle prior to rejection. The required core power density is specified as P watts/cm³ and this is required to be the average power density of a fuel particle from first entry into the core to final rejection. Since the system is in equilibrium the model is equivalent to a fuel particle burning up under conditions of constant flux level and spectrum, which is a problem which makes full use of the analytic solution technique used in GYMEA (6.2). The problem is to find the equilibrium state core concentrations of component nuclides which (1) determine the spectrum and level through the specified power and (2) represent the average of all states of burnup of one fuel particle burning in this spectrum. An iterative procedure is available in GYMEA which

- (i) calculates flux level and spectrum from a guess at the equilibrium concentrations, $\hat{N}_A(\text{AFT.N})$;
- (ii) calculates the time of burnup of a fuel particle ($N_S - \text{BEF.N}$) necessary to achieve a specified number of fissions (F, FIFAS)— inner iteration;
- (iii) calculates the average composition of the burning fuel particle $\hat{N}_B(\text{AVE.N})$ and its average power density,
- (iv) if the power density is not that specified for the core, repeats the procedure from (i) using \hat{N}_B in place of \hat{N}_A — outer iteration.

JOB4 (15) illustrates the manner in which GYMEA is set up to do the problem. For this application it should be noted that the requirement for the inner iteration is specified with a SEARCH card (5.5.2) which must be punched with NO0 followed by the 3 eigenvalue parameters, the required FIFAS, and the required accuracy in FIFAS. The outer iteration is set to a required accuracy using a card, punched, for example,

ITERATE 9,0.1

which sets the outer iteration to terminate after 9 attempts if the power from one iteration to the next is not within the specified limit of 0.1 watt/cm³.

On completion of an equilibrium search, the concentrations AFT.N are normally overwritten with the concentrations AVE.N— step (iv) above. This may be prevented by setting

X(7)=-1.

13.2 Westcott Flux

When the subroutine WEST is used thus

CALL WEST

a group flux is calculated based on the Westcott convention with an r value (ratio of epithermal to thermal flux) specified using the variable $X(0)$ (5.5.3). In addition to the Maxwellian thermal region (temperature set by SCAT or TEMP card (5.2.4)), joining region, and 1/E region, a terminating region is added based on the fission source thus

$$\int_0^u S(u') du'$$

The flux per unit lethargy is then given by (Westcott 1957)

$$\phi(u) \propto (x^2 e^{-x} + \beta \Delta \int_0^u S(u') du')$$

where $x = E/kT$,

$$\beta = r/(1-1.01r),$$

and $\Delta = (1-0.26(1+(x/16.4)^5)^{-1} + (4.75/x)^7)^{-1}$.

This is integrated to give group flux following Cook (1964).

WEST is also used as a guess for group flux which is to be calculated from neutron balance of a homogeneous system (7.2) provided a group flux is not otherwise defined.

13.3 Saving Output on Tape 4

Normally, using the feature mentioned in Section 5.4.1 for saving output on tape 4 ($X(9) \leq -10.$), a file of cross section information is saved as a second file of binary records, maximum 256 words. This file consists of condensed and mixed cross section data given/reaction/group/material. If however $X(49)$ is N , then, not counting the first file, N files will be skipped before information is written or read. With this feature it is possible to use tape 4 as a library in another programme or to list the tape later using the option

CALL OUTPUT

In this application the additional information such as nuclide concentration, etc., is obtained from the data currently held in core and may therefore not relate to the cross section data.

13.4 Expanded Error Diagnostics

A card is available which sets the 'width' of programme COMMON (21) to be printed when an error is detected (normally zero). This feature was designed to assist programme debugging; however it may be of use to people trying to probe the manner in which GYMEA functions. An example of the use of the feature is

ERROR 1, 16000, 1,10000

which will result in a print out, when an error is detected, of COMMON 1 to 16000, excluding the array WH, and elements 1 to 10000 of the array WH.

14. CHECKOUT

14.0 Summary

Checkout of a code such as GYMEA must, of necessity, be divided into three categories: (1) theory checkout, (2) coding checkout, and (3) library checkout. Brief details are given below of some of the tests which have been carried out. These tests were all completely satisfactory.

14.1 Theory Checkout

- (1) Resonance theory checked against PEAS (Pollard 1964) and LUBRA (Kletzmayer 1966).
- (2) Analytic burnup theory checked with DEMON (Bennett 1965).

14.2 Coding Checkout

- (1) Group data preparation checked against MULGA (Clancy et al. 1963).
- (2) Burnup performance compared with FEVER (Todt 1962).
- (3) Neutron diffusion calculation of critical buckling for a bare system compared with CRAM (Hassitt 1962).

14.3 Library Checkout

- (1) Cross section data for dominant nuclides checked against MULGA data.
- (2) Selected cross section data plotted to the same scale as that given by Hughes and Schwartz(1958).

14.4 Maintenance

Maintenance based on these three main aspects will continue for the life of the GYMEA project, which is expected to be at least five years. From time to time further documents will be published to give the GYMEA user confidence in the ability of GYMEA to provide data which will adequately predict neutron behaviour in a reactor system. For example, correlation of A.A.E.C. measurements carried out on exponential assemblies composed of BeO - Pu - Th - U^{233} - U^{235} - U^{238} with predictions based on GYMEA data is at present being prepared.

At the time of preparation of this report the code GYMEA has been operational for half a year.


```

$IBSYS
$ID 30073
$JOB 2210 J.POLLARD 'GYMEA'
$* TAPES REQUIRED.
$PAUSE 295 ON 1, 260 ON 2 (VIA 1401), SCRATCH ON 3 AND 4
$TIME 60
$EXECUTE GYMEA
*GYMEA

```

```

,JOB1 A SIMPLE JOB AS ILLUSTRATION.
CALL STAN
USE LIB X-S=1 OF 100 NUCLIDES AND SCAT=2
'BE0'=1,'BE9',1,'O16',DENSITY=1.8,'
SCAT 900. 'BE0' CRY '
N(1,100)=0.,'BE0'=4.,-2,U233=4.,-5,'
OUTPUT 3 PRINTX PRINTS CRAM
MATS 'BE0',U233,PU239,'
GROUPS 5 1 40 50 71 90 120
T=0,DAYS,20 MAX FLUX ITERATIONS,0 CRITICALITY
END
CEASE

```

```

* *****
*

```

```

,DISCARD SUBROUTINES
SUBROUTINE PFLUX
PRINTS GROUP FLUX
NOTE THAT ONLY 3 CHARS OF NAME ARE UNIQUE,PFL..
FORMAT (6HOGFLUX/ (1X1P10E12.3))
WX 551,670,1
IE WRITE X(551) TO X(670) IN STEPS OF 1
RETURN
END

```

```

* *****
*

```

```

,JOB2 CRITICALITY
CREATION OF A LIB. OF BE9,O16, AND U235 TO SAVE TIME
CALL STAN
USE 1,70,-1 (A SCAT MATRIX LIB. IS NOT REQD)
CREATE 3 XDXD 100*-2 I.E. SAME STRUCTURE AS MAIN LIBRARY
BE9 AND O16 AS LIGHT NUCS AND J235 AS A RESONANCE NUC
MATS BE9,O16,U235,'

```

```

*
```

```

* ONLY ABOVE MATS ARE CREATED FOR NON-BURN UP USER
GFLUX=120*1.
FLUX=1.
NO16=1.
THE ABOVE ARE REQD TO MAKE A DEFINED PROBLEM
HEAD BE9,016 AND U235 FROM NDXD
T=0,1(GROUP FLUX IS NOT CALC. HOWEVER),0

*
CALL STAN
USE 3,3,2
=GYMEA XDXD NDSC
HEAD CF TABLE 3-98 ANL 5800(2ND ED.)
BE9'=1, BE9,0;'
O16'=1, O16,0;'
BE9' AND O16' ARE NOT REQD AS SUCH FOR THE MIXTURE BEING STUDIED
HENCE A DENSITY OF ZERO IS GIVEN
NOTE ALSO THAT THE 2 MATERIALS 'BE9' AND 'BEO' SHOULD NOT HAVE
THE SAME DOMINANT NUCLEIDE, BE9, IF BOTH ARE REQUIRED FOR VOLUME
NORMALIZATION USING CALL NJC
SCAT 300, BE9, GAS O16, GAS ''
GB2=120*8.-3
SEARCH GB2 0.01 0.8 100. 1.000 0.001
BUCKLING IS TO BE ADJUSTED TO MAKE K=1.000 + OR - 0.001
BEO'=1, BE9, 1.016, 1.+10 ''
DENSITY OF 'BEO' IS REQD TO BE EFFECTIVELY INFINITE
AS U5' IS GIVEN FOR THE MIXTURE
U5'=1, U235,-1;'
IE Z(1) CONTAINS DENSITY OF U5 IN MIXTURE (G/CM**3)
Z(1)=0.0140
N(1,100)=0., U5'=1., 'BEO'=1920.
CALL CALC
Z(1)=0.00702
N(1,100)=0., U5'=1., 'BEO'=3830.
CALL CALC
END

*****
*
SUBROUTINE CALC
FORCES CALCULATION OF CRITICALITY AND COMPUTES REQD QUANTITIES
CALL NUC
T=0.,20,6 MAX CRIT IT
RESCALES NUCLEIDE CONC TO UNIT VOL

```

```

Z(3)=X(1971)*X(1971)
Z(3)=Z(3)*X(1971)
Z(3)=4.1888*Z(3)
Z(4)=Z(3)*Z(1)
Z(4)=Z(4)*1.E-3
Z(5)=X(51)/X(53)
WRITE10  RATIO      MASS
WRITE1  BEO/U5      KG-U5
FORMAT(1H0,F7.0,F10.2)
WZS 2 5,4
* IE WRITE 2 Z'S, NAMELY Z(5),Z(4)
CALL PFLUX
PRINT GROUP FLUX
RETURN
END

```

```

*****
*

```

```

,JOB3  FIRST BURN UP RUN
CALL STAN
USE 1,61,2
NUCLIDES 62-70 ARE NOT REQD.
=GYMEA NDXD NDSC
'BE0'=1.016,1.016,1.8,
'PU'=0.76,PU239,0.17,PU240,0.07,PU241,0,
'PU02'=1.016,2.016,11,
'TH02'=1.016,2.016,9.5,
'FP'=0.023,61,0,
Z(0)=1
AFT.N WEIGHTED X-S REQD.
SCAT 900. 'BE0' CRY
HEAD FIRST BURN UP RUN
OUTPUT 1 PRINTX
MATS 'BE0','TH232','U233','PU239','PU240','PU241','FP'
GROUPS 2,1,71,120
X(6)=4.9
X(9)=4.9

```

```

*

```

```

OUTPUT IS SET TO PRODUCE MICRO X-S
N(1,100)=0.0, 'PU02'=1.0, 'TH02'=10.0, 'BE0'=2000.0,
CALL NUC
PART 0.01, 'PU02', 'TH02'
200 MICRON PARTICLES OF SEPARATED FUEL ('PU02', 'TH02')

```

```

*

```

CRIT VOL
MASS OF U5
KG-U5
NUCLIDE RATIO

```

* * *
FIFA PU239,PU241,,
THIS CARD OPENS THE WAY FOR A BURN UP USER AS ITS USE
DISTINGUISHES HIM FROM A USER ONLY INTERESTED IN T=0.
POWER=10. WATTS/CM**3
IE BURN UP AT STEP-WISE CONSTANT POWER
T=0.,20,0
DT=100.,20,0,2
2-GROUP OUTPUT IS NOT PRODUCED AFTER 100 DAYS
END
CEASE

```

```

* *****
*

```

```

,RETAIN SUBROUTINES
,JOB4 EQUILIBRIUM SEARCH
USE 1,61,2
=GYMEA NDXD NDSC
,BE0'=1.,BE9,1.,O16,0,,
,016'=-1.016,-2.016, 0 ,
,016'= 016 ASSOCIATED WITH HEAVY NUCLIDES
      = Z(1)+Z(2) (=1 LATER) PROPORTION OF 016 FOR CROSS SECTIONS
CONCENTRATION OF ,016' HOWEVER = CONCENTRATION OF 016/Z(1)
AND THIS WEIGHTS SCATTERING MATRIX OF ,016'
Z(1) AND Z(2) ARE DEFINED LATER
SCAT 900. ,BEO' CRY ,016' GAS ,
ITERATE 9,0.1 (WATTS/CM**3)
THE MAX. NO. OF (OUTER) ITERATIONS IS SET AND REQD POWER ACCURACY
OUTER REFERS TO SPECTRUM DET. FROM EQU. GUESS
SEARCH NOO .1 .9 10. 1.4 FIFA ACCURATE TO 0.005
INNER ITERATION-ADJUSTMENT OF TIME OF BURN UP
TO ACHIEVE A SPECIFIED FIFA
HEAD EQUILIBRIUM SEARCH
N(1,100)=0.,BE9=4.,-2.,O16=4.,046-2.,TH232=2.,-4.,PU239=2.,34-5,
PU240=5.,1-6.,PU241=1.,5-6.,
NOTE CONTINUATION CARD ABOVE(A NON-DESCENDING INTEGER IN COL 1)
Z(1)=X(52)-X(51)
Z(1)=X(52)/Z(1)
Z(2)=1.-Z(1)
INVERSE FRACTION OF 016
EXCESS 016

```

```

* * *
FIFA PU239,PU241,,
AFT.N ARRAY IS ALSO TRANSFERRED TO BEF.N
AND THIS MAT IS BURNT FOR THIS MODE OF OPERATION
N(3,100)=0.,TH232=1.,9-4.,U233=4.,3-6.,PU239=2.,8-6.,PU240=2.,-6,

```

```

* *

```

```
1 * PU241=2.2-6M  
* THIS SET IS AN EQU. GUESS AND DETS FIRST SPECTRUM  
* POWER=11.  
* CALL EQU  
* A FINAL SPECTRUM DET MUST BE FORCED IF MULTIGROUP OUTPUT  
* IS REQD--USE T=0.,20,0.,SAY  
CALL PFLUX  
UNLOAD 1  
UNLOAD 2  
TYPE TAPES UNLOADED  
END  
CEASE  
$IBSYS
```


16. COMMENTS ON OUTPUT

Heading information is widely used to clarify the meaning of output information (see Section 12.3 for the correspondence of printed labels with variables of the theory). The data entries given thus

```

      4      3      5      11
IT-FLUX,EIGEN.XS-0T

```

for the sample output (17) in order refer to

(1) the actual number of iterations taken by the subroutine calculating group flux from neutron balance (7.2) if its use has been specified (7.1) (minus the maximum number of iterations specified if convergence has not been achieved), or the number given on the time card, say $T=0,-1,0$, otherwise (5.5.3),

(2) the number of inner iterations associated with convergence of the SEARCH procedure (5.5.2) (negative if not converged),

(3) the number of outer iterations associated with the search for an equilibrium composition (13.1) (negative if not converged), and

(4) a number which indicates the overlap and transport averaging method adopted by the code (negative if the specified options cannot be accommodated) (5.4.2).

Columns 1 to 60 of the actual user's card which initiates the neutronics calculation are given to help label the output and suitable comments may be appended to this card if required.

The library nuclear data is under continual revision and hence the output given here may not be completely up to date. The source and type of data is shortly to be documented (Cook 1966).

VI TUPYUO EJMAZ

2210 J.POLLARD 'GYMEA'

\$ID 30073

\$JOB 2210 J.POLLARD 'GYMEA'

\$* TAPES REQUIRED.

\$PAUSE 295 ON 1, 260 ON 2 (VIA 1401), SCRATCH ON 3 AND 4

\$TIME 60

\$EXECUTE GYMEA

GYMEA DATA LIST * 23/12/65 *

ISN

SOURCE STATEMENT

```
1 ;JOB1 A SIMPLE JOB AS ILLUSTRATION.  
2 CALL STAN  
3 USE LIB X-S=1 OF 100 NUCLIDES AND SCAT=2  
4 ;BEO'=1.8E9,1.016,DENSITY=1.8''  
5 SCAT 900. ;BEO' CRY ''  
6 N(1,100)=0.;;BEO'=4.-2,U233=4.-5''  
7 OUTPUT 3 PRINTX PRINTS CRAM  
8 MATS ;BEO',U233,PU239''  
9 GROUPS 5 1 40 50 71 90 120  
10 T=0.DAYS,20 MAX FLUX ITERATIONS,0 CRITICALITY  
11 END  
12 CEASE
```

GYMEA JOB JOB1 X-SEC LIB. NDXD (120, 70, 0.K) ON 1 * 23/12/65 * OUTPUT NO. 1 RESONANCE TEMP= 900.K

(900.K) * 2 ----- ISN 10

UNIT SCE POWER

FLUX FLUX SCALING

FIFA

POWER-T

FLUX-T

T=0.DAYS,20 MAX FLUX ITERATIONS,0 CRITICALITY

(AFTER)
(AVERAGE)
(BEFORE)
(SPEC.)

A/S

1.0000101E+00 4 1 1 11
IT-FLUX,EIGEN,XS-OT

B2AV CRIT.-R-EXTRAP.D(CM)

B2 4.089E-03

2.2433

4.089E-03

4.089E-03 4.651E+01 2.623E+00

4.089E-03

2.2433

4.089E-03

SCAT BY 4.000E-02 'BED' CRY

AFT.N*SIGMA*FLUX/SOURCE NEUTRON

OUTPUT NO. 1 JOB1

T(RANGE)=	NUCLIDE	AFT.N	-DAYS- AVE.N	BEF.N	BURN	CAP1	FISS	ABS	NU-FISS	TRANS	CAP2	DECAY
0 LEAK	1 BE9	4.000E-02	4.000E-02	0.	0	3.344E-02	0.	3.041E-13	5.389E+01	8.280E-02	0.	1
2 O16	4.000E-02	4.000E-02	0.	0	0.	0.	8.119E-03	-3.813E-02	3.807E+01	0.	0.	2
13 U233	4.000E-05	4.000E-05	0.	0	0	1.330E-01	8.970E-01	1.030E+00	2.243E+00	9.735E-01	0.	13
18 PU239	4.000E-05	4.000E-05	0.	0	0	0.	0.	0.	0.	0.	0.	18

=====

MATERIAL GROUP	'BED' *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB1,* 23/12/65 *	NU-FISS	TRANS	CAP2	AS AFT.N*SIGMA*FLUX/SOURCE	SCATT XS-OT= 11	NEUTRON
1	0.	11.75	3.3442E-02	0.	-4.1151E-02	0.	0.	0.	4.9365E+01	0.	8.2799E-02	6.0057E+01	0.	0.
2	11.75	14.00	0.	0.	3.9708E-04	0.	0.	0.	1.1517E+01	0.	0.	1.2270E+01	0.	0.
3	14.00	16.10	0.	0.	9.1839E-04	0.	0.	0.	9.2726E+00	0.	0.	9.8793E+00	0.	0.
4	16.10	18.00	0.	0.	4.3350E-03	0.	0.	0.	1.4481E+01	0.	0.	1.5425E+01	0.	0.
5	18.00	23.00	0.	0.	5.4888E-03	0.	0.	0.	8.7726E+00	0.	0.	9.3411E+00	0.	0.

MATERIAL GROUP	U233 *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB1,* 23/12/65 *	NU-FISS	TRANS	CAP2	AS AFT.N*SIGMA*FLUX/SOURCE	SCATT XS-OT= 11	NEUTRON
1	0.	11.75	1.6346E-02	6.6815E-02	8.3162E-02	1.6784E-01	0.	0.	1.2224E-01	0.	0.	0.	0.	0.
2	11.75	14.00	2.0943E-02	7.3965E-02	9.4908E-02	1.8492E-01	0.	0.	1.0715E-01	0.	0.	0.	0.	0.
3	14.00	16.10	3.1014E-02	1.9827E-01	2.2928E-01	4.9566E-01	0.	0.	2.3532E-01	0.	0.	0.	0.	0.
4	16.10	18.00	3.4436E-02	2.7102E-01	3.0545E-01	6.7750E-01	0.	0.	3.2498E-01	0.	0.	0.	0.	0.
5	18.00	23.00	3.0236E-02	2.8698E-01	3.1717E-01	7.1738E-01	0.	0.	3.2568E-01	0.	0.	0.	0.	0.

MATERIAL GROUP	PU239 *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB1,* 23/12/65 *	NU-FISS	TRANS	CAP2	AS AFT.N*SIGMA*FLUX/SOURCE	SCATT XS-OT= 11	NEUTRON
1	0.	11.75	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
2	11.75	14.00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
3	14.00	16.10	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
4	16.10	18.00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
5	18.00	23.00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.

SUMMARY GROUP	INFORMATION	AFT.N	NORMALIZED ENERGY(EV)	OUTPUT NO	E(AVE,LETH)	FLUX/SOURCE	FISS.SPECT	TOT.LOSS	TOT.NU-FISS	TOT.DB#2
1	0.	11.75	1.0000E+07	7.8893E+01	2.8088E+04	1.7142E+02	1.00000	4.2011E-02	1.6784E-01	1.9794E-13
2	11.75	14.00	7.8893E+01	8.3153E+00	2.5613E+01	3.1366E+01	0.	9.5308E-02	1.8492E-01	2.8214E-14
3	14.00	16.10	8.3153E+00	1.0183E+00	2.9098E+00	2.5227E+01	0.	2.3024E-01	4.9570E-01	2.2312E-14
4	16.10	18.00	1.0183E+00	1.5230E-01	3.9380E-01	3.9562E+01	0.	3.0978E-01	6.7750E-01	3.5237E-14
5	18.00	23.00	1.5230E-01	1.0262E-03	1.2502E-02	2.3598E+01	0.	3.2266E-01	7.1738E-01	2.0402E-14
1 - 5						2.9118E+02	1.00000	1.0000E+00	2.2434E+00	3.0410E-13

AS AFT.N*SIGMA*FLUX/SOURCE NEUTRON

MATERIAL 'BEO' AFT.N= 4.0000E-02 OUTPUT NO 1, JOB1,* 23/12/65 *

OUTSCATTERS FROM GROUP 1

-9.5798E-01 9.5798E-01 0. 0. 0.

OUTSCATTERS FROM GROUP 2

0. -8.6402E-01 8.6402E-01 0. 0.

OUTSCATTERS FROM GROUP 3

0. 1.3525E-03 -7.1344E-01 7.1208E-01 6.0225E-07

OUTSCATTERS FROM GROUP 4

0. 7.9638E-02 -9.5901E-01 8.7937E-01

OUTSCATTERS FROM GROUP 5

0. 0. 5.5671E-01 -5.5671E-01

MATERIAL U233 AFT.N= 4.0000E-05 OUTPUT NO 1, JOB1,* 23/12/65 *

NO SCATTERING MATRIX

MATERIAL PU239 AFT.N= 0. OUTPUT NO 1, JOB1,* 23/12/65 *

NO SCATTERING MATRIX

		FND CF		JOB1	
TOTAL	JOB	TIME	TAKEN	IN	MINS
1.759E+00	1.656E+00	1.708E-01	9.672E-01	LINKS 3,4	LINKS 5
				4.297E-01	0.
BURNUP	FLUX	R-RATES	RUREDY	PAUSE	
0.	6.167E-01	2.10CE-01	2.194E-02	0.	

NEW LOAD CALLED

... CARD OUTPUT ...

```

=GYMEA JOB JOB1 CRAM OUTPUT NO. 1 * 23/12/65 *
=GYMEA X-S NDXD * 24/11/65 * MAIN GENERAL PURPOSE LIB. OF 70 NUC.
=GYMEA SCAT NDSC 3/12/65 BE01BE910161H201C121C123009006001200
SP JOB1 1
I 1 1.00000E+00 0. 0. 0. 0. 0.
I 1 7.19930E+00 9.17908E+00 9.18896E+00 9.15082E+00 9.29381E+00
1.33708E-01 6.88975E-01 7.07914E-01 6.08765E-01 5.95600E-01
0. 0. 0. 0. 0.
0. 1.39710E-01 0. 0. 0. 0.
0. 6.88659E-01. 0. 0. 0.
0. 1.34033E-03 0. 7.05663E-01 5.96820E-07
0. 0. 5.03257E-02 0. 5.57700E-01
0. 0. 0. 5.89785E-01 0.
I 2 1.78271E+01 8.54057E+01 2.33202E+02 2.05363E+02 3.45034E+02
1.21281E+01 7.56455E+01 2.27215E+02 1.93021E+02 3.36017E+02
2.44780E+01 1.47389E+02 4.91194E+02 4.28131E+02 7.60000E+02
0. 0. 0. 0. 0.
0. 0. 0. 0. 0.
0. 0. 0. 0. 0.
0. 0. 0. 0. 0.
I 3 1.41353E+01 9.44861E+01 4.24192E+01 1.65160E+03 8.15896E+02
7.15697E+00 8.32656E+01 3.06373E+01 1.63685E+03 8.09863E+02
1.34275E+01 1.33457E+02 5.37102E+01 2.87482E+03 1.61275E+03
0. 0. 0. 0. 0.
0. 0. 0. 0. 0.
0. 0. 0. 0. 0.
0. 0. 0. 0. 0.
***** END OF CRAM DATA. FOR JOB1 OUTPUT NO 1

```

SOURCE STATEMENT

ISN

```

1 SUBROUTINE PFLUX
2 PRINTS GRGUP FLUX
3 *
4 * NOTE THAT ONLY 3 CHARS OF NAME ARE UNIQUE,PFL..
5 *
6 * FORMAT (6HOGFLUX/ (1XIP10E12.3))
7 *
8 * WX 551,670,1
9 *
10 * IE WRITE X(551) TO X(670) IN STEPS OF 1
11 *
12 * RETURN
13 *
14 * END
15 *
16 *****
17 *
18 * ,JOB2 CRITICALITY
19 * CREATION OF A LIB. OF BE9,016, AND U235 TO SAVE TIME
20 * CALL STAN
21 *
22 * USE 1,70,-1 (A SCAT MATRIX LIB. IS NOT REQD)
23 * CREATE 3 XDXD 100*-2 I.E. SAME STRUCTURE AS MAIN LIBRARY
24 * BE9 AND 016 AS LIGHT NUCLS AND U235 AS A RESONANCE NUC
25 * MATS BE9,016,U235,,
26 * ONLY ABOVE MATS ARE CREATED FOR NON-BURN UP USER
27 * GFLUX=120*1.
28 * FLUX=1.
29 *
30 * NO16=1.,,
31 * THE ABOVE ARE REQD TO MAKE A DEFINED PROBLEM
32 * HEAD BE9,016 AND U235 FROM NDXD
33 * T=0,1(GROUP FLUX IS NOT CALC. HOWEVER),0
34 *
35 * CALL STAN
36 * USE 3,3,2
37 * =GYMEA XDXD NDSC
38 * HEAD CF TABLE 3-98 ANL 5800(2ND ED.)
39 *
40 * 'BE9'=1.,BE9,0,,
41 * '016'=1.,016,0,,
42 * 'BE9' AND '016' ARE NOT REQD AS SUCH FOR THE MIXTURE BEING STUDIED
43 * HENCE A DENSITY OF ZERO IS GIVEN
44 * NOTE ALSO THAT THE 2 MATERIALS 'BE9' AND 'BE0' SHOULD NOT HAVE
45 * THE SAME DOMINANT NUCLIDE,BE9, IF BOTH ARE REQUIRED FOR VOLUME
46 * NORMALIZATION USING CALL NUC
47 * SCAT 300. 'BE9' GAS '016' GAS **
48 * G82=120*8.-3
49 * SEARCH G82 0.01 0.8 100. 1.000 0.001
50 * BUCKLING IS TO BE ADJUSTED TO MAKE K=1.000 + OR - 0.001
51 * 'BE0'=1.,BE9,1.,016,1.,+10 **
52 * DENSITY OF 'BE0' IS REQD TO BE EFFECTIVELY INFINITE
53 * AS 'U5' IS GIVEN FOR THE MIXTURE
54 * 'U5'=1.,U235,-1,,
55 * IE Z(1) CONTAINS DENSITY OF U5 IN MIXTURE (G/CM**3)
56 * Z(1)=0.,0140
57 * N(1,100)=0., 'U5'=1., 'BE0'=1920.,,
58 * CALL CALC
59 * Z(1)=0.00702
60 * N(1,100)=0., 'U5'=1., 'BE0'=3830.,,
61 * CALL CALC
62 * END

```

```

34 *****
34 *
34 * SUBROUTINE CALC
34 * FORCES CALCULATION OF CRITICALITY AND COMPUTES REQD QUANTITIES
35 * CALL NUC
35 * T=0.,20,6 MAX CRIT IT
36 * Z(3)=X(1971)*X(1971)
37 * Z(3)=Z(3)*X(1971)
38 * Z(3)=4.1888*Z(3)
39 * Z(4)=Z(3)*Z(1)
40 * Z(4)=Z(4)*1.E-3
41 * Z(5)=X(51)/X(53)
42 * WRITE(10,RATIO) MASS
43 * WRITE(10,BE0/U5 KG-U5
44 * FORMAT(1H0,F7.0,F10.2)
45 * WZS 2 5,4
46 * IE WRITE 2 Z'S, NAMELY Z(5),Z(4)
47 * CALL PELUX
48 * PRINT GROUP FLUX
49 * RETURN
50 * END
50 *****
50 *
50 * ,JOB3 FIRST BURN UP RUN
51 * CALL STAN
51 * USE 1,61,2
52 * NUCLIDES 62-70 ARE NOT REQD.
53 * =CYMEA NDXD NOSC
54 * 'BE0'=1.8E9,1.016,1.8,,
55 * 'PU'=0.76*PU239,0.17*PU240,0.07*PU241,0,,
56 * 'PU02'=1.,'PU',2.016,11.,
57 * 'TH02'=1.,'TH232',2.016,9.5,,
58 * 'FP'=0.(23,61),0,,
59 * Z(10)=1
60 * AFT.N WEIGHTED X-S REQD.
61 * SCAT 900. 'BE0' CRY ,,
62 * HEAD FIRST BURN UP RUN
63 * OUTPUT 1 PRINTX
64 * MATS 'BE0','TH232','U233','PU239','PU240','PU241','FP' ,,
65 * GROUPS 2,1,71,120
66 * X(6)=4.9
67 * X(9)=4.9
67 * OUTPUT IS SET TO PRODUCE MICRO X-S
67 * N(1,100)=0.,'PU02'=1.,'TH02'=10.,'BE0'=2000.,'
68 * CALL NUC
69 * PART 0.01,'PU02','TH02' ,,
70 * 200 MICRON PARTICLES OF SEPARATED FUEL ('PU02','TH02')
70 * FIFA PU239,PU241,,
71 * THIS CARD OPENS THE WAY FOR A BURN UP USER AS ITS USE
71 * DISTINGUISHES HIM FROM A USER ONLY INTERESTED IN T=0.
71 * POWER=10. WATTS/CM**3
72 * IE BURN UP AT STEP-WISE CONSTANT POWER
72 * T=0.,20,0
73 * DT=100.,20,0,2
74 * 2-GROUP OUTPUT IS NOT PRODUCED AFTER 100 DAYS
74 * END
75 * CEASE

```

GYMEA JOB JOB2 X-SEC LIB. NDXD (120, 70, 0.K) ON 1 * 23/12/65 * ISN 15 OUTPUT NO. 1 RESONANCE TEMP= 0.K
 SCATT : 00000 (0.K) *-1
 T(DAYS) FLUX-T POWER-T FIF A POWER CONSTANT (AFTER)
 0. 0. 0. 0. 1.000E+00 0. (AVERAGE)
 (AFT.N) K 82 82AV CRIT.R-EXTRAP.D(CM) A/S (BEFORE)
 (EFF.) 0. -2.253E-03 (INF.) 0. -2.253E-03-2.253E-03 0. 2.069E-01 1.0000001E+00 (SPEC.)
 II-FLUX,EIGEN.XS-OT 1 1 1 11

SCAT BY 0. N00 AT ALL
 OUTPUT NO. 1 JOB2 BF9,016 AND U235 FROM NDXD AFT.N*SIGMA*FLUX/SOURCE NEUTRON
 T(RANGE)= 0. -DAYS-
 NUCLIDE AFT.N BEF.N BURN CAP1 FISS ABS NU-FISS TRANS CAP2 DECAY
 0 LEAK 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 4.438E-13 0.000E+00 0.000E+00 0.000E+00
 1 BE9 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00
 2 U16 1.000E+00 1.000E+00 0.000E+00 0.000E+00 0.000E+00 1.000E+00 1.704E+04 0.000E+00 0.000E+00
 15 U235 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00 0.000E+00
 =====

-----GYMEA-----* 23/12/65 *-----GYMEA-----

ISN 6 +JOB2 CRITICALITY

LIB. 3 =GYMEA X-S XCDX * 23/12/65 * 8E9,016 AND U235 FROM NDXD

LIB. 2 =GYMEA SCAT NCSC 3/12/65 'BEU'BE9''016''H20''C12''300,900,600,1200

GYMEA JOB JOB2 X-SEC LIB. XDXD (120, 3, 0.K) ON 3 * 23/12/65 * OUTPUT NO. 1 RESONANCE TEMP= 300.K
 SCATT NDSC (300.K) * 2 ISN 36
 T(DAYS) FLUX-T POWER-T FIF A UNIT SCE FLUX POWER (AFTER)
 0. 0. 0. 0. 1.326E+02 0. (AVERAGE)
 (AFT.N) K B2 B2AV CRIT.R-EXTRAP.D(CM) A/S (BEFORE)
 (EFF.) 1.0007 8.459E-06 82AV CRIT.R-EXTRAP.D(CM) A/S (SPEC.)
 (INF.) 2.0712 6.341E-03 6.341E-03 3.789E+01 1.559E+00 9.9999224E-01 3 3 1 11
 IT-FLUX,EIGEN.XS-OT

SCAT BY 6.889E-02 'BE9' GAS 6.889E-02 'D16' GAS
 OUTPUT NO. 1 JOB2 CF TABLE 3-98 ANL 5800(2ND ED.) AFT.N*SIGMA*FLUX/SOURCE NEUTRON
 T(RANGE)= 0. -DAYS- 0.
 NUCLIDE AFT.N 86F.N BURN CAP1 FISS ABS NU-FISS TRANS CAP2 DECAY
 O LEAK 0.0000 5.168E-01

RATIO MASS 1920. 3.19
 BEO/U5 KG-U5

GFLUX

5.631E+00	4.501E+00	2.107E+00	2.565E+00	2.237E+00	1.679E+00	2.183E+00	2.180E+00	2.079E+00	2.087E+00
2.083E+00	1.888E+00	1.869E+00	1.861E+00	1.843E+00	1.828E+00	1.816E+00	1.802E+00	1.789E+00	1.775E+00
1.761E+00	1.747E+00	1.733E+00	1.720E+00	1.706E+00	1.693E+00	1.679E+00	1.666E+00	1.652E+00	1.638E+00
1.624E+00	1.610E+00	1.596E+00	1.582E+00	1.568E+00	1.552E+00	1.537E+00	1.522E+00	1.508E+00	1.493E+00
1.474E+00	1.458E+00	1.444E+00	1.430E+00	1.414E+00	1.396E+00	1.383E+00	1.353E+00	1.332E+00	1.319E+00
5.198E-01	5.168E-01	5.119E-01	5.128E-01	5.120E-01	5.097E-01	5.080E-01	5.039E-01	5.036E-01	5.036E-01
5.032E-01	5.020E-01	5.035E-01	5.001E-01	5.023E-01	4.990E-01	5.011E-01	4.984E-01	5.010E-01	4.944E-01
4.974E-01	4.964E-01	4.987E-01	5.000E-01	5.003E-01	5.029E-01	5.048E-01	5.107E-01	5.107E-01	5.127E-01
5.172E-01	5.226E-01	5.267E-01	5.358E-01	5.474E-01	5.657E-01	5.903E-01	6.258E-01	6.679E-01	7.196E-01
7.813E-01	8.462E-01	9.178E-01	9.772E-01	1.040E+00	1.085E+00	1.117E+00	1.132E+00	1.129E+00	1.108E+00
1.072E+00	1.024E+00	9.616E-01	8.948E-01	8.234E-01	7.463E-01	6.734E-01	6.003E-01	5.293E-01	4.663E-01
4.058E-01	3.510E-01	3.040E-01	2.594E-01	2.227E-01	6.912E-01	2.868E-01	1.130E-01	4.313E-02	1.622E-02

-----GYMEA-----* 23/12/65 *-----GYMEA-----

ISN 50 , JOB3 FIRST BURN UP RUN

LIB. 1 =GYMEA X-S NDXD * 24/11/65 * MAIN GENERAL PURPOSE LIB. OF 70 NUC.

LIB. 2 =GYMEA SCAT NDSC 3/12/65 'BEO'BE9'D16'H20'C12'C12'300,900,600,1200

T (DAYS) 0.
 FLUX-T 0. POWER-T 0. FIF A 0. FLUX 1.778E+14
 CONSTANT
 POWER 1.000E+01 (AFTER)
 1.000E+01 (AVERAGE)
 1.000E+01 (BEFORE)
 1.000E+01 (SPEC.)
 A/S
 CRIT.-R-EXTRAP.D(CM) 9.998929E-01
 K B2 82AV
 8.833E-04
 8.833E-04 8.833E-04 1.032E+02 2.486E+00
 IT-FLUX,EIGEN.XS-OT 4 1 1 11

SCAT BY 4.290E-02 'BEO' CRY

OUTPUT NO. 1 JOB3 FIRST BURN UP RUN

T (RANGE) =	NUCLIDE	AFT.N	0.	-DAYS- AVE.N	BEF.N	BURN	CAP1	FISS	ABS	NU-FISS	TRANS	CAP2	DECAY
0 LEAK													
1 BE9	4.290E-02	4.290E-02	4.290E-02	4.290E-02	4.290E-02	0	3.219E-03	0.	-4.158E-03	0.	4.533E+00	7.969E-03	0.
2 O16	4.337E-02	4.337E-02	4.337E-02	4.337E-02	4.337E-02	0	0.	0.	7.721E-04	0.	3.231E+00	0.	0.
3 C12	0.	0.	0.	0.	0.	0	1.056E-03	0.	1.056E-03	0.	3.723E+00	0.	0.
4 H1	0.	0.	0.	0.	0.	0	2.039E-02	0.	2.039E-02	0.	5.110E+00	0.	0.
5 O2	0.	0.	0.	0.	0.	0	3.500E-05	0.	-1.260E-03	0.	2.178E+00	1.295E-03	0.
6 AL27	0.	0.	0.	0.	0.	0	1.746E-02	0.	1.782E-02	0.	2.214E+00	0.	0.
7 LI6	0.	0.	0.	0.	0.	0	5.694E+01	0.	5.694E+01	0.	4.719E+01	0.	0.
8 T3	0.	0.	0.	0.	0.	0	0.	0.	0.	0.	1.000E+00	0.	1.009E+01
9 HE3	0.	0.	0.	0.	0.	0	3.353E+02	0.	3.353E+02	0.	2.749E+02	0.	0.
10 TH232	2.145E-04	2.145E-04	2.145E-04	2.145E-04	2.145E-04	0	3.276E+00	7.615E-03	3.282E+00	1.736E-02	1.226E+01	1.017E-03	0.
11 PA233	0.	0.	0.	0.	0.	0	5.020E+01	6.370E-02	5.027E+01	1.539E-01	5.221E+01	0.	1.647E+03
12 U232	0.	0.	0.	0.	0.	0	1.842E+01	0.	2.333E+01	0.	1.913E+01	0.	0.
13 U233	0.	0.	0.	0.	0.	0	9.555E+00	6.009E+01	6.965E+01	1.503E+02	6.717E+01	0.	0.
14 U234	0.	0.	0.	0.	0.	0	3.603E+01	1.754E-01	3.620E+01	3.950E-01	3.926E+01	0.	0.
15 U235	0.	0.	0.	0.	0.	0	1.091E+01	3.196E+01	4.287E+01	7.802E+01	4.564E+01	0.	0.
16 U236	0.	0.	0.	0.	0.	0	1.637E+01	6.370E-02	1.644E+01	1.539E-01	2.215E+01	0.	0.
17 U238	0.	0.	0.	0.	0.	0	1.358E+01	3.203E-02	1.361E+01	8.518E-02	2.076E+01	0.	0.
18 PU239	1.630E-05	1.630E-05	1.630E-05	1.630E-05	1.630E-05	0	5.840E+01	9.600E+01	1.544E+02	2.784E+02	1.308E+02	0.	0.
19 PU240	3.646E-06	3.646E-06	3.646E-06	3.646E-06	3.646E-06	0	2.321E+02	3.209E-02	2.322E+02	9.408E-02	1.354E+02	0.	0.
20 PU241	1.501E-06	1.501E-06	1.501E-06	1.501E-06	1.501E-06	0	3.738E+01	9.758E+01	1.349E+02	2.986E+02	1.182E+02	0.	9.506E+00
21 PU242	0.	0.	0.	0.	0.	0	2.633E+01	0.	2.633E+01	0.	3.342E+01	0.	0.
22 Z2999	0.	0.	0.	0.	0.	0	6.141E-02	0.	6.141E-02	0.	5.034E-02	0.	0.
23 MO95	0.	0.	0.	0.	0.	0	6.366E+00	0.	6.366E+00	0.	1.382E+01	0.	0.
24 TC99	0.	0.	0.	0.	0.	0	1.081E+01	0.	1.081E+01	0.	1.523E+01	0.	0.
25 RU101	0.	0.	0.	0.	0.	0	1.060E+01	0.	1.060E+01	0.	1.716E+01	0.	0.
26 RU102	0.	0.	0.	0.	0.	0	8.995E+00	0.	8.995E+00	0.	1.560E+01	0.	0.
27 RH103	0.	0.	0.	0.	0.	0	6.818E+01	0.	6.818E+01	0.	5.514E+01	0.	0.
28 RU104	0.	0.	0.	0.	0.	0	1.359E+01	0.	1.359E+01	0.	1.961E+01	0.	0.
29 RH105	0.	0.	0.	0.	0.	0	3.408E+03	0.	3.408E+03	0.	1.000E+01	0.	3.008E+04
30 PD105	0.	0.	0.	0.	0.	0	4.379E+00	0.	4.379E+00	0.	9.014E+00	0.	0.
31 PD106	0.	0.	0.	0.	0.	0	1.828E+00	0.	1.828E+00	0.	6.241E+00	0.	0.
32 PD107	0.	0.	0.	0.	0.	0	4.001E+00	0.	4.001E+00	0.	8.743E+00	0.	0.
33 PD108	0.	0.	0.	0.	0.	0	1.025E+01	0.	1.025E+01	0.	1.495E+01	0.	0.
34 AG109	0.	0.	0.	0.	0.	0	7.189E+01	0.	7.189E+01	0.	7.000E+01	0.	0.
35 CD113	0.	0.	0.	0.	0.	0	2.384E+03	0.	2.384E+03	0.	1.940E+03	0.	0.
36 I131	0.	0.	0.	0.	0.	0	3.507E+01	0.	3.507E+01	0.	3.806E+01	0.	5.606E+03
37 XE131	0.	0.	0.	0.	0.	0	4.428E+01	0.	4.428E+01	0.	4.913E+01	0.	0.
38 XE133	0.	0.	0.	0.	0.	0	9.168E+01	0.	9.168E+01	0.	9.298E+01	0.	8.561E+03
39 CS133	0.	0.	0.	0.	0.	0	1.979E+01	0.	1.979E+01	0.	2.623E+01	0.	0.

SIGMA

MATERIAL GROUP	'BED' *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	SCATT	AS	SIGMA	XN-OT=
1	0.	16.10	23.00	3.6940E-03	0.	0.	-4.4001E-03	0.	7.5955E+00	9.1459E-03	9.0074E+00	9.7957E+00	0.	0.	0.	0.	11
2	16.10	23.00	0.	0.	0.	0.	3.4780E-03	0.	9.1978E+00	0.	0.	0.	0.	0.	0.	0.	11

MATERIAL GROUP	TH232 *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	SCATT	AS	SIGMA	XN-OT=
1	0.	16.10	23.00	3.4327E+00	8.7396E-03	0.	3.4403E+00	1.9921E-02	1.2114E+01	1.1668E-03	0.	0.	0.	0.	0.	0.	11
2	16.10	23.00	0.	2.2121E+00	0.	0.	2.2121E+00	0.	1.3464E+01	0.	0.	0.	0.	0.	0.	0.	11

MATERIAL GROUP	U233 *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	SCATT	AS	SIGMA	XN-OT=
1	0.	16.10	23.00	7.5525E+00	3.9002E+01	0.	4.6555E+01	9.7595E+01	4.7352E+01	0.	0.	0.	0.	0.	0.	0.	11
2	16.10	23.00	0.	2.3115E+01	2.0288E+02	0.	2.2597E+02	5.0715E+02	2.3654E+02	0.	0.	0.	0.	0.	0.	0.	11

MATERIAL GROUP	PU239 *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	SCATT	AS	SIGMA	XN-OT=
1	0.	16.10	23.00	7.8998E+00	1.1226E+01	0.	1.9125E+01	3.2630E+01	2.4824E+01	0.	0.	0.	0.	0.	0.	0.	11
2	16.10	23.00	0.	4.0029E+02	6.6991E+02	0.	1.0702E+03	1.9426E+03	1.0363E+03	0.	0.	0.	0.	0.	0.	0.	11

MATERIAL GROUP	PU240 *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	SCATT	AS	SIGMA	XN-OT=
1	0.	16.10	23.00	2.5009E+02	3.6828E-02	0.	2.5012E+02	1.0797E-01	1.3703E+02	0.	0.	0.	0.	0.	0.	0.	11
2	16.10	23.00	0.	1.1112E+02	0.	0.	1.1112E+02	0.	1.2138E+02	0.	0.	0.	0.	0.	0.	0.	11

MATERIAL GROUP	PU241 *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	SCATT	AS	SIGMA	XN-OT=
1	0.	16.10	23.00	1.0870E+01	3.3247E+01	0.	4.4117E+01	1.0177E+02	4.5534E+01	0.	0.	0.	0.	0.	0.	0.	11
2	16.10	23.00	0.	2.1688E+02	5.3313E+02	0.	7.4990E+02	1.6312E+03	7.3875E+02	0.	0.	0.	0.	0.	0.	0.	11

MATERIAL GROUP	'FP' *****	AFT.N=	LETH	CAPI	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	SCATT	AS	SIGMA	XN-OT=
1	0.	16.10	23.00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	11
2	16.10	23.00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.	11

SUMMARY GROUP	INFORMATION, AFT.N	LETH	NORMALIZED ENERGY(EV)	OUTPUT NO	FISS	ABS	1, JOB3 ,*	23/12/65 *	NU-FISS	FIRST BURN UP	TRANS	CAP2	TOT.LOSS	TOT.NU-FISS	TOT.DB**2
1	0.	16.10	1.0000E+07	1.0183E+00	1.0183E+00	3.1910E+03	2.1121E+02	1.00000	3.8943E-01	1.4637E-01	2.1276E-13	0.	0.	0.	0.
2	16.10	23.00	1.0183E+00	1.0262E-03	3.2325E-02	3.1197E+01	0.	0.	6.1146E-01	1.0643E+00	2.4900E-14	0.	0.	0.	0.
1 - 2									2.4241E+02	1.00000	0.	0.	1.0009E+00	1.2107E+00	2.3766E-13

ISN 73

DT=100.,20,0,2

CONSTANT
 FLUX POWER
 1.536E+21 1.006E+03 0.1487 1.778E+14 9.447E+00 (AFTER)
 1.006E+01 (AVERAGE)
 1.000E+01 (BEFORE)
 1.000E+01 (SPEC.)
 A/S

(AFT.N) K B2 B2AV CRIT.R-EXTRAP.D(CM) 9.9999911E-01 4 1 1 11
 (EFF.) 1.1511 6.324E-04
 (INF.) 1.1511 6.324E-04 7.835E-04 1.224E+02 2.482E+00
 II-FLUX,EIGEN,XS-UT

SCAT BY 4.290E-02 *8E0* CRY

OUTPUT NO. 2 JOB3 FIRST BURN UP RUN

SIGMA

T(RANGE)=	NUCLIDE	AFT.N	-DAYS- AVE.N	BEF.N	BURN	CAP1	FISS	ABS	NU-FISS	TRANS	CAP2	DECAY
0 LEAK												
1 BE9	4.290E-02	4.290E-02	4.290E-02	4.290E-02	0	3.199E-03	0.	-4.104E-03	0.	4.538E+00	7.921E-03	0.
2 O16	4.337E-02	4.337E-02	4.337E-02	4.337E-02	0	0.	0.	7.681E-04	0.	3.233E+00	0.	0.
3 C12	0.	0.	0.	0.	0	1.060E-03	0.	1.060E-03	0.	3.727E+00	0.	0.
4 H1	0.	0.	0.	0.	0	2.128E-02	0.	2.128E-02	0.	5.141E+00	0.	0.
5 D2	0.	0.	0.	0.	0	3.654E-05	0.	-1.251E-03	0.	2.180E+00	1.287E-03	0.
6 AL27	0.	0.	0.	0.	0	1.808E-02	0.	1.844E-02	0.	2.211E+00	0.	0.
7 LI6	2.031E-07	1.030E-07	0.	0.	0	5.943E+01	0.	5.943E+01	0.	4.931E+01	0.	0.
8 T3	8.808E-09	3.252E-09	0.	0.	0	0.	0.	0.	0.	1.000E+00	0.	1.009E+01
9 ME3	4.002E-11	7.221E-12	0.	0.	0	3.500E+02	0.	3.500E+02	0.	2.874E+02	0.	0.
10 TH232	2.134E-04	2.140E-04	2.145E-04	2.145E-04	0	3.279E+00	7.570E-03	3.285E+00	1.725E-02	1.226E+01	1.011E-03	0.
11 PA233	3.821E-07	2.664E-07	0.	0.	0	5.055E+01	6.332E-02	5.062E+01	1.529E-01	5.261E+01	0.	1.647E+03
12 U232	3.282E-10	1.652E-10	0.	0.	0	1.923E+01	0.	2.436E+01	0.	2.000E+01	0.	0.
13 U233	6.464E-07	2.566E-07	0.	0.	0	9.668E+00	6.113E+01	7.079E+01	1.529E+02	6.820E+01	0.	0.
14 U234	2.381E-08	8.968E-09	0.	0.	0	3.575E+01	1.743E-01	3.593E+01	3.926E-01	3.904E+01	0.	0.
15 U235	4.879E-10	1.314E-10	0.	0.	0	1.118E+01	3.348E+01	4.466E+01	8.173E+01	4.717E+01	0.	0.
16 U236	2.046E-12	1.251E-12	0.	0.	0	1.618E+01	6.332E-02	1.625E+01	1.529E-01	2.201E+01	0.	0.
17 U238	0.	0.	0.	0.	0	1.349E+01	3.184E-02	1.352E+01	8.468E-02	2.069E+01	0.	0.
18 PU239	1.286E-05	1.451E-05	1.630E-05	1.630E-05	0	6.312E+01	1.036E+02	1.667E+02	3.005E+02	1.419E+02	0.	0.
19 PU240	3.640E-06	3.668E-06	3.646E-06	3.646E-06	0	2.275E+02	3.190E-02	2.276E+02	9.352E-02	1.334E+02	0.	0.
20 PU241	2.376E-06	1.953E-06	1.501E-06	1.501E-06	0	3.976E+01	1.033E+02	1.430E+02	3.160E+02	1.252E+02	0.	9.508E+00
21 PU242	1.102E-07	5.128E-08	0.	0.	0	2.586E+01	0.	2.586E+01	0.	3.301E+01	0.	0.
22 Z7999	0.	0.	0.	0.	0	6.411E-02	0.	6.411E-02	0.	5.264E-02	0.	0.
23 MO95	1.196E-07	6.123E-08	0.	0.	0	6.374E+00	0.	6.374E+00	0.	1.382E+01	0.	0.
24 TC99	1.486E-07	7.616E-08	0.	0.	0	1.074E+01	0.	1.074E+01	0.	1.517E+01	0.	0.
25 RU101	1.445E-07	7.408E-08	0.	0.	0	1.054E+01	0.	1.054E+01	0.	1.712E+01	0.	0.
26 RU102	1.454E-07	7.438E-08	0.	0.	0	8.936E+00	0.	8.936E+00	0.	1.555E+01	0.	0.
27 RH103	1.322E-07	6.870E-08	0.	0.	0	6.708E+01	0.	6.708E+01	0.	5.373E+01	0.	0.
28 RU104	1.425E-07	7.330E-08	0.	0.	0	1.351E+01	0.	1.351E+01	0.	1.953E+01	0.	0.
29 RH105	2.498E-09	2.602E-09	0.	0.	0	3.400E+03	0.	3.400E+03	0.	1.000E+01	0.	3.008E+04
30 PD105	1.198E-07	6.003E-08	0.	0.	0	4.386E+00	0.	4.386E+00	0.	9.014E+00	0.	0.
31 PD106	1.263E-07	6.441E-08	0.	0.	0	1.835E+00	0.	1.835E+00	0.	6.243E+00	0.	0.
32 PD107	9.160E-08	4.679E-08	0.	0.	0	4.006E+00	0.	4.006E+00	0.	8.744E+00	0.	0.
33 PD108	6.712E-08	3.426E-08	0.	0.	0	1.021E+01	0.	1.021E+01	0.	1.491E+01	0.	0.
34 AG109	4.050E-08	2.085E-08	0.	0.	0	7.111E+01	0.	7.111E+01	0.	6.929E+01	0.	0.
35 CD113	7.237E-10	5.327E-10	0.	0.	0	2.592E+03	0.	2.592E+03	0.	2.117E+03	0.	0.
36 I131	1.010E-08	9.466E-09	0.	0.	0	3.480E+01	0.	3.480E+01	0.	3.783E+01	0.	5.606E+03
37 XE131	7.901E-08	3.664E-08	0.	0.	0	4.434E+01	0.	4.434E+01	0.	4.923E+01	0.	0.
38 XE133	1.182E-08	1.164E-08	0.	0.	0	9.091E+01	0.	9.091E+01	0.	9.231E+01	0.	8.561E+03
39 CS133	1.509E-07	7.214E-08	0.	0.	0	1.961E+01	0.	1.961E+01	0.	2.607E+01	0.	0.

40	CS134	2.128E-09	6.783E-10	0.	6.179E+01	0.	6.508E+01	0.	0.	40
41	II135	6.477E-10	6.904E-10	0.	1.503E-02	0.	1.000E+01	0.	1.617E+05	41
42	XE135	4.483E-10	4.765E-10	0.	1.346E+05	0.	1.429E+05	0.	1.177E+05	42
43	CS135	8.602E-08	4.376E-08	0.	3.177E+00	0.	1.533E+01	0.	0.	43
44	ND143	1.149E-07	5.869E-08	0.	1.654E+01	0.	2.982E+01	0.	0.	44
45	ND144	9.788E-08	4.954E-08	0.	6.698E-01	0.	1.000E+01	0.	0.	45
46	ND145	7.649E-08	3.917E-08	0.	2.169E+01	0.	4.616E+01	0.	0.	46
47	ND146	6.570E-08	3.330E-08	0.	1.697E+00	0.	1.000E+01	0.	0.	47
48	PM147	4.647E-08	2.444E-08	0.	6.152E+01	0.	1.337E+02	5.524E+01	0.	48
49	PM148	2.772E-10	1.363E-10	0.	1.399E+03	0.	2.057E+03	0.	8.370E+03	49
50	PM8M	3.485E-10	1.667E-10	0.	5.994E+03	0.	5.146E+03	0.	1.111E+03	50
51	SM149	6.547E-09	5.340E-09	0.	3.542E+03	0.	2.932E+03	0.	0.	51
52	SM150	2.634E-08	1.133E-08	0.	4.659E+01	0.	7.487E+01	0.	0.	52
53	SM151	1.533E-08	8.390E-09	0.	3.376E+02	0.	3.266E+02	0.	0.	53
54	SM152	1.650E-08	8.074E-09	0.	1.152E+02	0.	1.286E+02	0.	0.	54
55	EU153	8.676E-09	4.309E-09	0.	9.567E+01	0.	1.050E+02	0.	0.	55
56	EU154	4.985E-10	1.720E-10	0.	4.883E+02	0.	4.233E+02	0.	7.723E+00	56
57	EU155	1.652E-09	1.103E-09	0.	1.361E+03	0.	1.123E+03	0.	3.089E+01	57
58	PS1	2.916E-07	1.487E-07	0.	1.743E+00	0.	1.178E+01	0.	0.	58
59	PS2	6.715E-07	3.432E-07	0.	7.202E-01	0.	9.056E+00	0.	0.	59
60	PS3	2.173E-07	1.109E-07	0.	2.621E+00	0.	3.165E+01	0.	1.144E+00	60
61	PS4	1.569E-06	8.018E-07	0.	7.580E-02	0.	9.599E+00	0.	6.711E+00	61
					0.		0.		0.	

0- 61
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ISN 73

DT=100.,20,0,2

CONSTANT

POWER

FLUX

FIFA

FLUX

POWER-T

2.051E+03

0.3033

1.882E+14

9.794E+00

(AFTER)

1.045E+01

(AVERAGE)

1.000E+01

(BEFORE)

1.000E+01

(SPEC.)

A/S

9.9998939E-01

IT-FLUX,EIGEN-XS-OT

4 1 1 11

SCAT BY

4.290E-02 'BEO' CRY

FIRST BURN UP RUN

200.0

100.0

BEF-N

200.0

AVE-N

4.290E-02

4.290E-02

4.290E-02

4.337E-02

4.337E-02

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SIGMA

OUTPUT NO. 3

NUCLIDE

0 LEAK

1 BE9

2 O16

3 C12

4 H1

5 D2

6 AL27

7 LI6

8 T3

9 HE3

10 TH232

11 PA233

12 U233

13 U235

14 U238

15 U235

16 U238

17 U236

18 PU239

19 PU240

20 PU241

21 PU242

22 Z999

23 MU95

24 TC99

25 RUI01

26 RUI02

27 RH103

28 RH104

29 RH105

30 PD105

31 PD106

32 PD107

33 PD108

34 AG109

35 CD113

36 I131

37 XE131

38 XE133

39 CS133

200.0

4.290E-02

4.290E-02

4.290E-02

4.337E-02

4.337E-02

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0.

3.162E+21

2.051E+03

0.3033

1.882E+14

9.794E+00

(AFTER)

1.045E+01

(AVERAGE)

1.000E+01

(BEFORE)

1.000E+01

(SPEC.)

A/S

9.9998939E-01

IT-FLUX,EIGEN-XS-OT

4 1 1 11

SCAT BY

4.290E-02 'BEO' CRY

FIRST BURN UP RUN

200.0

100.0

BEF-N

200.0

AVE-N

4.290E-02

4.290E-02

4.290E-02

4.337E-02

4.337E-02

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40	CS134	8.958E-09	5.185E-09	2.128E-09	0	6.166E+01	0.	6.503E+01	0.	1.000E+01	0.	0.	40
41	I135	6.386E-10	6.850E-10	6.477E-10	0	1.581E-02	0.	1.000E+01	0.	1.527E+05	0.	0.	41
42	XE135	4.123E-10	4.423E-10	4.483E-10	0	1.570E+05	0.	1.675E+05	0.	1.112E+05	0.	0.	42
43	CS135	1.658E-07	1.268E-07	8.602E-08	0	3.189E+00	0.	1.532E+01	0.	0.	0.	0.	43
44	ND143	2.300E-07	1.737E-07	1.149E-07	0	1.850E+01	0.	3.189E+01	0.	0.	0.	0.	44
45	ND144	2.017E-07	1.504E-07	9.788E-08	0	6.934E-01	0.	1.000E+01	0.	0.	0.	0.	45
46	ND145	1.518E-07	1.151E-07	7.649E-08	0	2.176E+01	0.	4.603E+01	0.	0.	0.	0.	46
47	ND146	1.346E-07	1.005E-07	6.570E-08	0	1.740E+00	0.	1.000E+01	0.	0.	0.	0.	47
48	PM147	8.522E-08	6.696E-08	4.647E-08	0	6.112E+01	0.	1.328E+02	5.493E+01	0.	0.	0.	48
49	PM148	5.492E-10	4.249E-10	2.772E-10	0	1.376E+03	0.	2.029E+03	0.	7.907E+03	0.	0.	49
50	PM8M*	6.454E-10	4.992E-10	3.485E-10	0	6.645E+03	0.	5.711E+03	0.	1.050E+03	0.	0.	50
51	SM149	6.401E-09	6.490E-09	6.547E-09	0	4.096E+03	0.	3.495E+03	0.	0.	0.	0.	51
52	SM150	6.041E-08	4.367E-08	2.634E-08	0	4.636E+01	0.	7.457E+01	0.	0.	0.	0.	52
53	SM151	2.510E-08	2.072E-08	1.533E-08	0	3.773E+02	0.	3.589E+02	0.	0.	0.	0.	53
54	SM152	3.595E-08	2.629E-08	1.650E-08	0	1.156E+02	0.	1.290E+02	0.	0.	0.	0.	54
55	EUL53	1.894E-08	1.378E-08	8.676E-09	0	9.704E+01	0.	1.062E+02	0.	0.	0.	0.	55
56	EUL54	1.764E-09	1.091E-09	4.985E-10	0	4.985E+02	0.	4.332E+02	0.	7.296E+00	0.	0.	56
57	EUL55	1.926E-09	1.813E-09	1.652E-09	0	1.567E+03	0.	1.299E+03	0.	2.918E+01	0.	0.	57
58	PS1*	5.863E-07	4.418E-07	2.916E-07	0	1.817E+00	0.	1.181E+01	0.	0.	0.	0.	58
59	PS2*	1.341E-06	1.014E-06	6.715E-07	0	7.151E-01	0.	9.042E+00	0.	0.	0.	0.	59
60	PS3*	4.382E-07	3.297E-07	2.173E-07	0	2.661E+00	0.	3.142E+01	0.	1.080E+00	0.	0.	60
61	PS4*	3.141E-06	2.371E-06	1.569E-06	0	7.843E-02	0.	9.607E+00	0.	6.340E+00	0.	0.	61

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MATERIAL 'BEO' AFT.N= 4.2899E-02 OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN XS-OT= 11
 ***** LETH CAPI FISS ABS NU-FISS TRANS CAP2 AS SIGMA
 1 0. 16.10 3.7022E-03 0. -4.4123E-03 0. 7.5924E+00 9.1661E-03 9.0056E+00
 2 16.10 23.00 0. 0. 3.6290E-03 0. 9.2003E+00 0. 9.7987E+00

MATERIAL TH232 AFT.N= 2.1228E-04 OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN XS-OT= 11
 ***** LETH CAPI FISS ABS NU-FISS TRANS CAP2 AS SIGMA
 1 0. 16.10 3.4445E+00 8.7590E-03 3.4521E+00 1.9966E-02 1.2117E+01 1.1693E-03 0.
 2 16.10 23.00 2.3423E+00 0. 2.3423E+00 0. 1.3464E+01 0. 0.

MATERIAL U233 AFT.N= 1.5712E-06 OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN XS-OT= 11
 ***** LETH CAPI FISS ABS NU-FISS TRANS CAP2 AS SIGMA
 1 0. 16.10 7.4901E+00 3.8434E+01 4.5924E+01 9.6174E+01 4.6781E+01 0. 0.
 2 16.10 23.00 2.4026E+01 2.0955E+02 2.3356E+02 5.2383E+02 2.4405E+02 0. 0.

MATERIAL PU239 AFT.N= 9.8061E-06 OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN XS-OT= 11
 ***** LETH CAPI FISS ABS NU-FISS TRANS CAP2 AS SIGMA
 1 0. 16.10 7.9583E+00 1.1269E+01 1.9227E+01 3.2755E+01 2.4926E+01 0. 0.
 2 16.10 23.00 4.3298E+02 7.2307E+02 1.1560E+03 2.0967E+03 1.1376E+03 0. 0.

MATERIAL PU240 AFT.N= 3.4721E-06 OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN XS-OT= 11
 ***** LETH CAPI FISS ABS NU-FISS TRANS CAP2 AS SIGMA
 1 0. 16.10 2.4802E+02 3.6910E-02 2.4806E+02 1.0821E-01 1.3810E+02 0. 0.
 2 16.10 23.00 1.1055E+02 0. 1.1055E+02 0. 1.2044E+02 0. 0.

MATERIAL PU241 AFT.N= 3.0270E-06 OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN XS-OT= 11
 ***** LETH CAPI FISS ABS NU-FISS TRANS CAP2 AS SIGMA
 1 0. 16.10 1.0828E+01 3.3077E+01 4.3906E+01 1.0125E+02 4.5293E+01 0. 0.
 2 16.10 23.00 2.3614E+02 5.7980E+02 8.1582E+02 1.7740E+03 8.0993E+02 0. 0.

MATERIAL 'FP' AFT.N= 1.0000E+00 OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN XS-OT= 11
 ***** LETH CAPI FISS ABS NU-FISS TRANS CAP2 AS SIGMA
 1 0. 16.10 8.9311E-05 0. 9.4205E-05 0. 1.7982E-04 4.9240E-06 0.
 2 16.10 23.00 8.4425E-04 0. 8.4753E-04 0. 1.0196E-03 3.2786E-06 0.

SUMMARY INFORMATION: AFT.N NORMALIZED OUTPUT NO 3, JOB3 , * 23/12/65 * FIRST BURN UP RUN
 (AVE.LETH) FLUX/SOURCE FISS.SPECT TOT.LOSS TOT.NU-FISS TOT.O8**2
 GROUP LETH ENERGY(EV) E(AVE.LETH) 3.1910E+03 2.1088E+02 1.00000 4.0746E-01 1.6622E-01 2.1241E-13
 1 0. 16.10 1.0000E+07 1.0183E+00 3.2325E-02 3.6527E+01 0. 5.9395E-01 9.7729E-01 2.9347E-14
 2 16.10 23.00 1.0183E+00 1.0262E-03 3.2325E-02 3.6527E+01 0. 1.0014E+00 1.1435E+00 2.4176E-13
 1 - 2

		END OF JOB3		TIME TAKEN IN MINS			
TOTAL	JOB	LINKS 1,2	LINKS 3,4	LINKS 5	LINKS 6	LINKS 7	LINKS 8
1.927E+01	1.125E+01	1.158E+00	1.664E+01	6.239E-01	0.		
BURNUP	FLUX	R-RATES	RUREDY	PAUSE			
2.719E-01	8.844E+00	5.743E+00	2.025E-01	0.			

NEW LOAD CALLED

-----GYMEA-----* 23/12/65 *-----GYMEA-----

ISN 6 ,JOB4 EQUILIBRIUM SEARCH
LIB. 1 =GYMEA X-S NDXD * 24/11/65 * MAIN GENERAL PURPOSE LIB. OF 70 NUC.
LIB. 2 =GYMEA SCAT NDSC 3/12/65 'BEO''BE9''O16''H20''C12''C12''300,900,600,1200

T (DAYS) FLUX-T POWER-T FLUX POWER
 1184.2 2.787E+22 1.311E+04 1.3999 2.723E+14 4.747E+00 (AFTER)
 1.107E+01 (AVERAGE)
 4.107E+01 (BEFORE)
 1.100E+01 (SPEC.)
 A/S
 (AVE.N) K B2 82AV CRIT-R-EXTRAP.D(CM) A/S
 (EFF.) 1.0145 4.918E-05
 (INF.) 1.0145 4.918E-05 4.918E-05 4.454E+02 2.597E+00 0.

SCAT BY 0. N00 AT ALL

OUTPUT NO. 1 JOB4 EQUILIBRIUM SEARCH AVE.N*SIGMA*FLUX/SOURCE NEUTRON

NUCLIDE	AFT.N	-DAYS- AVE.N	BEF.N	BURN	CAP1	FISS	ABS	NU-FISS	TRANS	CAP2	DECAY
0 LEAK	1184.2	0.0	0.0				2.957E-13				
1 BE9	4.000E-02	4.000E-02	4.000E-02	0	3.329E-02	0.	-3.775E-02	0.	5.280E+01	8.243E-02	0.
2 O16	4.046E-02	4.046E-02	4.046E-02	0	0.	0.	8.179E-03	0.	3.778E+01	0.	0.
3 C12	0.	0.	0.	0	0.	0.	0.	0.	0.	0.	0.
4 H1	0.	0.	0.	0	0.	0.	0.	0.	0.	0.	0.
5 D2	0.	0.	0.	0	0.	0.	-0.	0.	0.	0.	0.
6 AL27	0.	0.	0.	0	0.	0.	0.	0.	0.	0.	0.
7 LI6	1.126E-06	7.890E-07	0.	0	2.174E-02	0.	2.174E-02	0.	1.838E-02	0.	0.
8 T3	2.096E-06	8.249E-07	0.	0	0.	0.	0.	0.	2.358E-04	0.	0.
9 HE3	2.214E-08	8.160E-09	0.	0	1.325E-03	0.	1.325E-03	0.	1.114E-03	0.	0.
10 TH232	1.815E-04	1.906E-04	2.000E-04	0	1.894E-01	3.752E-04	1.898E-01	8.552E-04	6.830E-01	5.010E-05	0.
11 PA233	5.600E-07	5.685E-07	0.	0	8.956E-03	9.360E-06	8.966E-03	2.261E-05	9.335E-03	0.	0.
12 U232	2.934E-09	1.770E-09	0.	0	1.580E-05	0.	2.001E-05	0.	1.682E-05	0.	0.
13 U233	6.273E-06	4.282E-06	0.	0	1.401E-02	9.638E-02	1.104E-01	2.410E-01	1.049E-01	0.	0.
14 U234	1.526E-06	7.278E-07	0.	0	7.272E-03	3.299E-05	7.305E-03	7.429E-05	8.022E-03	0.	0.
15 U235	4.191E-07	1.484E-07	0.	0	6.310E-04	2.342E-03	2.973E-03	5.717E-03	2.962E-03	0.	0.
16 U236	5.569E-08	1.437E-08	0.	0	5.947E-05	2.366E-07	5.971E-05	5.714E-07	8.544E-05	0.	0.
17 U238	0.	0.	0.	0	0.	0.	0.	0.	0.	0.	0.
18 PU239	6.097E-09	2.835E-06	2.340E-05	0	9.002E-02	1.500E-01	2.400E-01	4.349E-01	2.095E-01	0.	0.
19 PU240	4.189E-08	1.972E-06	5.100E-06	0	1.419E-01	1.636E-05	1.419E-01	4.795E-05	9.613E-02	0.	0.
20 PU241	2.003E-07	2.196E-06	1.500E-06	0	4.289E-02	1.085E-01	1.513E-01	3.319E-01	1.331E-01	0.	3.895E-03
21 PU242	2.774E-06	2.250E-06	0.	0	1.443E-02	0.	1.443E-02	0.	1.918E-02	0.	0.
22 Z2999	0.	0.	0.	0	0.	0.	0.	0.	0.	0.	0.
23 MO95	1.550E-06	1.072E-06	0.	0	1.966E-03	0.	1.966E-03	0.	4.238E-03	0.	0.
24 TC99	1.663E-06	1.254E-06	0.	0	3.782E-03	0.	3.782E-03	0.	5.366E-03	0.	0.
25 RU101	1.503E-06	1.184E-06	0.	0	3.288E-03	0.	3.288E-03	0.	5.566E-03	0.	0.
26 RU102	1.773E-06	1.301E-06	0.	0	3.002E-03	0.	3.002E-03	0.	5.489E-03	0.	0.
27 RH103	6.257E-07	7.190E-07	0.	0	1.332E-02	0.	1.332E-02	0.	9.372E-03	0.	0.
28 RU104	1.227E-06	1.061E-06	0.	0	3.731E-03	0.	3.731E-03	0.	5.597E-03	0.	0.
29 RH105	8.462E-10	2.810E-09	0.	0	2.969E-03	0.	2.969E-03	0.	8.034E-06	0.	0.
30 PD105	1.410E-06	1.034E-06	0.	0	1.310E-03	0.	1.310E-03	0.	2.650E-03	0.	1.578E-02
31 PD106	1.613E-06	1.206E-06	0.	0	6.684E-04	0.	6.684E-04	0.	2.160E-03	0.	0.
32 PD107	1.037E-06	8.090E-07	0.	0	9.428E-04	0.	9.428E-04	0.	2.025E-03	0.	0.
33 PD108	7.400E-07	5.910E-07	0.	0	1.639E-03	0.	1.639E-03	0.	2.426E-03	0.	0.
34 AG109	2.376E-07	2.648E-07	0.	0	4.923E-03	0.	4.923E-03	0.	4.843E-03	0.	0.
35 CD113	4.576E-11	3.698E-10	0.	0	5.570E-04	0.	5.570E-04	0.	4.725E-04	0.	0.
36 I131	5.203E-09	1.210E-08	0.	0	1.136E-04	0.	1.136E-04	0.	1.253E-04	0.	1.266E-02
37 XE131	5.604E-07	5.345E-07	0.	0	6.912E-03	0.	6.912E-03	0.	7.778E-03	0.	0.
38 XE133	5.615E-09	1.383E-08	0.	0	3.452E-04	0.	3.452E-04	0.	3.540E-04	0.	0.
39 CS133	1.523E-06	1.217E-06	0.	0	6.481E-03	0.	6.481E-03	0.	8.733E-03	0.	0.

18. STRUCTURE OF BINARY LIBRARIES

18.0 Summary

Details are given of the structure of GYMEA libraries which consist of binary magnetic tape files. The cards from which the libraries are prepared are discussed elsewhere (19.4). For the cross section library, four types of data exist associated with data for nuclides designated as (1) light - σ_s is given rather than $\nu\sigma_f$ and these nuclides are moderators used for background scattering in resonance theory, (2) unshielded - only cross sections given, (3) shielded - shielding parameters given, and (4) resonance - resonance parameters given.

18.1 GYMEA X-S Library (no scattering matrices) (Max. 120 groups 100 nuclides (or materials))

18.1.0 Summary

A library of NN nuclides contains NN+1 files each separated by a file mark. The library consists of blocked binary records only (≤ 256 words/block). From burnup considerations, the nuclides must be ordered in such a way that, except for a few nuclides, all nuclides are derived from nuclides earlier in the library. Nuclides may be derived from the nuclide immediately following provided no two such requirements appear together. Also from resonance absorption calculation considerations the light nuclides must be first in the library. In general this ordering of nuclides is no handicap to the functioning of the programme as an edit programme is available (Ford 1966 - Section 19). The library may be considered to consist of five parts:

- File 1 - (Library identification 18.1.1
- (Burnup data 18.1.2
- (Group and nuclide information 18.1.3
- Files - (Nuclear data (excluding scattering matrices which are available as part of
- 2-(NN+1) (the SCAT library) 18.1.4 and 18.1.5

File 1

18.1.1 Library identification

1. Heading record of 12 words (ST(I), I = 1, 12)

word 1 = 6H=GYMEA
 word 2 = 6HbX-Sbb
 word 3 = 6HNAMEbb, where NAME = library identifier
 word 4)
 word 12) = 9A6 characters used to describe the library.

2. Second heading record of 5 words (ST(I), I = 1, 5)

word 1 = number of groups (NG)
 word 2 = number of nuclides (NN)
 word 3 = temperature in $^{\circ}\text{K}$ (this is only meaningful for libraries containing shielding parameters rather than resonance parameters)
 word 4 = the group at which resonance theory starts
 word 5 = the group at which resonance theory terminates.

18.1.2 Burnup data

NN records of 10 words each $\begin{pmatrix} \text{DOI} = 1, \text{NN} \\ \vdots \\ \text{ST}(K), K = 1, 10 \end{pmatrix}$. The I^{th} record relates to nuclide I in the library (with nuclear data file (I+1)). The words have the following meanings:

- word 1 = 6HPA233b = nuclide I, say.
- word 2 = 0-4 to indicate fuel type,
(6HNO0bbb if no decay)
- word 3 = (6HU233bb if decay of nuclide I produces U233b
(6HYS0bbb if decay of nuclide I produces a nuclide not included in the
nuclide set (actually any label will do other than NO0bbb or
a library nuclide label).
- word 4 = as above but for neutron capture via $\sigma^{(1)}\phi$
- word 5 = as above but for neutron capture via $\sigma^{(6)}\phi$
- word 6 = λ , the decay constant ($1/10^{25}$ sec)
- word 7 = yield (1)
- word 8 = yield (2)
- word 9 = yield (3)
- word 10 = yield (4)

18.1.3 Group and nuclide information (ST(K), K=1, NG(+1) or
(ST(K), K=1, NN)

1. A record of NG+1 words of the lethargies of the upper boundary of the first group and the lower boundaries of all groups.
2. A record of NG words of the normalized fission spectrum (unit sum).
3. A record of NG words of the group velocities (cm/10⁻⁸ sec)
4. A record of NN fixed point words (IS) used to identify each nuclide in relation to resonance theory:
 - 1 indicates a light nuclide (maximum 10 will be used)
 - 0 indicates a nuclide not requiring resonance theory
 - 1 indicates a nuclide with shielding parameters given
 - ≥2 indicates a nuclide with resonance parameters given.
5. A record of NN words of the nuclide masses (for example 9. for Be⁹).
6. A record of NN words of the material's potential scattering cross section. σ_s .
7. A record of NN words of the fission energy releases (joule/fission).
8. A record of NN words of N's to reconstruct shielded cross sections (1).
9. A record of NN words of N's to reconstruct shielded cross sections (2).
10. A record of NN words of N's to reconstruct shielded cross sections (3).

END OF FILE

18.1.4 Heading record, Ith nuclide (file I+1)

1. A 3-word record to identify the nuclide
 - word 1 = label, for example BE9bbb
 - word 2 = mass, for example 9.
 - word 3 = IS(18.1.3) for example -1
2. A record of NG words indicating the number of trailer records associated with data for each group. In the resonance region each entry indicates the number of resonance parameter sets associated with each group.

18.1.5 Nuclear data, Ith nuclide, given a group at a time (NG items in all)

1. For non-resonance groups

Records of 6 words each (ST(K), K=1,6)

$$\text{word 1} = \sigma^{(1)} = \sigma_{\text{cap 1}} - \text{usually } (n, \gamma)$$

$$\text{word 2} = \sigma^{(2)} = \sigma_f$$

$$\text{word 3} = \sigma^{(3)} = \sigma_a$$

$$\text{word 4} = \sigma^{(4)} = \begin{cases} \nu \sigma_f & \text{for non-light nuclides IS } \neq -1 \\ \sigma_s & \text{for light nuclides IS } = -1 \end{cases}$$

$$\text{word 5} = \sigma^{(5)} = \sigma_{\text{tr}}$$

$$\text{word 6} = \sigma^{(6)} = \sigma_{\text{cap 2}} - \text{usually } (n, 2n).$$

2. For resonance groups

(i) light nuclides, IS = -1,
as for Section 18.1.5 part 1,

(ii) nuclides not affected by resonance theory, IS = 0
as for Section 18.1.5 part 1,

(iii) nuclides with shielding parameters given, IS = 1 :
Records of 6 words each (ST(K), K=1, 6),

$$\text{word 1} = (\sigma_{n, \gamma})_{\infty}$$

$$\text{word 2} = (\sigma_f)_{\infty}$$

word 3 = σ_s , scattering cross section (not used at present)

$$\text{word 4} = (\nu \sigma_f)_{\infty}$$

word 5 = σ_{tr}^* absorption removed ($\simeq \sigma_s$)

$$\text{word 6} = \sigma^{(6)} = \sigma_{\text{cap 2}} ;$$

Trailer records of 6 words each (the trailer record follows the cross section record)
(ST(K), K=1, 6)

$$\text{word 1} = A^{(1)}$$

$$\text{word 2} = A^{(2)}$$

$$\text{word 3} = A^{(4)}$$

$$\text{word 4} = B^{(1)}$$

$$\text{word 5} = B^{(2)}$$

$$\text{word 6} = B^{(4)} ;$$

Then

$$\sigma^{(1)} = (\sigma_{n,\gamma})_{\infty} / \left(1 + \frac{A^{(1)}}{\sigma_p} + \frac{B^{(1)}}{\sigma_p^2} \right)^{1/2}$$

$$\sigma^{(2)} = (\sigma_f)_{\infty} / \left(1 + \frac{A^{(2)}}{\sigma_p} + \frac{B^{(2)}}{\sigma_p^2} \right)^{1/2}$$

$$\sigma^{(3)} = \sigma^{(1)} + \sigma^{(2)}$$

$$\sigma^{(4)} = (\nu\sigma_f)_{\infty} / \left(1 + \frac{A^{(4)}}{\sigma_p} + \frac{B^{(4)}}{\sigma_p^2} \right)^{1/2}$$

$$\sigma^{(5)} = \sigma^{(3)} + \sigma_{tr}^* \text{ absorption removed}$$

$$\sigma^{(6)} = \text{word 6.}$$

(iv) nuclides with resonance parameters given, $IS \geq 2$:
Records of 6 words each (ST(K), K=1, 6),

word 1 = $\sigma_{n,\gamma}^*$ resonance removed

word 2 = σ_f^* resonance removed

word 3 = σ_s (not used at present)

word 4 = $(\nu\sigma_f^*)$ resonance removed

word 5 = σ_{tr}^* absorption removed

word 6 = $\sigma^{(6)} = \sigma_{cap 2}$;

M trailer records of 6 words each (the trailer records follow the cross sections record) as indicated in the file heading record (Section 18.1.4 part 2),

DO1J =1,M
:
:
... ST (K), K=1, 6
:
:

word 1 = E_r (eV)

word 2 = Γ_n "

word 3 = Γ_{γ} "

word 4 = Γ_f "

word 5 = $\nu\Gamma_f$ "

word 6 = g ;

then the cross sections are calculated as for (iii) except that resonance theory is used instead of shielding theory (7.4.2).

END OF FILE

18.2 GYMEA SCAT Library (scattering matrices only)
 (Max.120 groups, 80 max.length of scatter)

18.2.0 Summary

A library with data for NS materials and given for NT temperatures contains $1+NS \times NT$ files, each separated by a file mark. The library consists of blocked binary records only (≤ 256 words/block). This library will be positioned by the IBM 1401 while the IBM 7040 is loading up cross section data (using Richardson's SEARCH 1965). As a result of this, machine time will not be wasted with a library containing many different temperatures. The library may be considered to consist of 4 parts:

- (Library identification (18.2.1)
- (
- File 1 - (Scattering matrix identifiers (18.2.2)
- (
- (Temperature list (18.2.3)
- Files - Matrices (18.2.4)
- 2-($1+NS \times NT$)

File 1

18.2.1 Library identification

1. Heading record of 12 words (ST(K), K=1, 12)

word 1 = 6H=GYMEA

word 2 = 6HbSCATb

word 3 = 6HNAMEbb (not necessarily the same as the cross section library)

word 4)= 9A6 characters used to describe the library.
 word 12)

2. Second heading record of 3 words (ST(K), K=1, 3)

word 1 = number of groups (NG - must agree with the cross section library)

word 2 = number of scattering matrices/temperature (NS)

word 3 = number of temperatures (NT).

(Note - it is assumed that every temperature has a complete set of scattering matrices).

18.2.2 Scattering matrix identifiers

A record of $2 \times NS$ words (ST(K), K=1, $2*NS$). Each group of two words identifies the material and its manner of preparation; for example a library may read

- word 1 = 'BEO'
- word 2 = CRY
- word 3 = BE9 (stored as 6HBE9bbb)
- word 4 = GAS

18.2.3 Temperature list

A record of NT words (ST(K), K=1,NT). Each word is a temperature in °K and is given in any order. For example the list may read

word 1 = 900.
word 2 = 300.
word 3 = 600.
word 4 = 1200.

END OF FILE

18.2.4 Heading record, Ith material Jth temperature (file 1+IxJ) (the next file would be for the (I+1)th material at the Jth temperature).

A 3-word record which identifies the matrix and which must agree with information given previously (ST(K), K=1, 3).

word 1 = material, for example 'BEO'
word 2 = preparation, for example CRY
word 3 = temperature, for example 300.

18.2.5 Scattering data, Ith material, Jth temperature

After the heading record follow NG records of a variable number of words thus

```
DO 1 J=1, NG
  :
  :
  . . . N, LGG, ST(K), K=1, N
  :
  :
```

where N is the length of scatters (max.81) and LGG is the position of the 'self-scatter' term (actually the negative sum of the outscatters).

END OF FILE

19. LIBRARIES AND EDIT PROGRAMMES IN BRIEF19.0 Summary

The content of two main libraries used at the A.A.E.C. is given. Other libraries are discussed as well as the programmes used to edit GYMEA libraries.

19.1 GYMEA X-S Library NDXD,

(120 groups, 70 nuclides)

19.1.1 Group boundaries

GROUP	LETH	ENERGY (EV)	E (AVE. LETH)	
1	0.	1.75	1.0000E+07	
2	1.75	2.25	1.7377E+06	
3	2.25	2.50	1.0540E+06	
4	2.50	2.75	8.2085E+05	
5	2.75	3.00	6.3928E+05	
6	3.00	3.25	4.9787E+05	
7	3.25	3.50	3.8774E+05	
8	3.50	3.75	3.0197E+05	
9	3.75	4.00	2.3518E+05	
10	4.00	4.25	1.8316E+05	
11	4.25	4.50	1.4264E+05	
12	4.50	4.75	1.1109E+05	
13	4.75	5.00	8.6517E+04	
14	5.00	5.25	6.7379E+04	
15	5.25	5.50	5.2475E+04	
16	5.50	5.75	4.0868E+04	
17	5.75	6.00	3.1828E+04	
18	6.00	6.25	2.4788E+04	
19	6.25	6.50	1.9305E+04	
20	6.50	6.75	1.5034E+04	
21	6.75	7.00	1.1709E+04	
22	7.00	7.25	9.1188E+03	
23	7.25	7.50	7.1017E+03	
24	7.50	7.75	5.5308E+03	
25	7.75	8.00	4.3074E+03	
26	8.00	8.25	3.3546E+03	
27	8.25	8.50	2.6126E+03	
28	8.50	8.75	2.0347E+03	
29	8.75	9.00	1.5846E+03	
30	9.00	9.25	1.2341E+03	
31	9.25	9.50	9.6112E+02	
32	9.50	9.75	7.4852E+02	
33	9.75	10.00	5.8295E+02	
start of resonance region	34	10.00	10.25	4.5400E+02
	35	10.25	10.50	3.5358E+02
	36	10.50	10.75	2.7536E+02
	37	10.75	11.00	2.1445E+02
	38	11.00	11.25	1.6702E+02
	39	11.25	11.50	1.3007E+02
	40	11.50	11.75	1.0130E+02
	41	11.75	12.00	7.8893E+01
	42	12.00	12.25	6.1442E+01
	43	12.25	12.50	4.7851E+01
	44	12.50	12.75	3.7267E+01
	45	12.75	13.00	2.9023E+01
	46	13.00	13.25	2.2603E+01

47	13.25	13.50	1.7603E+01	1.3710E+01	1.5535E+01
48	13.50	13.75	1.3710E+01	1.0677E+01	1.2099E+01
49	13.75	13.90	1.0677E+01	9.1898E+00	9.9056E+00
50	13.90	14.00	9.1898E+00	8.3153E+00	8.7416E+00
51	14.00	14.10	8.3153E+00	7.5240E+00	7.9097E+00
52	14.10	14.20	7.5240E+00	6.8080E+00	7.1570E+00
53	14.20	14.30	6.8080E+00	6.1601E+00	6.4760E+00
54	14.30	14.40	6.1601E+00	5.5739E+00	5.8597E+00
55	14.40	14.50	5.5739E+00	5.0435E+00	5.3021E+00
56	14.50	14.60	5.0435E+00	4.5635E+00	4.7975E+00
57	14.60	14.70	4.5635E+00	4.1292E+00	4.3410E+00
58	14.70	14.80	4.1292E+00	3.7363E+00	3.9279E+00
59	14.80	14.90	3.7363E+00	3.3807E+00	3.5541E+00
60	14.90	15.00	3.3807E+00	3.0590E+00	3.2159E+00
61	15.00	15.10	3.0590E+00	2.7679E+00	2.9098E+00
62	15.10	15.20	2.7679E+00	2.5045E+00	2.6329E+00
63	15.20	15.30	2.5045E+00	2.2662E+00	2.3824E+00
64	15.30	15.40	2.2662E+00	2.0505E+00	2.1557E+00
65	15.40	15.50	2.0505E+00	1.8554E+00	1.9505E+00
66	15.50	15.60	1.8554E+00	1.6788E+00	1.7649E+00
67	15.60	15.70	1.6788E+00	1.5191E+00	1.5970E+00
68	15.70	15.80	1.5191E+00	1.3745E+00	1.4450E+00
69	15.80	15.90	1.3745E+00	1.2437E+00	1.3075E+00
end of resonance region	70	15.90	16.00	1.2437E+00	1.1254E+00
	71	16.00	16.10	1.1254E+00	1.0183E+00
	72	16.10	16.20	1.0183E+00	9.2136E-01
	73	16.20	16.30	9.2136E-01	8.3368E-01
	74	16.30	16.40	8.3368E-01	7.5435E-01
	75	16.40	16.50	7.5435E-01	6.8256E-01
	76	16.50	16.60	6.8256E-01	6.1761E-01
	77	16.60	16.70	6.1761E-01	5.5883E-01
	78	16.70	16.80	5.5883E-01	5.0565E-01
	79	16.80	16.90	5.0565E-01	4.5753E-01
	80	16.90	17.00	4.5753E-01	4.1399E-01
	81	17.00	17.10	4.1399E-01	3.7460E-01
	82	17.10	17.20	3.7460E-01	3.3895E-01
	83	17.20	17.30	3.3895E-01	3.0669E-01
	84	17.30	17.40	3.0669E-01	2.7751E-01
	85	17.40	17.50	2.7751E-01	2.5110E-01
	86	17.50	17.60	2.5110E-01	2.2720E-01
	87	17.60	17.70	2.2720E-01	2.0558E-01
	88	17.70	17.80	2.0558E-01	1.8602E-01
	89	17.80	17.90	1.8602E-01	1.6832E-01
	90	17.90	18.00	1.6832E-01	1.5230E-01
	91	18.00	18.10	1.5230E-01	1.3781E-01
	92	18.10	18.20	1.3781E-01	1.2469E-01
	93	18.20	18.30	1.2469E-01	1.1283E-01
	94	18.30	18.40	1.1283E-01	1.0209E-01
	95	18.40	18.50	1.0209E-01	9.2375E-02
	96	18.50	18.60	9.2375E-02	8.3584E-02
	97	18.60	18.70	8.3584E-02	7.5630E-02
	98	18.70	18.80	7.5630E-02	6.8433E-02
	99	18.80	18.90	6.8433E-02	6.1920E-02
	100	18.90	19.00	6.1920E-02	5.6028E-02
	101	19.00	19.10	5.6028E-02	5.0696E-02
	102	19.10	19.20	5.0696E-02	4.5872E-02
	103	19.20	19.30	4.5872E-02	4.1507E-02
	104	19.30	19.40	4.1507E-02	3.7557E-02

105	19.40	19.50	3.7557E-02	3.3983E-02	3.5725E-02
106	19.50	19.60	3.3983E-02	3.0749E-02	3.2325E-02
107	19.60	19.70	3.0749E-02	2.7823E-02	2.9249E-02
108	19.70	19.80	2.7823E-02	2.5175E-02	2.6466E-02
109	19.80	19.90	2.5175E-02	2.2779E-02	2.3947E-02
110	19.90	20.00	2.2779E-02	2.0612E-02	2.1668E-02
111	20.00	20.10	2.0612E-02	1.8650E-02	1.9606E-02
112	20.10	20.20	1.8650E-02	1.6875E-02	1.7741E-02
113	20.20	20.30	1.6875E-02	1.5269E-02	1.6052E-02
114	20.30	20.40	1.5269E-02	1.3816E-02	1.4525E-02
115	20.40	20.50	1.3816E-02	1.2502E-02	1.3142E-02
116	20.50	21.00	1.2502E-02	7.5826E-03	9.7362E-03
117	21.00	21.50	7.5826E-03	4.5991E-03	5.9053E-03
118	21.50	22.00	4.5991E-03	2.7895E-03	3.5817E-03
119	22.00	22.50	2.7895E-03	1.6919E-03	2.1724E-03
120	22.50	23.00	1.6919E-03	1.0262E-03	1.3177E-03

19.1.2 Nuclides

6 light nuclides, 16 other main nuclides, 35 individual fission products, 4 pseudo fission products, 9 control nuclides, (R = resonance parameters given).

1	Be ⁹		27	Rh ¹⁰³	54	Sm ¹⁵²	
2	O ¹⁶		28	Ru ¹⁰⁴	55	Eu ¹⁵³	
3	C ¹²		29	Rh ¹⁰⁵	56	Eu ¹⁵⁴	
4	H ¹		30	Pd ¹⁰⁵	57	Eu ¹⁵⁵	
5	D ²		31	Pd ¹⁰⁶			
6	Al ²⁷		32	Pd ¹⁰⁷	58	'PS1'	
7	Li ⁶		33	Pd ¹⁰⁸	59	'PS2'	
8	T ³		34	Ag ¹⁰⁹	60	'PS3'	
9	He ²		35	Cd ¹¹³	61	'PS4'	
10	Th ²³²	R	36	I ¹³¹			
11	Pa ²³³	R	37	Xe ¹³¹	62	B ¹⁰	
12	U ²³²		38	Xe ¹³³	63	Cr ⁵²	
13	U ²³³	R	39	Cs ¹³³	64	Mn ⁵⁵	R
14	U ²³⁴	R	40	Cs ¹³⁴	65	Fe ⁵⁶	
15	U ²³⁵	R	41	I ¹³⁵	66	Ni ⁵⁸	
16	U ²³⁶	R	42	Xe ¹³⁵	67	Co ⁵⁹	
17	U ²³⁸	R	43	Cs ¹³⁵	68	Cu ⁶³	
18	Pu ²³⁹	R	44	Nd ¹⁴³	69	W ¹⁸⁶	R
19	Pu ²⁴⁰	R	45	Nd ¹⁴⁴	70	Au ¹⁹⁷	R
20	Pu ²⁴¹	R	46	Nd ¹⁴⁵			
21	Pu ²⁴²		47	Nd ¹⁴⁶			
22	ZZ ⁹⁹⁹		48	Pm ¹⁴⁷			
			49	Pm ¹⁴⁸			
23	Mo ⁹⁵		50	'Pm ^{6M1}			
24	Tc ⁹⁹		51	Sm ¹⁴⁹			
25	Ru ¹⁰¹		52	Sm ¹⁵⁰			
26	Ru ¹⁰²		53	Sm ¹⁵¹			

19.2 GYMEA SCAT Library NDSC,

120 groups 6 matrices at 4 temperatures

19.2.1 Materials

'REO'	CRY	crystal scattering
'BE9'	GAS	monatomic gas scattering
'O16'	GAS	" " "
'H2O'	EFF	effective width model scattering
'C12'	CRY	crystal scattering
'C12'	GAS	monatomic gas scattering

19.2.2 Temperatures

- 300 °K
- 900 °K
- 600 °K
- 1200 °K

19.3 Other Libraries

The cross section library NDXD was prepared from a library consisting of 100 nuclides (NDXC- Cook 1966) by mixing 43 fission product dregs into 4 pseudo fission products. It has been found that little loss of accuracy results from this procedure for long-term reactivity predictions. Other libraries of both fewer nuclides (more mixed fission products) and fewer groups have been prepared from time to time to suit the particular needs of people engaged in reactor assessment.

Additional scattering matrix libraries are available. The data for all scattering libraries was prepared by McGregor (1965a) using the programmes LEAP (McLatchie 1962), PIXSE (Macdougall 1963), and GAM-1 (Joanou and Dudek 1961).

19.4 Edit Programmes in Brief (EDITOR and SCATLD)

Details have been given in the previous sections for the structure of libraries as required by the code GYMEA. However those engaged in library maintenance are necessarily more interested in the structure of BCD cards from which an original library is loaded or which comprise corrections to be made to an existing library. Two library edit programmes are available- EDITOR (Ford 1966) for the cross section library, and SCATLD (McGregor 1965b) for the scattering matrix library. EDITOR has been made very versatile to cope with the detailed needs of those maintaining the cross section library and uses free input. SCATLD does not need this flexibility, as maintenance is normally only required at the level of one complete scattering matrix, and fixed format input is therefore used.

20. PROGRAMME SOURCE STRUCTURE

The structure of the programme is set out by chains below. Each of the subroutines required in a chain is given, with entries in columns of those subroutines which it calls. Subroutines given in inverted commas have more than one entry point and an (M) following a subroutine means that it is included in chain (0).

<u>CHAIN (0)</u>	<u>CHAIN (0)</u>	<u>CHAIN (1)</u>
'CLOCK'	FPT	CALL
INTVLL	-	HOLLRH (M)
CHAIN (1)	HOLLRH	CALL
CHAIN (2)	-	CHAIN (2)
PSLIB		ERROR
RUREDY		ERROR 3
'SEARCH'	<u>CHAIN (1)</u>	ERROR 3
SKIP	'CLOCK' (M)	-
RUREDY	SKIP (M)	STAN
'SEARCH'	SRCH	HOLLRH (M)
'CLOCK'	COMPIL	PACK
SKIP	ERROR	-
CHAIN (1)	CHAIN (2)	DATE
SEARCH	SRCH	-
-	-	
COLLEC	COMPIL	
CONC	'FREE'	<u>CHAIN (2)</u>
CONC	PACK	'CLOCK' (M)
-	DATE	SKIP (M)
LOOK	ERROR	MAIN
TEST	STAN	CALL
TEST	'FREE'	MAIN
-	LHFREE	CHAIN (1)
NUCL	RHFREE	SCEASE
-	ERFREE	'FREE'
CLOCK	LHFREE	PACK
-	ERROR	'CLOCK'
CHNRTN	RHFREE	TEST
-	ERROR	HOLLRH (M)
	ERFREE	ERROR

CHAIN (2) (contd.)

UNDO	LIBER	CALL
SDAT	STAN	MAIN
CALLCH	DUMP	SDAT
GETL	CALLCH	'FREE'
MAKET	FINISH	PACK
PUTN	WEST	ERROR
SORTN	NUC	SEND
TYPE	NSS	'CLOCK'(M)
GNAME	MAKET	MAIN
CALL	GETL	SCEASE
RETURN	'FREE'	'SEARCH'
COLLEC (M)	PUTN	'FREE'
UNLOAD	ERROR	PACK
BUT	GETL	MAIN
SEND	ERROR	LIBER
PART	PACK	CALL
ORDER	LOOK (M)	DAT
PSLIB (M)	'FREE'	CALL
'SEARCH' (M)	DUMP	ERROR
NUCL (M)	-	CHAIN (1)
CONC (M)	STAN	FINISH
SKIP (M)	HOLLRH (M)	DUMP
TEST (M)	GNAME	BUT
CALL	'FREE'	LOOK
ERROR	PACK	ERROR
MAIN	SORTN	PART
SDAT	PUTN	ERROR
SEND	PUTN	COLLEC (M)
SCEASE	-	AETC
'SEARCH' (M)	NSS	NUC
'CLOCK' (M)	TEST (M)	COLLEC (M)
DAT	RETURN	AETC

CHAIN (2) (contd.)

SORTN
 AETC
 -
 WEST
 WESTSP
 WESTSP
 -
 CALLCH
 CLOCK
 CHAIN (3)
 ERROR
 UNLOAD
 -
 UNDO
 -
 TYPE
 -
 ORDER
 -
 PACK
 -
 'FREE'
 as for CHAIN (1)
 EXTRA2
 -

CHAIN (3)

WATCH
 GYMEA2
 EQU
 OUTPUT
 SAVE
 CHAIN (2)
 CHAIN (5)
 CHAIN (6)
 SKIP (M)
 GYMEA2
 ERROR
 BURNUP
 CHAIN (4)
 KETC
 EIGENV
 TGRALS
 OUTPUT
 EIGENV
 ERROR
 COLLEC (M)
 SORTN
 EXPL
 ERROR
 TGRALS
 ERROR
 KETC
 -
 EQU
 ERROR
 GYMEA2
 OUTPUT
 BURNUP

WATCH
 CODE
 CODE
 -
 OUTPUT
 NUCL (M)
 ANORM
 HOLLRH (M)
 WDS
 -
 ANORM
 -
 SAVE
 COLLEC (M)
 ANORM
 WDS
 WATCH
 CLOCK (M)
 ERROR
 CHAIN (1)
 SORTN
 PUTN
 PUTN
 -

CHAIN (4)

FLUXBL
RATES
CHAIN (3)
FLUXBL
PSLIB (M)
ERROR
STOD
WEST
SKIP (M)
FIRSTD
BLOCKD
RUREDY (M)
SCATD
SEDL
WATCH
RATES
WATCH
COLLEC (M)
SKIP
FIRSTD
BLOCKD
VERLAP
CAM
ERROR
SCATD
RUREDY (M)
COLLEC (M)
LOOK (M)
ERROR
SKIP (M)
PSLIB (M)

FIRSTD
ERROR
SAME
WISIT
CAM
NUCL (M)
SKIP (M)
BLOCKD
SHIELD
RES
RESH
SHIELD
-
RESH
-
RES
RESJ
RESJ
-
STOD
MUGN
MUGN
-
CAM
-
WISIT
ERROR
VERLAP
-
SEDL
SIDW

SIDW
-
WEST
WESTSP
WESTSP
-
SAME
-
WATCH
CLOCK (M)
ERROR
CHAIN (1)

CHAIN (5)

CLOCK (M)
 COLLEC (M)
 PRINTX
 PCRAM
 PZOOM
 PDSN
 PWDSN
 PGMCM9
 PLSCD
 CHAIN (2)
 CHAIN (6)
 'SGG '
 PSLIB (M)
 RMAT
 SORTS
 RUREDY (M)
 ERROR
 HOLLRH (M)
 PRINTX
 HOLLRH (M)
 GCOL
 ANORM
 COLLEC (M)
 RMAT
 CONC (M)
 NUCL (M)
 PRINTS
 PRINTN
 WDS
 PCRAM
 CHEAD
 GCOL

RMAT
 WDS
 'SGG '
 NUCL (M)
 PZOOM
 CHEAD
 GCOL
 HOLLRH (M)
 RMAT
 NUCL (M)
 WDS
 'SGG '
 CONC (M)
 PDSN
 CHEAD
 HOLLRH (M)
 'SGG '
 ERROR
 GCOL
 RMAT
 NUCL (M)
 PWDSN
 CHEAD
 RMAT
 NUCL (M)
 WDS
 'SGG '
 PGMCM9
 CHEAD
 RMAT
 NUCL (M)
 WDS

PRINTS
 PRINTN
 'SGG '
 PRINTN
 GCOL
 RMAT
 NUCL
 SORTS
 SLOOK
 NUCL
 CONC
 SLOOK
 ERROR
 'SGG '
 PLSCD
 WDS
 ERROR
 CHAIN (1)
 ANORM
 CHEAD

CHAIN (6)

SHIELD
MATS
NUCL (M)
CHAIN (2)
LOT13
WATCH
RESH
SAME
LOT14
PSLIB (M)
RUREDY (M)
SKIP (M)
WDS
ANORM
WRT
MIXEM
BUT
ERROR
UNLOAD
TYPE
ANORM
-
WRT
MIXEM
MATS
MIXEM
-
BUT
LOOK (M)
ERROR
UNLOAD
-

ERROR
CHAIN (1)
WDS
-
TYPE
(not the same as the
subroutine used in CHAIN (2))
-
LOT13
-
LOT14
-
MATS
-
SHIELD
-
RESH
-
SAME
-
WATCH
'CLOCK' (M)

21. PROGRAMME COMMON

A list of the locations of all variables in blank common is given below. This is very useful in implementing two features of the programme which the general user should not require.

The first use is in making accessible variables not included in the user's common. Use of the loop-round feature of core storage gives the relationships

// ≡ X (19491)

// ≡ Z (17491)

// ≡ J (19501)

where // is the symbol used in the listing for the beginning of common. Thus any variables in common can be used, but they should be altered only with great caution. As an example, the lechary boundaries U(1) to U(121) are accessible as X(32277) to X(32397).

The other use is in conjunction with the expanded error dump facility for programme debugging (13.4).

161	SED	EQU	//+11739	BSED	EQU	//+11742	SIDAC	EQU	//+11751
164	NSED	EQU	//+11752	ISED	EQU	//+11753	KSED	EQU	//+11754
167	JSED	EQU	//+11755	LSED	EQU	//+11756	ICSED	EQU	//+11757
170	JSED1	EQU	//+11758	JSED2	EQU	//+11759	SEDI	EQU	//+11760
173	ETC3	EQU	//+11761	NUCLA	EQU	//+11764	LIBX	EQU	//+11905
176	ITRAN	EQU	//+11906	RPART1	EQU	//+11907	LL	EQU	//+11908
179	KTR	EQU	//+11909	NOT	EQU	//+11910	ICH2	EQU	//+11911
182	NMD	EQU	//+11912	NY	EQU	//+11913	NGD2	EQU	//+11914
185	NR	EQU	//+11915	NS	EQU	//+11916	NT	EQU	//+11917
188	NPART	EQU	//+11918	NYS	EQU	//+11919	LCR	EQU	//+11920
191	ICRT	EQU	//+11921	IPAG	EQU	//+11922	STEMP	EQU	//+11923
194	PTEMP	EQU	//+11924	RTEMP	EQU	//+11925	NTEMP	EQU	//+11926
197	NSCAT	EQU	//+11927	NLIB	EQU	//+11928	NSLIB	EQU	//+11929
200	TEV	EQU	//+11930	LGL	EQU	//+11931	LXG	EQU	//+11932
203	LXGL	EQU	//+11933	LSG	EQU	//+11934	LSGL	EQU	//+11935
206	PMIC	EQU	//+11936	IWT	EQU	//+11937	TCARD	EQU	//+11938
209	IHEAD	EQU	//+11948	DEFN	EQU	//+11955	NHIS	EQU	//+12155
212	LMAT	EQU	//+12195	NYN	EQU	//+12245	LPART	EQU	//+12345
215	KSCATL	EQU	//+12365	KSCATP	EQU	//+12385	LSCAT	EQU	//+12405
218	ISCR	EQU	//+12425	IBURN	EQU	//+12525	NGS	EQU	//+12625
221	XAN	EQU	//+12746	U	EQU	//+12786	V1	EQU	//+12907
224	SETP	EQU	//+13027	AM	EQU	//+13037	JSAVE	EQU	//+13137
227	MSAVE	EQU	//+13143	LOT	EQU	//+13193	MOT	EQU	//+13208
230	ITL1E1	EQU	//+13223	ITL2E	EQU	//+13235	ETC5	EQU	//+13247
233	KG2	EQU	//+13251	IOLA	EQU	//+13252	ACETA	EQU	//+13253
236	ITSA	EQU	//+13254	IFPA	EQU	//+13255	X7	EQU	//+13256
239	X9	EQU	//+13257	SAV	EQU	//+13258	G1	EQU	//+13259
242	G2	EQU	//+13260	G	EQU	//+13261	LSGC	EQU	//+13262
245	NGDC	EQU	//+13263	NSC	EQU	//+13264	NTC	EQU	//+13265
248	NTEMPC	EQU	//+13266	JINO	EQU	//+13267	JIN	EQU	//+13268
251	X0	EQU	//+13277	X	EQU	//+13278	TA	EQU	//+13288
254	DTD	EQU	//+13289	AFL	EQU	//+13290	APR	EQU	//+13291
257	RPR	EQU	//+13292	VPR	EQU	//+13293	BPR	EQU	//+13294
260	FIN	EQU	//+13295	PIM	EQU	//+13296	VIN	EQU	//+13297
263	XIN	EQU	//+13298	FIFA	EQU	//+13299	AKI	EQU	//+13300
266	AKE	EQU	//+13301	XI	EQU	//+13302	XE	EQU	//+13303
269	XAV	EQU	//+13304	CNA	EQU	//+13305	SKI	EQU	//+13306
272	SKE	EQU	//+13307	SXA	EQU	//+13308	STR	EQU	//+13309
275	GL	EQU	//+13310	CG	EQU	//+13311	GU	EQU	//+13312
278	AER	EQU	//+13313	ACER	EQU	//+13314	ACEI	EQU	//+13315
281	FIFC	EQU	//+13316	TEMP	EQU	//+13317	TLAST	EQU	//+13318
284	BPRS	EQU	//+13319	XY	EQU	//+13320	EPART	EQU	//+13323
287	VOID	EQU	//+13324	ETC1	EQU	//+13325	AN	EQU	//+13328
290	AV	EQU	//+13428	BN	EQU	//+13528	SXS	EQU	//+13628
293	FER	EQU	//+13728	GFLUX	EQU	//+13828	GSCE	EQU	//+13948
296	GB2	EQU	//+14068	WU	EQU	//+14188	WV	EQU	//+14308
299	DB2	EQU	//+14428	SIDC	EQU	//+14548	SEDET	EQU	//+14668
302	DSED	EQU	//+14787	SEDAC	EQU	//+15144	ISED1	EQU	//+15145
305	ISED3	EQU	//+15146	INSED	EQU	//+15147	DC	EQU	//+15148
308	ETC2	EQU	//+15248	Z0	EQU	//+15277	Z	EQU	//+15278
311	IA	EQU	//+161	W	EQU	//+1764	NH	EQU	//+1764
314	MH	EQU	//+1764	IST	EQU	//+1764	AST	EQU	//+1764
317	CRIT	EQU	//+13310	SIDF	EQU	//+13828	R	EQU	//+14548
320	WX	EQU	//+15778	IN	EQU	//+15778	WY	EQU	//+15898
323	JW	EQU	//+15898	SPA	EQU	//+15898	WZ	EQU	//+16018

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```

326 EQU          //+13278
329 IDEFN       EQU          //+11955
332 LCODE      EQU          CONST+30
335 IDIS       EQU          CONST+86
338 ILC        EQU          CONST+89
341 ILF        EQU          CONST+92
344 ILJ        EQU          CONST+95
347 ILP        EQU          CONST+98
350 ILT        EQU          CONST+101
353 ILZ        EQU          CONST+104
356 ILAST      EQU          CONST+107
359 IBRAK      EQU          CONST+110
362 ILDADA     EQU          CONST+113
365 IGYMEA     EQU          CONST+116
368 SFMT      EQU          CONST+119
371 EVEN       EQU          CONST+121
374 //         CONTRL //
377 *          USE          PRGCT.
380 PROLOGUE   EXTERN     SETFP.
383          BEGIN     SFLCT.,*
386          USE          SFLCT.,*

326 EQU          //+258
329 JSUB       EQU          //+14188
332 ICEA       EQU          CONST+45
335 KSUB       EQU          CONST+87
338 ILE        EQU          CONST+90
341 ILI        EQU          CONST+93
344 ILO        EQU          CONST+96
347 ILS        EQU          CONST+99
350 ILX        EQU          CONST+102
353 ILDOL      EQU          CONST+105
356 ICOM       EQU          CONST+108
359 ILDA       EQU          CONST+111
362 ILG82      EQU          CONST+114
365 ILSCAT     EQU          CONST+117
368          CONST+121
371          00203
374          00203
377          00001 P.0010
380          00002
383          USE          USE
386          BEGIN     BEGIN STRCT.,*

326 EQU          //+428
329 JSUB       EQU          CONST+0
332 ICEA       EQU          CONST+85
335 KSUB       EQU          CONST+88
338 ILE        EQU          CONST+91
341 ILI        EQU          CONST+94
344 ILO        EQU          CONST+97
347 ILS        EQU          CONST+100
350 ILX        EQU          CONST+103
353 ILDOL      EQU          CONST+106
356 ICOM       EQU          CONST+109
359 ILDA       EQU          CONST+112
362 ILG82      EQU          CONST+115
365 ILSCAT     EQU          CONST+118
368          CONST+CONST+122
371          //
374          USE          STRCT.
377          S-JXIT
380          EXTERN   S-JXIT
383          USE          PRGCT.
386          BEGIN     STRCT.,*

```

\$DKEND COMMON

NO MESSAGES FOR ABOVE ASSEMBLY

END

22. FREE INPUT SUBROUTINE

22.0 Summary

The routine used has four main features:

- (1) numeric data may be in any form readable by eye and alphabetic data may be ignored,
- (2) a limited alphabetic mode is available in which blanks only are ignored,
- (3) as much control of the reading and interpretation process as possible is left under programmer control,
- (4) data can be read from any input device (up to column 72 of a record) or from storage.

For the reading of floating point data, exponents may be simplified by omitting either the 'E' or the sign; blanks may be present in the exponent. Thus .01, 1.E-b2, 1.-2, 2.E2, 2.E+02 are all acceptable. If no decimal point is present, one is assumed after the number.

For any numeric data, the repeat notation, for example 5*.1 meaning a set of five .1's may be used. No blanks are allowed between the numbers here. Thus 5* and *.1 are both completely ignored. A card with an * in column one is ignored.

When alphameric data is read, each character from the first non-blank one is stored in the given vector in the lower part of each word. Thus on the IBM 7040 the character A is equivalent to 21 octal, that is 17 decimal, and if 3 characters are requested the string AbB becomes 17,48,18.

22.1 Use

The routine consists of two MAP subroutines. The first has the 3 normal entry points and the second has the auxiliary routines used, which may be replaced.

The three entry points to the routine for reading floating point, fixed point, and alphameric data respectively are:

```
CALL  EFREE (IT,IC,A,N)
CALL  IFREE (IT,IC,IA,N)
CALL  AFREE (IT,IC,IA,N).
```

In each of these calling sequences IT is the FORTRAN unit number (if negative this implies reading from storage; see later), IC is the column from which interpretation is to begin, A or IA is the variable or vector to be read, and N is the number of elements to be read. The column count may be set as desired and the routine then continues to update it. IC must be greater than 72 for a new record to be read, in which case the (IC/72) card is read and IC is used mod 72 for the column count. If when column 72 is reached a request is not satisfied, a new record is read and IC set back to 1. Thus if no control is required, IC is set to 73 at the beginning and then left for the routine to use.

When data is being read from storage, records are assumed to be stored in blank common under format 12A6. Then if IT is negative the next record used is the Mth stored card where M = -IT. Thus IT is updated by the routine when each new card is read if this feature is used. The routine PACK (IA,N) can be used to pack the right hand character of the N(≤6) elements of the vector IA into the first element of IA in normal A format form.

To provide further control of the routine by the programmer, any of the routines ERFREE, LHFREE, and RHFREE may be rewritten by him.

When an error is detected,

ERFREE (IT,IND,I,ICC,BUF)

is called. Here IT is the unit number; IND is one of the alphameric words bb.E, bbIM, bbIS, bEND, for an exponent or decimal error, a mode error, a scaling error, or a '\$' control card respectively; I is the number of elements of the vector already filled; ICC is the present column count; and BUF contains the card image (format 12A6). From this information an appropriate message is printed.

When column 72 of a card is reached,

RHFREE (IT, IMODE, I, N)

is called. Here $IMODE = \begin{Bmatrix} 1 \\ 0 \\ -1 \end{Bmatrix}$ for a call from $\begin{Bmatrix} EFREE \\ IFREE \\ AFREE \end{Bmatrix}$,

and IT, N, I are as above.

When a new record is placed in the working buffer for FREE,

LHFREE (IT, BR, BUF)

is called, where RR is negative for a carry over from a previous record.

In the standard version these last two do nothing. The working buffer, which consists of the 72 characters stored in the right hand end of the words 2-73 of the labelled common area /FR/ of length 76, is accessible as long as the area in FREE is actually used.

22.2 Adaptations for GYMEA

This routine is used to give input to the code the following features:

- (1) All data may be punched anywhere on a card except in the few cases noted.
- (2) On FORTRAN-like cards the repeat notation is discarded (as * means multiplication), and the data for each statement must be on one card only.
- (3) For all other cards, repeat notation may be used, and a set of data may be continued over many cards by having a non-descending integer in column 1 for cards other than the first of a set.
- (4) Input cards are stored and read from storage as a programme
- (5) Cards with an '*' in column one are simply listed.
- (6) Where the input cards RX etc.(10) are included, the standard version of FREE is used for reading.

23. ACKNOWLEDGEMENTS

The code GYMEA forms part of a larger entity now called the GYMEA PROJECT. This project is to maintain and extend a reliable nuclear data file for reactor calculations to be carried out at the A.A.E.C. Research Establishment. To this end, the talents, inspirations, and hard work of many people have gone into the project. Work of this project has been carried out over a period of several years, although the code GYMEA is of more recent origin. The authors therefore thank the following people for contributions to the project.

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- Mr. J. Cook – nuclear data, particularly fission product data,
- Mr. B. McGregor – scattering matrix preparation and editing programme SCATLD,
- Mr. K. Maher – data transcription and preliminary correlation with experiment,
- Miss E. Kletzmayer – comparison of various resonance theories,
- Mrs. L. Wall – comparison of GYMEA data with the data of Hughes and Schwartz (1958),
- Miss J. Robertson – organization of computer runs;

Dr. G. Ford, IBM – library editing programme EDITOR;

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